## Integrating the Wigner Distribution on subsets of the phase space

NICOLAS LERNER Sorbonne Université

Conference in honour of Jorge HOUNIE, on the occasion of his 75th birthday, January 9-11, 2023, University of São Carlos, Brazil.

January 9, 2023, 10:30-11:20

イロト イポト イヨト イヨト

## 1. Introduction

The Wigner distribution Weyl quantization Integrals of the Wigner distribution

## The Wigner distribution

Let  $u, v \in L^2(\mathbb{R}^n)$ . We define the Wigner distribution of u, v as

イロン イボン イヨン イヨン

## 1. Introduction

The Wigner distribution

## The Wigner distribution

Let  $u, v \in L^2(\mathbb{R}^n)$ . We define the Wigner distribution of u, v as

$$\mathcal{W}(u,v)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z\cdot\xi} u(x+\frac{z}{2})\bar{v}(x-\frac{z}{2})dz.$$

イロン イボン イヨン イヨン

2

The Wigner distribution Weyl quantization Integrals of the Wigner distribution

# 1. Introduction

## The Wigner distribution

Let  $u, v \in L^2(\mathbb{R}^n)$ . We define the Wigner distribution of u, v as

$$\mathcal{W}(u,v)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z\cdot\xi} u(x+\frac{z}{2})\overline{v}(x-\frac{z}{2})dz.$$

We note that  $\mathcal{W}(u, v)$  is the partial Fourier transform (wrt to z) of

$$\mathbb{R}^n \times \mathbb{R}^n \ni (z, x) \mapsto u(x + \frac{z}{2})\overline{v}(x - \frac{z}{2}) = \Omega(u, v)(x, z),$$

(ロ) (部) (注) (注) (注)

2

## 1. Introduction

### The Wigner distribution

Let  $u, v \in L^2(\mathbb{R}^n)$ . We define the Wigner distribution of u, v as

$$\mathcal{W}(u,v)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z\cdot\xi} u(x+\frac{z}{2})\bar{v}(x-\frac{z}{2})dz.$$

The Wigner distribution

We note that  $\mathcal{W}(u, v)$  is the partial Fourier transform (wrt to z) of

$$\mathbb{R}^n \times \mathbb{R}^n \ni (z,x) \mapsto u(x+\frac{z}{2})\overline{v}(x-\frac{z}{2}) = \Omega(u,v)(x,z),$$

so that

$$\|\mathcal{W}(u,v)\|_{L^{2}(\mathbb{R}^{2n})} = \|\Omega(u,v)\|_{L^{2}(\mathbb{R}^{2n})} = \|u\|_{L^{2}(\mathbb{R}^{n})} \|v\|_{L^{2}(\mathbb{R}^{n})}.$$

(ロ) (部) (注) (注) (注)

## 1. Introduction

## The Wigner distribution

Let  $u, v \in L^2(\mathbb{R}^n)$ . We define the Wigner distribution of u, v as

$$\mathcal{W}(u,v)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z \cdot \xi} u(x+\frac{z}{2})\overline{v}(x-\frac{z}{2})dz.$$

The Wigner distribution

We note that  $\mathcal{W}(u, v)$  is the partial Fourier transform (wrt to z) of

$$\mathbb{R}^n \times \mathbb{R}^n \ni (z,x) \mapsto u(x+\frac{z}{2})\overline{v}(x-\frac{z}{2}) = \Omega(u,v)(x,z),$$

so that

$$\|\mathcal{W}(u,v)\|_{L^{2}(\mathbb{R}^{2n})} = \|\Omega(u,v)\|_{L^{2}(\mathbb{R}^{2n})} = \|u\|_{L^{2}(\mathbb{R}^{n})} \|v\|_{L^{2}(\mathbb{R}^{n})}.$$

Moreover,  $\mathcal{W}(u, v)$  belongs to  $\mathscr{S}(\mathbb{R}^{2n})$  when u, v belong to  $\mathscr{S}(\mathbb{R}^{n})$ .

・ロト ・聞ト ・ヨト ・ヨト

1. Introduction	The Wigner distribution
2. Positive results, Examples and Counterexamples	Weyl quantization
3. More results and comments	Integrals of the Wigner distribution

### The Wigner function of u is

$$\mathcal{W}(u,u)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z \cdot \xi} u(x+\frac{z}{2})\bar{u}(x-\frac{z}{2})dz \tag{1}$$

・ロン ・御と ・注と ・注と ……注

#### The Wigner function of u is

$$\mathcal{W}(u,u)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z \cdot \xi} u(x+\frac{z}{2})\bar{u}(x-\frac{z}{2})dz \tag{1}$$

is real-valued and such that

3

$$\iint \mathcal{W}(u,u)(x,\xi)dxd\xi = \|u\|_{L^2(\mathbb{R}^n)}^2,$$

イロン イヨン イヨン イヨン

æ

#### The Wigner function of u is

$$\mathcal{W}(u,u)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z \cdot \xi} u(x+\frac{z}{2})\bar{u}(x-\frac{z}{2})dz \tag{1}$$

is real-valued and such that

3

$$\iint \mathcal{W}(u,u)(x,\xi)dxd\xi = \|u\|_{L^2(\mathbb{R}^n)}^2,$$

but may take negative values: choosing for instance  $u_1(x) = xe^{-\pi x^2}$  on the real line, we get

$$\mathcal{W}(u_1, u_1)(x, \xi) = 2^{1/2} e^{-2\pi (x^2 + \xi^2)} (x^2 + \xi^2 - \frac{1}{4\pi}).$$

イロン イヨン イヨン イヨン

#### The Wigner function of u is

$$\mathcal{W}(u,u)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z \cdot \xi} u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2})dz \tag{1}$$

is real-valued and such that

3

$$\iint \mathcal{W}(u,u)(x,\xi)dxd\xi = \|u\|_{L^2(\mathbb{R}^n)}^2,$$

but may take negative values: choosing for instance  $u_1(x) = xe^{-\pi x^2}$  on the real line, we get

$$\mathcal{W}(u_1, u_1)(x, \xi) = 2^{1/2} e^{-2\pi (x^2 + \xi^2)} (x^2 + \xi^2 - \frac{1}{4\pi}).$$

In fact the real-valued function W(u, u) will take negative values unless u is a Gaussian function, thanks to a theorem due to E. LIEB.

イロン イヨン イヨン イヨン

#### The Wigner function of u is

$$\mathcal{W}(u,u)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z \cdot \xi} u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2})dz \tag{1}$$

is real-valued and such that

3

$$\iint \mathcal{W}(u,u)(x,\xi)dxd\xi = \|u\|_{L^2(\mathbb{R}^n)}^2,$$

but may take negative values: choosing for instance  $u_1(x) = xe^{-\pi x^2}$  on the real line, we get

$$\mathcal{W}(u_1, u_1)(x, \xi) = 2^{1/2} e^{-2\pi (x^2 + \xi^2)} (x^2 + \xi^2 - \frac{1}{4\pi}).$$

In fact the real-valued function  $\mathcal{W}(u, u)$  will take negative values unless u is a Gaussian function, thanks to a theorem due to E. LIEB. As a matter of fact, this range of  $\mathcal{W}(u, u)$  intersecting  $\mathbb{R}_{-}$  for most "pulses" u in  $L^{2}(\mathbb{R}^{n})$  makes rather weird the qualification of  $\mathcal{W}(u, u)$  as a "quasi-probability"

#### The Wigner function of u is

$$\mathcal{W}(u,u)(x,\xi) = \int_{\mathbb{R}^n} e^{-2i\pi z \cdot \xi} u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2})dz \tag{1}$$

is real-valued and such that

3

$$\iint \mathcal{W}(u,u)(x,\xi)dxd\xi = \|u\|_{L^2(\mathbb{R}^n)}^2,$$

but may take negative values: choosing for instance  $u_1(x) = xe^{-\pi x^2}$  on the real line, we get

$$\mathcal{W}(u_1, u_1)(x, \xi) = 2^{1/2} e^{-2\pi (x^2 + \xi^2)} (x^2 + \xi^2 - \frac{1}{4\pi}).$$

In fact the real-valued function  $\mathcal{W}(u, u)$  will take negative values unless u is a Gaussian function, thanks to a theorem due to E. LIEB. As a matter of fact, this range of  $\mathcal{W}(u, u)$  intersecting  $\mathbb{R}_{-}$  for most "pulses" u in  $L^{2}(\mathbb{R}^{n})$  makes rather weird the qualification of  $\mathcal{W}(u, u)$  as a "quasi-probability" (anyhow the emphasis must be on *quasi*, not on *probability*).

・ロン ・回と ・ヨン・

 1. Introduction
 The Win

 2. Positive results, Examples and Counterexamples
 Weyl quint

 3. More results and comments
 Integral

The Wigner distribution Weyl quantization Integrals of the Wigner distribution

We have also by Fourier inversion formula

(b) 
$$u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2}) = \Omega(x,z) = \int \mathcal{W}(u,u)(x,\xi)e^{2i\pi z\cdot\xi}d\xi,$$

イロン イボン イモン イモン 三日

We have also by Fourier inversion formula

(b) 
$$u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2}) = \Omega(x,z) = \int \mathcal{W}(u,u)(x,\xi)e^{2i\pi z\cdot\xi}d\xi,$$

so that, with z = 2x = y, we get the Reconstruction Formula,

$$u(y)\overline{u}(0) = \int \mathcal{W}(u,u)(\frac{y}{2},\xi)e^{2i\pi y\cdot\xi}d\xi$$
(2)

・ロン ・回 と ・ ヨン ・ ヨ

We have also by Fourier inversion formula

4

$$(\flat) \qquad u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2}) = \Omega(x,z) = \int \mathcal{W}(u,u)(x,\xi)e^{2i\pi z\cdot\xi}d\xi,$$

so that, with z = 2x = y, we get the Reconstruction Formula,

$$u(y)\overline{u}(0) = \int \mathcal{W}(u,u)(\frac{y}{2},\xi)e^{2i\pi y\cdot\xi}d\xi$$
(2)

Your radar device is in principle providing you with the **knowledge of** W(u, u) and using (2), you want to reconstruct the unknown pulse u (assumed to be in  $L^2(\mathbb{R}^n)$ ).

We have also by Fourier inversion formula

4

(b) 
$$u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2}) = \Omega(x,z) = \int \mathcal{W}(u,u)(x,\xi)e^{2i\pi z\cdot\xi}d\xi,$$

so that, with z = 2x = y, we get the Reconstruction Formula,

$$u(y)\overline{u}(0) = \int \mathcal{W}(u,u)(\frac{y}{2},\xi)e^{2i\pi y\cdot\xi}d\xi$$
(2)

Your radar device is in principle providing you with the **knowledge of** W(u, u) and using (2), you want to **reconstruct the unknown pulse** u (assumed to be in  $L^2(\mathbb{R}^n)$ ). A typical difficulty: the integral above is in general converging slowly and in particular you should not expect that  $W(u, u)(\frac{y}{2}, \cdot)$  belongs to  $L^1(\mathbb{R}^n)$ .

We have also by Fourier inversion formula

(b) 
$$u(x+\frac{z}{2})\overline{u}(x-\frac{z}{2}) = \Omega(x,z) = \int \mathcal{W}(u,u)(x,\xi)e^{2i\pi z\cdot\xi}d\xi,$$

so that, with z = 2x = y, we get the Reconstruction Formula,

$$u(y)\overline{u}(0) = \int \mathcal{W}(u,u)(\frac{y}{2},\xi)e^{2i\pi y\cdot\xi}d\xi$$
(2)

Your radar device is in principle providing you with the **knowledge of** W(u, u) and using (2), you want to **reconstruct the unknown pulse** u (assumed to be in  $L^2(\mathbb{R}^n)$ ). A typical difficulty: the integral above is in general converging slowly and in particular you should not expect that  $W(u, u)(\frac{y}{2}, \cdot)$  belongs to  $L^1(\mathbb{R}^n)$ .

There are many other properties of the Wigner distribution and it turns out that most of these properties are closely linked to Weyl quantization, named after the German mathematician HERMANN WEYL (1885–1955).

1. Introduction	The Wigner distribution
2. Positive results, Examples and Counterexamples	Weyl quantization
3. More results and comments	Integrals of the Wigner distribution

### Weyl quantization

Let  $a(x,\xi)$  be a Hamiltonian defined on  $\mathbb{R}^n \times \mathbb{R}^n$ . We want to associate an operator to that function.

・ロト・(部)・・ヨト・ヨト ヨー・ショー

#### Weyl quantization

5

Let  $a(x,\xi)$  be a Hamiltonian defined on  $\mathbb{R}^n \times \mathbb{R}^n$ . We want to associate an operator to that function. For instance in one dimension, we have with the standard quantization formulas,

$$\begin{split} \xi \rightsquigarrow D_x &= \frac{1}{i} \frac{d}{dx}, \\ x \rightsquigarrow \text{multiplication by } x \\ x\xi \rightsquigarrow xD_x. \end{split}$$

・ロン ・回 と ・ ヨン ・ ヨ

#### Weyl quantization

5

Let  $a(x,\xi)$  be a Hamiltonian defined on  $\mathbb{R}^n \times \mathbb{R}^n$ . We want to associate an operator to that function. For instance in one dimension, we have with the standard quantization formulas,

 $\xi \rightsquigarrow D_x = \frac{1}{i} \frac{d}{dx},$ x \mapsto multiplication by x x\xeta \mapsto xD\_x.

The latter formula is not quite satisfactory, since we would like to quantize real Hamiltonians into formally self-adjoint operators, and we prefer

$$x\xi \rightsquigarrow \frac{1}{2}(xD_x + D_x x)$$
 ... (indeed selfadjoint).

#### Weyl quantization

Let  $a(x,\xi)$  be a Hamiltonian defined on  $\mathbb{R}^n \times \mathbb{R}^n$ . We want to associate an operator to that function. For instance in one dimension, we have with the standard quantization formulas,

$$\begin{split} \xi &\leadsto D_x = \frac{1}{i} \frac{d}{dx}, \\ x &\leadsto \text{ multiplication by } x \\ x \xi &\leadsto x D_x. \end{split}$$

The latter formula is not quite satisfactory, since we would like to quantize real Hamiltonians into formally self-adjoint operators, and we prefer

$$x\xi \rightsquigarrow \frac{1}{2}(xD_x + D_x x)$$
 ... (indeed selfadjoint).

We want to use

$$(a^{\mathrm{WEYL}}u)(x) = \iint e^{2i\pi(x-y)\cdot\xi}a(\frac{x+y}{2},\xi)u(y)dyd\xi,$$

#### Weyl quantization

Let  $a(x,\xi)$  be a Hamiltonian defined on  $\mathbb{R}^n \times \mathbb{R}^n$ . We want to associate an operator to that function. For instance in one dimension, we have with the standard quantization formulas,

$$\begin{split} \xi &\leadsto D_x = \frac{1}{i} \frac{d}{dx}, \\ x &\leadsto \text{multiplication by } x \\ x \xi &\leadsto x D_x. \end{split}$$

The latter formula is not quite satisfactory, since we would like to quantize real Hamiltonians into formally self-adjoint operators, and we prefer

$$x\xi \rightsquigarrow \frac{1}{2}(xD_x + D_x x)$$
 ... (indeed selfadjoint).

We want to use

$$(a^{\mathrm{WEYL}}u)(x) = \iint e^{2i\pi(x-y)\cdot\xi}a(\frac{x+y}{2},\xi)u(y)dyd\xi,$$

instead of the standard

$$(\operatorname{Op}(a)u)(x) = \iint e^{2i\pi(x-y)\cdot\xi} a(x,\xi)u(y)dyd\xi = \int e^{2i\pi x\cdot\xi} a(x,\xi)\hat{u}(\xi)d\xi.$$

Let a be a tempered distribution on  $\mathbb{R}^n\times\mathbb{R}^n;$  we would like to give a meaning to the integral

$$(a^{\mathrm{WEYL}}u)(x) = \iint e^{2i\pi(x-y)\cdot\xi}a(\frac{x+y}{2},\xi)u(y)dyd\xi$$

イロン イボン イモン イモン 三日

Let a be a tempered distribution on  $\mathbb{R}^n \times \mathbb{R}^n$ ; we would like to give a meaning to the integral

$$(a^{\mathrm{WEYL}}u)(x) = \iint e^{2i\pi(x-y)\cdot\xi}a(\frac{x+y}{2},\xi)u(y)dyd\xi,$$

where  $a^{\text{WEYL}}$  stands for the Weyl quantization of the Hamiltonian "function"  $a(x,\xi)$ . Assuming that  $u, v \in \mathscr{S}(\mathbb{R}^n)$ , and for a minute that  $a \in \mathscr{S}(\mathbb{R}^{2n})$ , we get

・ロン ・回 と ・ ヨン ・ ヨ

Let a be a tempered distribution on  $\mathbb{R}^n \times \mathbb{R}^n$ ; we would like to give a meaning to the integral

$$(a^{\mathrm{WEYL}}u)(x) = \iint e^{2i\pi(x-y)\cdot\xi}a(\frac{x+y}{2},\xi)u(y)dyd\xi,$$

where  $a^{\text{WEYL}}$  stands for the Weyl quantization of the Hamiltonian "function"  $a(x,\xi)$ . Assuming that  $u, v \in \mathscr{S}(\mathbb{R}^n)$ , and for a minute that  $a \in \mathscr{S}(\mathbb{R}^{2n})$ , we get

$$\langle a^{\text{WEYL}}u, v \rangle_{L^{2}(\mathbb{R}^{n})} = \iiint e^{2i\pi (x-y) \cdot \xi} a(\frac{x'}{2}, \xi) u(y) \overline{v}(x) dy d\xi dx = \iint a(x', \xi) \left[ \int e^{-2i\pi z\xi} u(x' + \frac{z}{2}) \overline{v}(x' - \frac{z}{2}) dz \right] dx' d\xi = \langle a, \qquad \mathcal{W}(u, v) \qquad \rangle_{\mathscr{S}'(\mathbb{R}^{2n}), \mathscr{S}(\mathbb{R}^{2n})},$$
  
This is the Wiener distribution

6

・ロン ・回 と ・ ヨン ・ ヨ

Let a be a tempered distribution on  $\mathbb{R}^n \times \mathbb{R}^n$ ; we would like to give a meaning to the integral

$$(a^{\mathrm{WEYL}}u)(x) = \iint e^{2i\pi(x-y)\cdot\xi}a(\frac{x+y}{2},\xi)u(y)dyd\xi,$$

where  $a^{\text{WEYL}}$  stands for the Weyl quantization of the Hamiltonian "function"  $a(x,\xi)$ . Assuming that  $u, v \in \mathscr{S}(\mathbb{R}^n)$ , and for a minute that  $a \in \mathscr{S}(\mathbb{R}^{2n})$ , we get

$$\langle a^{\mathrm{WEYL}}u, v \rangle_{L^{2}(\mathbb{R}^{n})} = \iiint e^{2i\pi (x-y) \cdot \xi} a(\frac{x'}{2}, \xi) u(y) \bar{v}(x) dy d\xi dx = \iint a(x', \xi) \left[ \int e^{-2i\pi z\xi} u(x' + \frac{z}{2}) \bar{v}(x' - \frac{z}{2}) dz \right] dx' d\xi = \langle a, \qquad \mathcal{W}(u, v) \qquad \rangle_{\mathscr{S}'(\mathbb{R}^{2n}), \mathscr{S}(\mathbb{R}^{2n})},$$
  
This is the Wigner distribution

providing a meaning for  $a^{\text{WEYL}}$  for  $a \in \mathscr{S}'(\mathbb{R}^{2n})$  as an operator from  $\mathscr{S}(\mathbb{R}^n)$  to  $\mathscr{S}'(\mathbb{R}^n)$ , with

6

Let a be a tempered distribution on  $\mathbb{R}^n \times \mathbb{R}^n$ ; we would like to give a meaning to the integral

$$(a^{\mathrm{WEYL}}u)(x) = \iint e^{2i\pi(x-y)\cdot\xi}a(\frac{x+y}{2},\xi)u(y)dyd\xi,$$

where  $a^{\text{WEYL}}$  stands for the Weyl quantization of the Hamiltonian "function"  $a(x,\xi)$ . Assuming that  $u, v \in \mathscr{S}(\mathbb{R}^n)$ , and for a minute that  $a \in \mathscr{S}(\mathbb{R}^{2n})$ , we get

$$\langle a^{\mathrm{W}\mathrm{E}\mathrm{YL}}u, v \rangle_{L^{2}(\mathbb{R}^{n})} = \iiint e^{2i\pi \overline{(x-y)} \cdot \xi} a(\overline{\frac{x'}{2}}, \xi) u(y) \overline{v}(x) dy d\xi dx = \iint a(x', \xi) \left[ \int e^{-2i\pi z\xi} u(x' + \frac{z}{2}) \overline{v}(x' - \frac{z}{2}) dz \right] dx' d\xi = \langle a, \underbrace{\mathcal{W}(u, v)}_{\text{This is the Wigner distribution}} \rangle_{\mathscr{S}'(\mathbb{R}^{2n}), \mathscr{S}(\mathbb{R}^{2n})},$$

providing a meaning for  $a^{\text{WEYL}}$  for  $a \in \mathscr{S}'(\mathbb{R}^{2n})$  as an operator from  $\mathscr{S}(\mathbb{R}^n)$  to  $\mathscr{S}'(\mathbb{R}^n)$ , with

6

$$\langle a^{\mathrm{WeyL}}u, \bar{v} \rangle_{\mathscr{S}'(\mathbb{R}^n), \mathscr{S}(\mathbb{R}^n)} = \langle a, \mathcal{W}(u, v) \rangle_{\mathscr{S}'(\mathbb{R}^{2n}), \mathscr{S}(\mathbb{R}^{2n})}.$$

1. Introduction	
2. Positive results, Examples and Counterexamples	Weyl quantization
3. More results and comments	

Here is a sufficient condition for  $L^2(\mathbb{R}^n)$  boundedness of  $a^{WEYL}$ : Let a be a tempered distribution on  $\mathbb{R}^{2n}$ . Then we have

$$\|a^{\text{WEYL}}\|_{\mathcal{B}(L^{2}(\mathbb{R}^{n}))} \leq \min(2^{n} \|a\|_{L^{1}(\mathbb{R}^{2n})}, \|\hat{a}\|_{L^{1}(\mathbb{R}^{2n})}).$$
(3)

・ロト・(部)・・ヨト・ヨト ヨー・ショー

1. Introduction	
2. Positive results, Examples and Counterexamples	Weyl quantization
3. More results and comments	

Here is a sufficient condition for  $L^2(\mathbb{R}^n)$  boundedness of  $a^{\text{WEYL}}$ : Let a be a tempered distribution on  $\mathbb{R}^{2n}$ . Then we have

$$\|a^{\text{WEYL}}\|_{\mathcal{B}(L^{2}(\mathbb{R}^{n}))} \leq \min(2^{n} \|a\|_{L^{1}(\mathbb{R}^{2n})}, \|\hat{a}\|_{L^{1}(\mathbb{R}^{2n})}).$$
(3)

In fact for  $u, v \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\langle a^{\mathrm{WEYL}}u,v\rangle_{L^2(\mathbb{R}^n)} = \iiint a(x,\xi)u(2x-y)\overline{v}(y)e^{-4i\pi(x-y)\cdot\xi}2^ndydxd\xi,$$

・ロト・日本・モト・モー ショー シック

Here is a sufficient condition for  $L^2(\mathbb{R}^n)$  boundedness of  $a^{WEYL}$ : Let a be a tempered distribution on  $\mathbb{R}^{2n}$ . Then we have

$$\|a^{W_{EYL}}\|_{\mathcal{B}(L^{2}(\mathbb{R}^{n}))} \leq \min(2^{n} \|a\|_{L^{1}(\mathbb{R}^{2n})}, \|\hat{a}\|_{L^{1}(\mathbb{R}^{2n})}).$$
(3)

In fact for  $u, v \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\langle a^{\mathrm{WeyL}}u,v\rangle_{L^2(\mathbb{R}^n)} = \iiint a(x,\xi)u(2x-y)\overline{v}(y)e^{-4i\pi(x-y)\cdot\xi}2^n dydxd\xi,$$

so that defining for  $(x,\xi)\in\mathbb{R}^{2n}$  the operator  $\sigma_{x,\xi}$  by

$$(\sigma_{x,\xi}u)(y) = u(2x-y)e^{-4i\pi(x-y)\cdot\xi},$$

Here is a sufficient condition for  $L^2(\mathbb{R}^n)$  boundedness of  $a^{WEYL}$ : Let a be a tempered distribution on  $\mathbb{R}^{2n}$ . Then we have

$$\|a^{W_{EYL}}\|_{\mathcal{B}(L^{2}(\mathbb{R}^{n}))} \leq \min(2^{n} \|a\|_{L^{1}(\mathbb{R}^{2n})}, \|\hat{a}\|_{L^{1}(\mathbb{R}^{2n})}).$$
(3)

In fact for  $u, v \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\langle a^{\mathrm{WeyL}}u,v\rangle_{L^2(\mathbb{R}^n)} = \iiint a(x,\xi)u(2x-y)\overline{v}(y)e^{-4i\pi(x-y)\cdot\xi}2^n dydxd\xi,$$

so that defining for  $(x,\xi)\in \mathbb{R}^{2n}$  the operator  $\sigma_{x,\xi}$  by

$$(\sigma_{x,\xi}u)(y) = u(2x-y)e^{-4i\pi(x-y)\cdot\xi},$$

we see that  $\sigma_{x,\xi}$  (phase symmetry) is unitary and self-adjoint and

$$a^{\text{WEYL}} = 2^n \iint a(x,\xi)\sigma_{x,\xi}dxd\xi = \iint \widehat{a}(\eta,y)e^{2i\pi(\eta\cdot x+y\cdot D_x)}dyd\eta, \tag{4}$$

・ロン ・回 と ・ ヨン ・ ヨ

Here is a sufficient condition for  $L^2(\mathbb{R}^n)$  boundedness of  $a^{WEYL}$ : Let a be a tempered distribution on  $\mathbb{R}^{2n}$ . Then we have

$$\|a^{W_{EYL}}\|_{\mathcal{B}(L^{2}(\mathbb{R}^{n}))} \leq \min(2^{n} \|a\|_{L^{1}(\mathbb{R}^{2n})}, \|\hat{a}\|_{L^{1}(\mathbb{R}^{2n})}).$$
(3)

In fact for  $u, v \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\langle a^{\mathrm{WeyL}}u,v\rangle_{L^2(\mathbb{R}^n)} = \iiint a(x,\xi)u(2x-y)\overline{v}(y)e^{-4i\pi(x-y)\cdot\xi}2^n dydxd\xi,$$

so that defining for  $(x,\xi)\in\mathbb{R}^{2n}$  the operator  $\sigma_{x,\xi}$  by

$$(\sigma_{x,\xi}u)(y) = u(2x-y)e^{-4i\pi(x-y)\cdot\xi},$$

we see that  $\sigma_{x,\xi}$  (phase symmetry) is unitary and self-adjoint and

$$a^{\text{WEYL}} = 2^n \iint a(x,\xi)\sigma_{x,\xi}dxd\xi = \iint \widehat{a}(\eta,y)e^{2i\pi(\eta\cdot x+y\cdot D_x)}dyd\eta, \tag{4}$$

proving the estimates of (3).

Here is a sufficient condition for  $L^2(\mathbb{R}^n)$  boundedness of  $a^{WEYL}$ : Let a be a tempered distribution on  $\mathbb{R}^{2n}$ . Then we have

$$\|a^{W_{EYL}}\|_{\mathcal{B}(L^{2}(\mathbb{R}^{n}))} \leq \min(2^{n} \|a\|_{L^{1}(\mathbb{R}^{2n})}, \|\hat{a}\|_{L^{1}(\mathbb{R}^{2n})}).$$
(3)

In fact for  $u, v \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\langle a^{\mathrm{WeyL}}u,v\rangle_{L^2(\mathbb{R}^n)} = \iiint a(x,\xi)u(2x-y)\overline{v}(y)e^{-4i\pi(x-y)\cdot\xi}2^n dy dx d\xi,$$

so that defining for  $(x,\xi)\in\mathbb{R}^{2n}$  the operator  $\sigma_{x,\xi}$  by

$$(\sigma_{x,\xi}u)(y) = u(2x-y)e^{-4i\pi(x-y)\cdot\xi},$$

we see that  $\sigma_{x,\xi}$  (phase symmetry) is unitary and self-adjoint and

$$a^{\text{WEYL}} = 2^n \iint a(x,\xi)\sigma_{x,\xi}dxd\xi = \iint \widehat{a}(\eta,y)e^{2i\pi(\eta\cdot x+y\cdot D_x)}dyd\eta, \tag{4}$$

proving the estimates of (3). As a consequence, we obtain that

$$(a^{\text{WeyL}})^* = (\overline{a})^{\text{WeyL}}$$
, so that for *a* real-valued,  $(a^{\text{WeyL}})^* = a^{\text{WeyL}}$ .

Another important property of the Weyl quantization is its symplectic covariance.

イロン イボン イモン イモン 三日

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n, \mathbb{R})$  is the subgroup of  $S \in Sl(2n, \mathbb{R})$  such that

 $\forall X, Y \in \mathbb{R}^{2n}, [SX, SY] = [X, Y],$  i.e.  $S^* \sigma S = \sigma,$ 

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n, \mathbb{R})$  is the subgroup of  $S \in Sl(2n, \mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}.$$

(ロ) (部) (注) (注) (注)
Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n,\mathbb{R})$  is the subgroup of  $S \in Sl(2n,\mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}$$

Now for  $S \in Sp(n, \mathbb{R})$ , the operator

$$(a \circ S)^{\mathrm{WEYL}} = \mathcal{M}^* a^{\mathrm{WEYL}} \mathcal{M}_{\mathrm{S}}$$

イロン イヨン イヨン イヨン

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n,\mathbb{R})$  is the subgroup of  $S \in Sl(2n,\mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}$$

Now for  $S \in Sp(n, \mathbb{R})$ , the operator

8

$$(a \circ S)^{WEYL} = \mathcal{M}^* a^{WEYL} \mathcal{M},$$

where  $\mathcal{M}$  belongs to the metaplectic group, which is a group of unitary transformations of  $L^2(\mathbb{R}^n)$ .

イロン イボン イヨン イヨン 三日

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n, \mathbb{R})$  is the subgroup of  $S \in Sl(2n, \mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}$$

Now for  $S \in Sp(n, \mathbb{R})$ , the operator

8

$$(a \circ S)^{\mathrm{WEYL}} = \mathcal{M}^* a^{\mathrm{WEYL}} \mathcal{M},$$

where  $\mathcal{M}$  belongs to the metaplectic group, which is a group of unitary transformations of  $L^2(\mathbb{R}^n)$ . This translates for the Wigner distribution as

$$\mathcal{W}(\mathcal{M}u, \mathcal{M}v) = \mathcal{W}(u, v) \circ S^{-1}.$$
(5)

・ロト ・聞ト ・ヨト ・ヨト

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n,\mathbb{R})$  is the subgroup of  $S \in Sl(2n,\mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}$$

Now for  $S \in Sp(n, \mathbb{R})$ , the operator

8

$$(a \circ S)^{ ext{WEYL}} = \mathcal{M}^* a^{ ext{WEYL}} \mathcal{M},$$

where  $\mathcal{M}$  belongs to the metaplectic group, which is a group of unitary transformations of  $L^2(\mathbb{R}^n)$ . This translates for the Wigner distribution as

$$\mathcal{W}(\mathcal{M}u,\mathcal{M}v) = \mathcal{W}(u,v) \circ S^{-1}.$$
(5)

3

$$(a(T^{-1}x, {}^tT\xi))^{W_{\mathrm{EYL}}} = \mathcal{M}_T^* a^{W_{\mathrm{EYL}}} \mathcal{M}_T, \quad \text{with} \quad (\mathcal{M}_T u)(x) = (\det T)^{1/2} u(Tx),$$

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n,\mathbb{R})$  is the subgroup of  $S \in Sl(2n,\mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}$$

Now for  $S \in Sp(n, \mathbb{R})$ , the operator

8

$$(a \circ S)^{\mathrm{WEYL}} = \mathcal{M}^* a^{\mathrm{WEYL}} \mathcal{M},$$

where  $\mathcal{M}$  belongs to the metaplectic group, which is a group of unitary transformations of  $L^2(\mathbb{R}^n)$ . This translates for the Wigner distribution as

$$\mathcal{W}(\mathcal{M}u,\mathcal{M}v) = \mathcal{W}(u,v) \circ S^{-1}.$$
(5)

3

$$\begin{split} \left(a(T^{-1}x,{}^{t}T\xi)\right)^{\text{WEYL}} &= \mathcal{M}_{T}^{*}a^{\text{WEYL}}\mathcal{M}_{T}, \quad \text{with} \quad (\mathcal{M}_{T}u)(x) = (\det T)^{1/2}u(Tx), \\ \left(a(\xi,-x)\right)^{\text{WEYL}} &= \mathcal{F}^{*}a^{\text{WEYL}}\mathcal{F}, \quad \text{where} \quad \mathcal{F} \text{ is the Fourier transformation,} \end{split}$$

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n,\mathbb{R})$  is the subgroup of  $S \in Sl(2n,\mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}$$

Now for  $S \in Sp(n, \mathbb{R})$ , the operator

8

$$({\it a} \circ {\it S})^{ ext{WEYL}} = {\cal M}^* {\it a}^{ ext{WEYL}} {\cal M},$$

where  $\mathcal{M}$  belongs to the metaplectic group, which is a group of unitary transformations of  $L^2(\mathbb{R}^n)$ . This translates for the Wigner distribution as

$$\mathcal{W}(\mathcal{M}u,\mathcal{M}v)=\mathcal{W}(u,v)\circ S^{-1}.$$
(5)

3

$$\begin{aligned} \left(a(T^{-1}x,{}^{t}T\xi)\right)^{\mathrm{WEYL}} &= \mathcal{M}_{T}^{*}a^{\mathrm{WEYL}}\mathcal{M}_{T}, \quad \text{with} \quad (\mathcal{M}_{T}u)(x) = (\det T)^{1/2}u(Tx), \\ \left(a(\xi,-x)\right)^{\mathrm{WEYL}} &= \mathcal{F}^{*}a^{\mathrm{WEYL}}\mathcal{F}, \quad \text{where} \quad \mathcal{F} \text{ is the Fourier transformation,} \\ \left(a(x,\xi+Ax)\right)^{\mathrm{WEYL}} &= \mathcal{L}_{A}^{*}a^{\mathrm{WEYL}}\mathcal{L}_{A}, \quad \text{where} \quad (\mathcal{L}_{A}v)(x) = e^{i\pi\langle Ax,x\rangle}v(x), \quad {}^{t}A = A \end{aligned}$$

Another important property of the Weyl quantization is its symplectic covariance. The symplectic group  $Sp(n, \mathbb{R})$  is the subgroup of  $S \in Sl(2n, \mathbb{R})$  such that

$$\forall X, Y \in \mathbb{R}^{2n}, \quad [SX, SY] = [X, Y], \quad \text{i.e.} \quad S^* \sigma S = \sigma,$$

with

$$\sigma = \begin{pmatrix} 0 & I_n \\ -I_n & 0 \end{pmatrix}$$

Now for  $S \in Sp(n, \mathbb{R})$ , the operator

8

$$(a \circ S)^{WEYL} = \mathcal{M}^* a^{WEYL} \mathcal{M},$$

where  $\mathcal{M}$  belongs to the metaplectic group, which is a group of unitary transformations of  $L^2(\mathbb{R}^n)$ . This translates for the Wigner distribution as

$$\mathcal{W}(\mathcal{M}u,\mathcal{M}v)=\mathcal{W}(u,v)\circ S^{-1}.$$
(5)

3

$$\begin{aligned} & \left(a(T^{-1}x,{}^{t}T\xi)\right)^{\mathrm{WEYL}} = \mathcal{M}_{T}^{*}a^{\mathrm{WEYL}}\mathcal{M}_{T}, & \text{with} \quad (\mathcal{M}_{T}u)(x) = (\det T)^{1/2}u(Tx), \\ & \left(a(\xi,-x)\right)^{\mathrm{WEYL}} = \mathcal{F}^{*}a^{\mathrm{WEYL}}\mathcal{F}, & \text{where} \quad \mathcal{F} \text{ is the Fourier transformation,} \\ & \left(a(x,\xi+Ax)\right)^{\mathrm{WEYL}} = \mathcal{L}_{A}^{*}a^{\mathrm{WEYL}}\mathcal{L}_{A}, & \text{where} \quad (\mathcal{L}_{A}v)(x) = e^{i\pi\langle Ax,x\rangle}v(x), \quad {}^{t}A = A, \\ & \left(a(x-C\xi,\xi)\right)^{\mathrm{WEYL}} = \mathcal{R}_{C}^{*}a^{\mathrm{WEYL}}\mathcal{R}_{C}, & \text{where} \quad (\mathcal{R}_{C}v)(x) = (e^{i\pi\langle CD,D\rangle}v)(x), \quad {}^{t}C = C. \end{aligned}$$

#### Integrals of the Wigner distribution

• Let *E* be a measurable subset with finite Lebesgue measure of the phase space  $\mathbb{R}^n \times \mathbb{R}^n$  and let  $\mathbf{1}_E$  be the indicator function of the set *E*. Then the operator with Weyl symbol  $\mathbf{1}_E$  is bounded (see (3)) self-adjoint (see (4)) on  $L^2(\mathbb{R}^n)$  and for any  $u \in L^2(\mathbb{R}^n)$ , we have

$$\langle \mathbf{1}_{E}^{\mathrm{WEYL}} u, u \rangle_{L^{2}(\mathbb{R}^{n})} = \iint_{E} \mathcal{W}(u, u)(x, \xi) dx d\xi.$$
(6)

#### Integrals of the Wigner distribution

9

• Let *E* be a measurable subset with finite Lebesgue measure of the phase space  $\mathbb{R}^n \times \mathbb{R}^n$  and let  $\mathbf{1}_E$  be the indicator function of the set *E*. Then the operator with Weyl symbol  $\mathbf{1}_E$  is bounded (see (3)) self-adjoint (see (4)) on  $L^2(\mathbb{R}^n)$  and for any  $u \in L^2(\mathbb{R}^n)$ , we have

$$\langle \mathbf{1}_{E}^{\mathrm{WeyL}} u, u \rangle_{L^{2}(\mathbb{R}^{n})} = \iint_{E} \mathcal{W}(u, u)(x, \xi) dx d\xi.$$
(6)

• The above formula provides the following equivalence:

#### Integrals of the Wigner distribution

9

• Let *E* be a measurable subset with finite Lebesgue measure of the phase space  $\mathbb{R}^n \times \mathbb{R}^n$  and let  $\mathbf{1}_E$  be the indicator function of the set *E*. Then the operator with Weyl symbol  $\mathbf{1}_E$  is bounded (see (3)) self-adjoint (see (4)) on  $L^2(\mathbb{R}^n)$  and for any  $u \in L^2(\mathbb{R}^n)$ , we have

$$\langle \mathbf{1}_{E}^{\mathrm{WEYL}} u, u \rangle_{L^{2}(\mathbb{R}^{n})} = \iint_{E} \mathcal{W}(u, u)(x, \xi) dx d\xi.$$
(6)

• The above formula provides the following equivalence:

$$\left\{ \forall u \in L^2(\mathbb{R}^n), \iint_E \mathcal{W}(u,u)(x,\xi) dx d\xi \leq \lambda \|u\|_{L^2(\mathbb{R}^n)}^2 \right\} \Longleftrightarrow \mathbf{1}_E^{\mathrm{WeyL}} \leq \lambda.$$

#### Integrals of the Wigner distribution

• Let *E* be a measurable subset with finite Lebesgue measure of the phase space  $\mathbb{R}^n \times \mathbb{R}^n$  and let  $\mathbf{1}_E$  be the indicator function of the set *E*. Then the operator with Weyl symbol  $\mathbf{1}_E$  is bounded (see (3)) self-adjoint (see (4)) on  $L^2(\mathbb{R}^n)$  and for any  $u \in L^2(\mathbb{R}^n)$ , we have

$$\langle \mathbf{1}_{E}^{\mathrm{WEYL}} u, u \rangle_{L^{2}(\mathbb{R}^{n})} = \iint_{E} \mathcal{W}(u, u)(x, \xi) dx d\xi.$$
(6)

• The above formula provides the following equivalence:

$$\left\{ \forall u \in L^2(\mathbb{R}^n), \iint_E \mathcal{W}(u, u)(x, \xi) dx d\xi \leq \lambda \|u\|_{L^2(\mathbb{R}^n)}^2 \right\} \Longleftrightarrow \mathbf{1}_E^{\text{Weyl}} \leq \lambda.$$

• Our main question: is it possible to determine

9

$$\lambda_{+}(E) = \sup\{ \texttt{spectrum } \mathbf{1}_{E}^{\text{Weyl}} \}$$

for a given set E or for a class of subsets enjoying specific properties?

#### Integrals of the Wigner distribution

• Let *E* be a measurable subset with finite Lebesgue measure of the phase space  $\mathbb{R}^n \times \mathbb{R}^n$  and let  $\mathbf{1}_E$  be the indicator function of the set *E*. Then the operator with Weyl symbol  $\mathbf{1}_E$  is bounded (see (3)) self-adjoint (see (4)) on  $L^2(\mathbb{R}^n)$  and for any  $u \in L^2(\mathbb{R}^n)$ , we have

$$\langle \mathbf{1}_{E}^{\mathrm{WEYL}} u, u \rangle_{L^{2}(\mathbb{R}^{n})} = \iint_{E} \mathcal{W}(u, u)(x, \xi) dx d\xi.$$
(6)

• The above formula provides the following equivalence:

$$\left\{ \forall u \in L^2(\mathbb{R}^n), \iint_E \mathcal{W}(u, u)(x, \xi) dx d\xi \leq \lambda \|u\|_{L^2(\mathbb{R}^n)}^2 \right\} \Longleftrightarrow \mathbf{1}_E^{\text{Weyl}} \leq \lambda.$$

• Our main question: is it possible to determine

9

$$\lambda_+(E) = \sup\{ \texttt{spectrum } \mathbf{1}_E^{_{ ext{Weyl}}} \}$$

for a given set E or for a class of subsets enjoying specific properties? We could raise the same question for

$$\lambda_{-}(E) = \inf\{ \texttt{spectrum } \mathbf{1}_{E}^{\text{Weyl}} \}$$

Positive results Convex subsets The quarter-plane

# 2. Positive results, Examples and Counterexamples

**Positive results** 

A trivial example.

・ロト ・回ト ・ヨト ・ヨト

Positive results Convex subsets The quarter-plane

## 2. Positive results, Examples and Counterexamples

### **Positive results**

A trivial example. If *E* is equal to a half-space  $E = \{(x, \xi) \in \mathbb{R}^{2n}, L(x, \xi) \ge 0\},\$ 

イロン イヨン イヨン イヨン

Positive results Convex subsets The quarter-plane

## 2. Positive results, Examples and Counterexamples

### **Positive results**

A trivial example. If E is equal to a half-space  $E = \{(x,\xi) \in \mathbb{R}^{2n}, L(x,\xi) \ge 0\}$ , where L is a linear form, we can find (if  $L \ne 0$ ) a unitary operator  $\mathcal{M}$  on  $L^2(\mathbb{R}^n)$  such that, with  $H = \mathbf{1}_{\mathbb{R}_+}$ ,

A D A A B A A B A A B A

Positive results Convex subsets The quarter-plane

## 2. Positive results, Examples and Counterexamples

### **Positive results**

A trivial example. If *E* is equal to a half-space  $E = \{(x, \xi) \in \mathbb{R}^{2n}, L(x, \xi) \ge 0\}$ , where *L* is a linear form, we can find (if  $L \ne 0$ ) a unitary operator  $\mathcal{M}$  on  $L^2(\mathbb{R}^n)$  such that, with  $H = \mathbf{1}_{\mathbb{R}_+}$ ,

$$\mathbf{1}_{E}^{\mathrm{WEYL}} = \left(H(L)\right)^{\mathrm{WEYL}} = \mathcal{M}^{*}H_{1}\mathcal{M},$$

A D A A B A A B A A B A

10

Positive results Convex subsets The quarter-plane

## 2. Positive results, Examples and Counterexamples

### **Positive results**

A trivial example. If *E* is equal to a half-space  $E = \{(x, \xi) \in \mathbb{R}^{2n}, L(x, \xi) \ge 0\}$ , where *L* is a linear form, we can find (if  $L \ne 0$ ) a unitary operator  $\mathcal{M}$  on  $L^2(\mathbb{R}^n)$  such that, with  $H = \mathbf{1}_{\mathbb{R}_+}$ ,

$$\mathbf{1}_{E}^{\mathrm{Weyl}} = \left(H(L)\right)^{\mathrm{Weyl}} = \mathcal{M}^{*}H_{1}\mathcal{M},$$

where  $H_1$  is the operator of multiplication by  $H(y_1)$ , which is an orthogonal projection (thus has norm 1): this is a consequence of the symplectic covariance properties of the Wigner distribution given above (see (5)).

Positive results Convex subsets The quarter-plane

## 2. Positive results, Examples and Counterexamples

### **Positive results**

A trivial example. If *E* is equal to a half-space  $E = \{(x, \xi) \in \mathbb{R}^{2n}, L(x, \xi) \ge 0\}$ , where *L* is a linear form, we can find (if  $L \neq 0$ ) a unitary operator  $\mathcal{M}$  on  $L^2(\mathbb{R}^n)$  such that, with  $H = \mathbf{1}_{\mathbb{R}_+}$ ,

$$\mathbf{1}_{E}^{\mathrm{Weyl}} = \left(H(L)\right)^{\mathrm{Weyl}} = \mathcal{M}^{*}H_{1}\mathcal{M},$$

where  $H_1$  is the operator of multiplication by  $H(y_1)$ , which is an orthogonal projection (thus has norm 1): this is a consequence of the symplectic covariance properties of the Wigner distribution given above (see (5)). We have here

 $\mathbf{1}_{\mathcal{E}}^{^{\mathrm{WEYL}}} \leq 1 \quad \text{and in fact} \quad \lambda_+(\mathcal{E}) = \sup\{\texttt{spectrum } \mathbf{1}_{\mathcal{E}}^{^{\mathrm{WEYL}}}\} = 1.$ 

Two highly non-trivial examples.

イロン イボン イヨン イヨン 三日

**Two highly non-trivial examples.** The above inequality is true for two-dimensional Euclidean disks and follows from a precise study of P. FLANDRIN:

イロン イヨン イヨン イヨン

**Two highly non-trivial examples.** The above inequality is true for two-dimensional Euclidean disks and follows from a precise study of P. FLANDRIN: for  $a \in \mathbb{R}_+$ , defining

$$D_a = \{(x,\xi) \in \mathbb{R}^{2n}, |x|^2 + |\xi|^2 \le rac{a}{2\pi}\},$$

we have for n = 1,

イロン イヨン イヨン イヨン

**Two highly non-trivial examples.** The above inequality is true for two-dimensional Euclidean disks and follows from a precise study of P. FLANDRIN: for  $a \in \mathbb{R}_+$ , defining

$$D_a = \{(x,\xi) \in \mathbb{R}^{2n}, |x|^2 + |\xi|^2 \le \frac{a}{2\pi}\},\$$

we have for n = 1,

$$\iint_{D_a} W(u, u)(x, \xi) dx d\xi \le (1 - e^{-a}) \|u\|_{L^2(\mathbb{R})}^2, \tag{7}$$

for any  $u \in L^2(\mathbb{R})$ .

イロン イヨン イヨン イヨン

**Two highly non-trivial examples.** The above inequality is true for two-dimensional Euclidean disks and follows from a precise study of P. FLANDRIN: for  $a \in \mathbb{R}_+$ , defining

$$D_a = \{(x,\xi) \in \mathbb{R}^{2n}, |x|^2 + |\xi|^2 \le \frac{a}{2\pi}\},\$$

we have for n = 1,

$$\iint_{D_a} W(u,u)(x,\xi) dx d\xi \le (1-e^{-a}) \|u\|_{L^2(\mathbb{R})}^2, \tag{7}$$

for any  $u \in L^2(\mathbb{R})$ . We have

$$\mathbf{1}_{D_{a}}^{\mathrm{WEYL}} \leq 1-e^{-a} \quad ext{and even} \quad \lambda_{+}(D_{a})=1-e^{-a}.$$

イロン イヨン イヨン イヨン

**Two highly non-trivial examples.** The above inequality is true for two-dimensional Euclidean disks and follows from a precise study of P. FLANDRIN: for  $a \in \mathbb{R}_+$ , defining

$$D_a = \{(x,\xi) \in \mathbb{R}^{2n}, |x|^2 + |\xi|^2 \le \frac{a}{2\pi}\},\$$

we have for n = 1,

$$\iint_{D_a} W(u,u)(x,\xi) dx d\xi \le (1-e^{-a}) \|u\|_{L^2(\mathbb{R})}^2, \tag{7}$$

for any  $u \in L^2(\mathbb{R})$ . We have

$$\mathbf{1}_{D_a}^{ ext{WEYL}} \leq 1-e^{-a} \quad ext{and even} \quad \lambda_+(D_a) = 1-e^{-a}.$$

The results for the disk in two dimensions are readily extendable to polydisks by tensorisation.

 1. Introduction
 Positive results, Examples and Counterexamples
 Positive results
 Convex subsets

 3. More results and comments
 The quarter-pla
 The guarter-pla

A non-trivial matter was to extend this study to 2n-dimensional Euclidean balls, a task performed by E. LIEB and Y. OSTROVER, who provided the case where E is chosen as an Euclidean ball.

イロン イヨン イヨン イヨン

 1. Introduction
 Positive results, Examples and Counterexamples
 Positive results
 Convex subsets

 3. More results and comments
 The quarter-pla
 The guarter-pla
 The guarter-pla

A non-trivial matter was to extend this study to 2*n*-dimensional Euclidean balls, a task performed by E. LIEB and Y. OSTROVER, who provided the case where *E* is chosen as an Euclidean ball. As for the 1*D* argument, a subtle inequality on Laguerre polynomials provides a proof of the estimate for  $n \ge 1$ ,

イロン イヨン イヨン イヨン

A non-trivial matter was to extend this study to 2*n*-dimensional Euclidean balls, a task performed by E. LIEB and Y. OSTROVER, who provided the case where *E* is chosen as an Euclidean ball. As for the 1*D* argument, a subtle inequality on Laguerre polynomials provides a proof of the estimate for  $n \ge 1$ ,

$$\iint_{D_a} W(u,u)(x,\xi) dx d\xi \le \left(1 - \frac{1}{(n-1)!} \int_a^{+\infty} e^{-t} t^{n-1} dt\right) \|u\|_{L^2(\mathbb{R}^n)}^2, \qquad (8)$$

for any  $u \in L^2(\mathbb{R}^n)$  and we have

$$\sup \bigl( \text{spectrum } \mathbf{1}_{D_a}^{\text{WEYL}} \bigr) = \lambda_+(D_a) = 1 - \frac{\Gamma(n,a)}{\Gamma(n)} = 1 - \frac{1}{\Gamma(n)} \int_a^{+\infty} e^{-t} t^{n-1} dt$$

・ロト ・聞ト ・ヨト ・ヨト

A non-trivial matter was to extend this study to 2n-dimensional Euclidean balls, a task performed by E. LIEB and Y. OSTROVER, who provided the case where E is chosen as an Euclidean ball. As for the 1D argument, a subtle inequality on Laguerre polynomials provides a proof of the estimate for  $n \ge 1$ ,

$$\iint_{D_a} W(u,u)(x,\xi) dx d\xi \le \left(1 - \frac{1}{(n-1)!} \int_a^{+\infty} e^{-t} t^{n-1} dt\right) \|u\|_{L^2(\mathbb{R}^n)}^2, \tag{8}$$

for any  $u \in L^2(\mathbb{R}^n)$  and we have

$$\sup \bigl( \texttt{spectrum } 1_{D_a}^{\text{WEYL}} \bigr) = \lambda_+(D_a) = 1 - \frac{\Gamma(n,a)}{\Gamma(n)} = 1 - \frac{1}{\Gamma(n)} \int_a^{+\infty} e^{-t} t^{n-1} dt$$

We have with  $\psi_k$  standing for the Hermite function at level k in one dimension

$$\mathcal{W}(\psi_k,\psi_k)(x,\xi) = (-1)^k 2e^{-2\pi(x^2+\xi^2)}L_k\big(4\pi(x^2+\xi^2)\big), \quad L_k \text{ is the Laguerre polynomial}$$

The Laguerre polynomials  $\{L_k\}_{k\in\mathbb{N}}$  are defined by

$$L_k(x) = e^x \frac{1}{k!} \left(\frac{d}{dx}\right)^k \left\{x^k e^{-x}\right\} = \left(\frac{d}{dx} - 1\right)^k \left\{\frac{x^k}{k!}\right\}.$$

A result due to E. Feldheim in 1940 states that

12

$$\forall k \in \mathbb{N}, \forall x \geq 0, \quad \sum_{0 \leq l \leq k} (-1)^l L_l(x) \geq 0.$$

(日本)(日本)(日本)

**Convex subsets.** The previous examples were (very particular) convex subsets of the phase space (half-spaces, discs, Euclidean balls).

イロン イボン イヨン イヨン 三日

**Convex subsets.** The previous examples were (very particular) convex subsets of the phase space (half-spaces, discs, Euclidean balls). It is rather natural to try an investigation of more general convex subsets.

イロン イヨン イヨン イヨン

**Convex subsets.** The previous examples were (very particular) convex subsets of the phase space (half-spaces, discs, Euclidean balls). It is rather natural to try an investigation of more general convex subsets. Let *C* be a convex bounded subset of  $\mathbb{R}^{2n}$  and  $\mathbf{1}_C$  be the indicator function of *C*. We have

・ロト ・回ト ・ヨト ・ヨト

**Convex subsets.** The previous examples were (very particular) convex subsets of the phase space (half-spaces, discs, Euclidean balls). It is rather natural to try an investigation of more general convex subsets. Let *C* be a convex bounded subset of  $\mathbb{R}^{2n}$  and  $\mathbf{1}_C$  be the indicator function of *C*. We have

$$\forall u \in L^2(\mathbb{R}^n), \quad \iint_C \mathcal{W}(u, u)(x, \xi) dx d\xi \leq \lambda_+(C) \|u\|_{L^2(\mathbb{R}^n)}^2.$$

・ロト ・回ト ・ヨト ・ヨト

**Convex subsets.** The previous examples were (very particular) convex subsets of the phase space (half-spaces, discs, Euclidean balls). It is rather natural to try an investigation of more general convex subsets. Let *C* be a convex bounded subset of  $\mathbb{R}^{2n}$  and  $\mathbf{1}_C$  be the indicator function of *C*. We have

$$\forall u \in L^2(\mathbb{R}^n), \quad \iint_C \mathcal{W}(u,u)(x,\xi) dx d\xi \leq \lambda_+(C) \|u\|_{L^2(\mathbb{R}^n)}^2.$$

On page 2178 of his 1988 article, *Maximum signal energy concentration in a time-frequency domain* (Proc. IEEE), P. Flandrin writes

・ロン ・回 と ・ ヨン ・ ヨン

13

**Convex subsets.** The previous examples were (very particular) convex subsets of the phase space (half-spaces, discs, Euclidean balls). It is rather natural to try an investigation of more general convex subsets. Let *C* be a convex bounded subset of  $\mathbb{R}^{2n}$  and  $\mathbf{1}_C$  be the indicator function of *C*. We have

$$\forall u \in L^2(\mathbb{R}^n), \quad \iint_C \mathcal{W}(u, u)(x, \xi) dx d\xi \leq \lambda_+(C) \|u\|_{L^2(\mathbb{R}^n)}^2.$$

On page 2178 of his 1988 article, *Maximum signal energy concentration in a time-frequency domain* (Proc. IEEE), P. Flandrin writes *"it is conjectured that the inequality* 

$$\lambda_+(C) \le 1$$
 is true for any convex domain C", (9)

a quite mild commitment for the validity of (9), although that statement was referred to later on as *Flandrin's conjecture* in the literature.

If Flandrin's conjecture were true, we would have for *C* convex subset of  $\mathbb{R}^{2n}$  (not necessarily bounded or with finite measure) and for all  $u \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\iint_{C} \mathcal{W}(u,u)(x,\xi) dx d\xi \leq \|u\|_{L^{2}(\mathbb{R}^{n})}^{2}.$$
(10)

イロン イヨン イヨン イヨン

If Flandrin's conjecture were true, we would have for C convex subset of  $\mathbb{R}^{2n}$ (not necessarily bounded or with finite measure) and for all  $u \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\iint_{C} \mathcal{W}(u,u)(x,\xi) dx d\xi \leq \|u\|_{L^{2}(\mathbb{R}^{n})}^{2}.$$
(10)

Indeed, when C is convex with infinite Lebesgue measure, we find that for  $u \in \mathscr{S}(\mathbb{R}^n)$  (implying  $\mathcal{W}(u, u) \in \mathscr{S}(\mathbb{R}^{2n})$ ),

・ロト ・回ト ・ヨト ・ヨト
If Flandrin's conjecture were true, we would have for C convex subset of  $\mathbb{R}^{2n}$ (not necessarily bounded or with finite measure) and for all  $u \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\iint_{C} \mathcal{W}(u,u)(x,\xi) dx d\xi \leq \|u\|_{L^{2}(\mathbb{R}^{n})}^{2}.$$
(10)

Indeed, when C is convex with infinite Lebesgue measure, we find that for  $u \in \mathscr{S}(\mathbb{R}^n)$  (implying  $\mathcal{W}(u, u) \in \mathscr{S}(\mathbb{R}^{2n})$ ), thanks to the Lebesgue Dominated Convergence Theorem, we have

$$\iint_{C} W(u, u)(x, \xi) dx d\xi = \lim_{\lambda \to +\infty} \iint_{C \cap \{(x, \xi), \max(|x|, |\xi|) \le \lambda\}} W(u, u)(x, \xi) dx d\xi,$$

$$\uparrow$$
convex

14

If Flandrin's conjecture were true, we would have for C convex subset of  $\mathbb{R}^{2n}$ (not necessarily bounded or with finite measure) and for all  $u \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\iint_{C} \mathcal{W}(u,u)(x,\xi) dx d\xi \leq \|u\|_{L^{2}(\mathbb{R}^{n})}^{2}.$$
(10)

Indeed, when C is convex with infinite Lebesgue measure, we find that for  $u \in \mathscr{S}(\mathbb{R}^n)$  (implying  $\mathcal{W}(u, u) \in \mathscr{S}(\mathbb{R}^{2n})$ ), thanks to the Lebesgue Dominated Convergence Theorem, we have

$$\iint_{C} W(u, u)(x, \xi) dx d\xi = \lim_{\lambda \to +\infty} \iint_{C \cap \{(x, \xi), \max(|x|, |\xi|) \le \lambda\}} W(u, u)(x, \xi) dx d\xi,$$

$$\uparrow$$
convex

and assuming Flandrin's conjecture, we get  $\iint_C W(u, u)(x, \xi) dx d\xi \leq ||u||^2_{L^2(\mathbb{R}^n)}$ .

If Flandrin's conjecture were true, we would have for C convex subset of  $\mathbb{R}^{2n}$ (not necessarily bounded or with finite measure) and for all  $u \in \mathscr{S}(\mathbb{R}^n)$ ,

$$\iint_{C} \mathcal{W}(u,u)(x,\xi) dx d\xi \leq \|u\|_{L^{2}(\mathbb{R}^{n})}^{2}.$$
(10)

Indeed, when C is convex with infinite Lebesgue measure, we find that for  $u \in \mathscr{S}(\mathbb{R}^n)$  (implying  $\mathcal{W}(u, u) \in \mathscr{S}(\mathbb{R}^{2n})$ ), thanks to the Lebesgue Dominated Convergence Theorem, we have

$$\iint_{C} W(u,u)(x,\xi) dx d\xi = \lim_{\lambda \to +\infty} \iint_{C \cap \{(x,\xi), \max(|x|,|\xi|) \le \lambda\}} W(u,u)(x,\xi) dx d\xi,$$

$$\uparrow$$
convex

and assuming Flandrin's conjecture, we get  $\iint_C W(u, u)(x, \xi) dx d\xi \leq ||u||_{L^2(\mathbb{R}^n)}^2$ . Conversely, you may also apply (10) to *C* convex bounded and recover Flandrin's conjecture for *C* via the boundedness of  $\mathbf{1}_C^{WEYL}$ .

1. Introduc	tion Positive results
2. Positive results, Examples and Counterexam	ples Convex subsets
3. More results and comm	ents The quarter-plane

### The quarter-plane

We choose now to focus our attention on a simple-looking case, when C is the "quarter-plane"  $C_0 = \{(x,\xi) \in \mathbb{R}^2, x \ge 0, \xi \ge 0\}.$ 

イロト イポト イヨト イヨト 二日

1. Introduction	Positive results
2. Positive results, Examples and Counterexamples	Convex subsets
3. More results and comments	The quarter-plane

### The quarter-plane

We choose now to focus our attention on a simple-looking case, when C is the "quarter-plane"  $C_0 = \{(x,\xi) \in \mathbb{R}^2, x \ge 0, \xi \ge 0\}$ . We shall study the operator

$$A_0 = (H(x)H(\xi))^{WEYL} = (\mathbf{1}_{C_0})^{WEYL}$$

イロン イヨン イヨン イヨン

### The quarter-plane

We choose now to focus our attention on a simple-looking case, when C is the "quarter-plane"  $C_0 = \{(x,\xi) \in \mathbb{R}^2, x \ge 0, \xi \ge 0\}$ . We shall study the operator

$$A_0 = \left(H(x)H(\xi)\right)^{\text{Weyl}} = \left(\mathbf{1}_{C_0}\right)^{\text{Weyl}}$$

where  $H = \mathbf{1}_{\mathbb{R}_+}$ , that is the Weyl quantization of the indicator function of the first quarter of the plane.

#### Theorem

15

Let  $A_0$  be the operator with Weyl symbol  $H(x)H(\xi)$ , where H is the Heaviside function. Then  $A_0$  is a bounded self-adjoint operator on  $L^2(\mathbb{R})$  such that

$$\inf(\operatorname{spectrum}(A_0)) < 0 < 1 < \sup(\operatorname{spectrum}(A_0)). \tag{11}$$

・ロト ・回 ト ・ヨト ・ヨト

### The quarter-plane

We choose now to focus our attention on a simple-looking case, when C is the "quarter-plane"  $C_0 = \{(x,\xi) \in \mathbb{R}^2, x \ge 0, \xi \ge 0\}$ . We shall study the operator

$$A_0 = \left(H(x)H(\xi)\right)^{\text{Weyl}} = \left(\mathbf{1}_{C_0}\right)^{\text{Weyl}}$$

where  $H = \mathbf{1}_{\mathbb{R}_+}$ , that is the Weyl quantization of the indicator function of the first quarter of the plane.

#### Theorem

15

Let  $A_0$  be the operator with Weyl symbol  $H(x)H(\xi)$ , where H is the Heaviside function. Then  $A_0$  is a bounded self-adjoint operator on  $L^2(\mathbb{R})$  such that

$$\inf(\operatorname{spectrum}(A_0)) < 0 < 1 < \sup(\operatorname{spectrum}(A_0)). \tag{11}$$

This theorem was proven in the paper entitled *On integrals over a convex set of the Wigner distribution,* by B. Delourme, T. Duyckaerts and N.L., published by the Journal of Fourier Analysis and Applications, volume 26, February 2020.

### Corollary (A counterexample to Flandrin's conjecture)

There exists a function  $\phi_0 \in \mathscr{S}(\mathbb{R})$ , with  $L^2(\mathbb{R})$  norm equal to 1 such that

$$\iint_{x\geq 0,\xi\geq 0}\mathcal{W}(\phi_0,\phi_0)(x,\xi)dxd\xi>1.$$
(12)

(ロ) (部) (注) (注) (注)

There exists a >0 such that  $\iint_{0 \le x \le a, 0 \le \xi \le a} \mathcal{W}(\phi_0, \phi_0)(x, \xi) dx d\xi > 1.$ 

### Corollary (A counterexample to Flandrin's conjecture)

There exists a function  $\phi_0 \in \mathscr{S}(\mathbb{R})$ , with  $L^2(\mathbb{R})$  norm equal to 1 such that

$$\iint_{x\geq 0,\xi\geq 0} \mathcal{W}(\phi_0,\phi_0)(x,\xi) dx d\xi > 1.$$
(12)

There exists a > 0 such that  $\iint_{0 \le x \le a, 0 \le \xi \le a} \mathcal{W}(\phi_0, \phi_0)(x, \xi) dx d\xi > 1.$ 

As a consequence, there exists a > 0 such that

 $\lambda_+([0,a]\times[0,a])>1,$ 

invalidating Flandrin's conjecture.

Rethinking the whole business Managing the quarter-plane Final comments and questions

# 3. More results and comments

### Rethinking the whole business

We have seen a simple example where E was a half-space

 $E = \{(x,\xi) \in \mathbb{R}^n \times \mathbb{R}^n, L(x,\xi) \ge \alpha\}, \text{ where } L \text{ is a linear form, } \alpha \in \mathbb{R},$ 

(ロ) (部) (注) (注) (注)

**Rethinking the whole business** Managing the quarter-plane Final comments and questions

# 3. More results and comments

### Rethinking the whole business

We have seen a simple example where E was a half-space

 $E = \{(x,\xi) \in \mathbb{R}^n \times \mathbb{R}^n, L(x,\xi) \ge \alpha\}, \text{ where } L \text{ is a linear form, } \alpha \in \mathbb{R},$ 

In that case (assuming  $L \neq 0$ ), we may find affine symplectic coordinates  $(y, \eta)$  on  $\mathbb{R}^{2n}$  such that  $L(x, \xi) - \alpha = y_1$ ,

イロン イボン イヨン トラ

Rethinking the whole business Managing the quarter-plane Final comments and questions

# 3. More results and comments

### Rethinking the whole business

We have seen a simple example where E was a half-space

 $E = \{(x,\xi) \in \mathbb{R}^n \times \mathbb{R}^n, L(x,\xi) \ge \alpha\}, \text{ where } L \text{ is a linear form, } \alpha \in \mathbb{R},$ 

In that case (assuming  $L \neq 0$ ), we may find affine symplectic coordinates  $(y, \eta)$  on  $\mathbb{R}^{2n}$  such that  $L(x, \xi) - \alpha = y_1$ , implying with the symplectic covariance of the Weyl calculus that  $\mathbf{1}_F^{WEYL}$  is unitarily equivalent to the orthogonal projection

**Rethinking the whole business** Managing the quarter-plane Final comments and questions

# 3. More results and comments

### Rethinking the whole business

We have seen a simple example where E was a half-space

$$E = \{(x,\xi) \in \mathbb{R}^n \times \mathbb{R}^n, L(x,\xi) \ge \alpha\}, \text{ where } L \text{ is a linear form, } \alpha \in \mathbb{R},$$

In that case (assuming  $L \neq 0$ ), we may find affine symplectic coordinates  $(y, \eta)$  on  $\mathbb{R}^{2n}$  such that  $L(x, \xi) - \alpha = y_1$ , implying with the symplectic covariance of the Weyl calculus that  $\mathbf{1}_F^{WEYL}$  is unitarily equivalent to the orthogonal projection

 $u\mapsto u(y)H(y_1).$ 

・ロト ・聞ト ・ヨト ・ヨト

Rethinking the whole business Managing the quarter-plane Final comments and questions

# 3. More results and comments

### Rethinking the whole business

We have seen a simple example where E was a half-space

 $E = \{(x,\xi) \in \mathbb{R}^n \times \mathbb{R}^n, L(x,\xi) \ge \alpha\}, \text{ where } L \text{ is a linear form, } \alpha \in \mathbb{R},$ 

In that case (assuming  $L \neq 0$ ), we may find affine symplectic coordinates  $(y, \eta)$  on  $\mathbb{R}^{2n}$  such that  $L(x, \xi) - \alpha = y_1$ , implying with the symplectic covariance of the Weyl calculus that  $\mathbf{1}_{F}^{WEYL}$  is unitarily equivalent to the orthogonal projection

 $u \mapsto u(y)H(y_1).$ 

### The simplicity of that first case is misleading

In many cases, including some rather explicit ones, the Weyl quantization of the rough Hamiltonian  $\mathbf{1}_E(x,\xi)$  could be far from a projection and may have a rather complicated spectrum

イロン イヨン イヨン イヨン

In many cases, including some rather explicit ones, the Weyl quantization of the rough Hamiltonian  $\mathbf{1}_E(x,\xi)$  could be far from a projection and may have a rather complicated spectrum with a supremum which could be strictly larger than 1 and an infimum which could be negative.

イロン イヨン イヨン イヨン

In many cases, including some rather explicit ones, the Weyl quantization of the rough Hamiltonian  $\mathbf{1}_E(x,\xi)$  could be far from a projection and may have a rather complicated spectrum with a supremum which could be strictly larger than 1 and an infimum which could be negative.

In some sense, although we have the trivial identity for functions

$$\mathbf{1}_E(x,\xi)^2 = \mathbf{1}_E(x,\xi),$$

we shall see that the quantization process by the Weyl formula is destroying that property. The Wigner distribution is not a probability density: although it takes real values and its integral is 1 (for a normalized  $L^2$  function), it can take negative values so that

$$\iint \mathbf{1}_E(x,\xi) \mathcal{W}(u,u)(x,\xi) dx d\xi$$
 does not necessarily belongs to [0,1].

The terminology "quasi-probability" is awfully misleading.

18

A D A A B A A B A A B A

In many cases, including some rather explicit ones, the Weyl quantization of the rough Hamiltonian  $\mathbf{1}_E(x,\xi)$  could be far from a projection and may have a rather complicated spectrum with a supremum which could be strictly larger than 1 and an infimum which could be negative.

In some sense, although we have the trivial identity for functions

$$\mathbf{1}_E(x,\xi)^2 = \mathbf{1}_E(x,\xi),$$

we shall see that the quantization process by the Weyl formula is destroying that property. The Wigner distribution is not a probability density: although it takes real values and its integral is 1 (for a normalized  $L^2$  function), it can take negative values so that

$$\iint \mathbf{1}_{E}(x,\xi)\mathcal{W}(u,u)(x,\xi)dxd\xi$$
 does not necessarily belongs to [0,1].

The terminology "quasi-probability" is awfully misleading.

Understand integrals of the Wigner distribution on subsets of the phase space forces us to consider the Weyl quantization of the function  $\mathbf{1}_E(x,\xi)$ . The Heisenberg Uncertainty Principle shows that non-commutation properties are governing operators whose symbols actually depend on conjugate variables (say  $x_1, \xi_1$ ) and these properties are of course distorting the classical identities satisfied by classical Hamiltonians.

We must point out as well that we do not have here at our disposal a semi-classical version of our quantization which could ensure some bridge between classical properties and operator-theoretic results as it is the case for the quantization of nice smooth semi-classical symbols depending on a small parameter h.

・ロト ・聞ト ・ヨト ・ヨト

We must point out as well that we do not have here at our disposal a semi-classical version of our quantization which could ensure some bridge between classical properties and operator-theoretic results as it is the case for the quantization of nice smooth semi-classical symbols depending on a small parameter h.

In particular for a symbol *a* such that  $a(x,\xi,h) = a_1(x,h\xi)$ ,  $a_1 \in C_b^{\infty}(\mathbb{R}^{2n})$ , we have the following result: if for all  $(x,\xi,h) \in \mathbb{R}^n \times \mathbb{R}^n \times (0,1]$  we have  $a(x,\xi,h) \leq 1$ , then there exists a semi-norm *C* of the symbol *a* such that

We must point out as well that we do not have here at our disposal a semi-classical version of our quantization which could ensure some bridge between classical properties and operator-theoretic results as it is the case for the quantization of nice smooth semi-classical symbols depending on a small parameter h.

In particular for a symbol *a* such that  $a(x,\xi,h) = a_1(x,h\xi)$ ,  $a_1 \in C_b^{\infty}(\mathbb{R}^{2n})$ , we have the following result: if for all  $(x,\xi,h) \in \mathbb{R}^n \times \mathbb{R}^n \times (0,1]$  we have  $a(x,\xi,h) \leq 1$ , then there exists a semi-norm *C* of the symbol *a* such that

$$\mathsf{Id} - a^{\mathsf{W}_{\mathrm{EYL}}} + Ch^2 \ge 0 \qquad \text{i.e.} \quad a^{\mathsf{W}_{\mathrm{EYL}}} \le \mathsf{Id} + Ch^2, \tag{13}$$

A D A A B A A B A A B A

an inequality following from the Fefferman-Phong Inequality which implies as well the following lemma.

### Lemma

Let a be a semi-classical symbol of order 0, e.g.  $a(x,\xi,h) = a_1(x,h\xi)$ ,  $a_1 \in C_b^{\infty}(\mathbb{R}^{2n})$ , such that for all  $(x,\xi,h) \in \mathbb{R}^n \times \mathbb{R}^n \times (0,1]$  we have

 $0 \leq a(x,\xi,h) \leq 1.$ 

Then there exists a semi-norm C of the symbol a such that

20

$$-Ch^2 \leq a^{\mathrm{W}_{\mathrm{EYL}}} \leq \mathrm{Id} + Ch^2.$$

イロン イヨン イヨン イヨン

Managing the quarter-plane. We want an explicit spectral decomposition for the operator

$$A_0 = (H(x)H(\xi))^{W_{\text{EYL}}}.$$

The kernel of  $A_0$  is

$$k_0(x,y) = H(x+y)\hat{H}(y-x) = H(x+y)\frac{1}{2}\Big(\delta_0(y-x) + \frac{1}{i\pi}pv\frac{1}{y-x}\Big).$$

First tool: use logarithmic coordinates on each half-line: The mapping  $\boldsymbol{\Psi}$  defined by

$$\begin{array}{rcl} \Psi: \ L^2(\mathbb{R}) & \longrightarrow & L^2(\mathbb{R}; \mathbb{C}^2) \\ u & \mapsto & \left(\phi_1(t) = u(e^t)e^{t/2}, \phi_2(t) = u(-e^t)e^{t/2}\right) \end{array}$$

is an isometric isomorphism of Hilbert spaces.

21

イロン イヨン イヨン イヨン

Let us use rather formally the following identities for an operator K with kernel  $k(\boldsymbol{x},\boldsymbol{y})$ : we get

Let us use rather formally the following identities for an operator K with kernel  $k(\boldsymbol{x},\boldsymbol{y})$ : we get

$$\langle KHu, Hv \rangle_{L^2(\mathbb{R})} = \iint k(x, y)H(y)u(y)H(x)\overline{v}(x)dydx$$

Let us use rather formally the following identities for an operator K with kernel k(x, y): we get

$$\langle KHu, Hv \rangle_{L^{2}(\mathbb{R})} = \iint k(x, y)H(y)u(y)H(x)\overline{v}(x)dydx$$
$$= \iint k(e^{s}, e^{t})u(e^{t})\overline{v}(e^{s})e^{s+t}dsdt$$

Let us use rather formally the following identities for an operator K with kernel k(x, y): we get

$$\langle KHu, Hv \rangle_{L^{2}(\mathbb{R})} = \iint k(x, y)H(y)u(y)H(x)\overline{v}(x)dydx$$

$$= \iint k(e^{s}, e^{t})u(e^{t})\overline{v}(e^{s})e^{s+t}dsdt$$

$$= \iint \underbrace{k(e^{s}, e^{t})e^{\frac{s+t}{2}}}_{\overline{k}(s,t)} \underbrace{u(e^{t})e^{\frac{t}{2}}\overline{v}(e^{s})e^{\frac{s}{2}}}_{\phi_{1}(t)} \underbrace{\overline{v}(e^{s})e^{\frac{s}{2}}}_{\overline{\psi}_{1}(s)} dsdt,$$

22

Let us use rather formally the following identities for an operator K with kernel k(x, y): we get

$$\langle KHu, Hv \rangle_{L^{2}(\mathbb{R})} = \iint k(x, y)H(y)u(y)H(x)\overline{v}(x)dydx$$

$$= \iint k(e^{s}, e^{t})u(e^{t})\overline{v}(e^{s})e^{s+t}dsdt$$

$$= \iint \underbrace{k(e^{s}, e^{t})e^{\frac{s+t}{2}}}_{\overline{k}(s,t)} \underbrace{u(e^{t})e^{\frac{t}{2}}\overline{v}(e^{s})e^{\frac{s}{2}}}_{\phi_{1}(t)} \underbrace{dsdt}_{\overline{\psi_{1}(s)}}$$

so that the new kernel for the operator HKH in logarithmic coordinates is  $\tilde{k}(s,t).$ 

22

1. Introduction	Rethinking the whole business
2. Positive results, Examples and Counterexamples	Managing the quarter-plane
3. More results and comments	Final comments and questions

We check first the "diagonal" terms

(日) (四) (三) (三) (三)

We check first the "diagonal" terms

$$H(x)H(y)H(x+y)\frac{1}{2}\left(\delta_0(y-x)+\frac{1}{i\pi}\mathsf{pv}\frac{1}{y-x}\right)$$

We check first the "diagonal" terms

23

$$H(x)H(y)H(x+y)\frac{1}{2}\left(\delta_{0}(y-x) + \frac{1}{i\pi}pv\frac{1}{y-x}\right)$$
  
=  $H(x)H(y)\frac{1}{2}\delta_{0}(y-x) + \frac{1}{2i\pi}pv\frac{H(x)H(y)}{y-x}$ 

We check first the "diagonal" terms

23

$$\begin{split} H(x)H(y)H(x+y)\frac{1}{2}\Big(\delta_0(y-x)+\frac{1}{i\pi}\mathsf{pv}\frac{1}{y-x}\Big) \\ &= H(x)H(y)\frac{1}{2}\delta_0(y-x)+\frac{1}{2i\pi}\mathsf{pv}\frac{H(x)H(y)}{y-x} \\ &= \frac{1}{2}\delta_0(t-s)+\frac{1}{2i\pi}\mathsf{pv}\frac{e^{t/2}e^{s/2}}{e^t-e^s} \end{split}$$

We check first the "diagonal" terms

23

$$\begin{split} H(x)H(y)H(x+y)\frac{1}{2}\Big(\delta_0(y-x)+\frac{1}{i\pi}\mathsf{pv}\frac{1}{y-x}\Big)\\ &=H(x)H(y)\frac{1}{2}\delta_0(y-x)+\frac{1}{2i\pi}\mathsf{pv}\frac{H(x)H(y)}{y-x}\\ &=\frac{1}{2}\delta_0(t-s)+\frac{1}{2i\pi}\mathsf{pv}\frac{e^{t/2}e^{s/2}}{e^t-e^s}\\ &=\frac{1}{2}\delta_0(t-s)+\frac{1}{2i\pi}\mathsf{pv}\frac{1}{e^{\frac{t-s}{2}}-e^{\frac{s-t}{2}}} \end{split}$$
 (this is a

this is a convolution in logarithmic coordinates)

We check first the "diagonal" terms

23

$$\begin{split} H(x)H(y)H(x+y)\frac{1}{2}\Big(\delta_0(y-x) + \frac{1}{i\pi}\mathsf{pv}\frac{1}{y-x}\Big) \\ &= H(x)H(y)\frac{1}{2}\delta_0(y-x) + \frac{1}{2i\pi}\mathsf{pv}\frac{H(x)H(y)}{y-x} \\ &= \frac{1}{2}\delta_0(t-s) + \frac{1}{2i\pi}\mathsf{pv}\frac{e^{t/2}e^{s/2}}{e^t - e^s} \\ &= \frac{1}{2}\delta_0(t-s) + \frac{1}{2i\pi}\mathsf{pv}\frac{1}{e^{\frac{t-s}{2}} - e^{\frac{s-t}{2}}} \quad \text{(this is a convolution in logarithmic coordinates)} \\ &= \frac{1}{2}\delta_0(t-s) + \frac{1}{2i\pi}\mathsf{pv}\frac{1}{2\sinh(\frac{t-s}{2})} \quad \text{(pretty explicit stuff)} \end{split}$$

イロン イヨン イヨン イヨン

The off-diagonal terms, now: we get with  $\check{H}(t) = H(-t)$ ,

The off-diagonal terms, now: we get with  $\check{H}(t) = H(-t)$ ,

$$\left(H(x)\check{H}(y)+\check{H}(x)H(y)\right)H(x+y)rac{1}{2}\left(\delta_{0}(y-x)+rac{1}{i\pi}\mathsf{pv}rac{1}{y-x}
ight)$$
$$(H(x)\check{H}(y) + \check{H}(x)H(y))H(x+y)\frac{1}{2}(\delta_0(y-x) + \frac{1}{i\pi}\mathsf{pv}\frac{1}{y-x})$$
$$= \frac{1}{2i\pi}\mathsf{pv}\frac{H(x)\check{H}(y)H(x+y)}{y-x} + \frac{1}{2i\pi}\mathsf{pv}\frac{H(y)\check{H}(x)H(x+y)}{y-x}$$

24

$$\begin{split} & (H(x)\check{H}(y) + \check{H}(x)H(y))H(x+y)\frac{1}{2}\Big(\delta_0(y-x) + \frac{1}{i\pi}\mathsf{pv}\frac{1}{y-x}\Big) \\ &= \frac{1}{2i\pi}\mathsf{pv}\frac{H(x)\check{H}(y)H(x+y)}{y-x} + \frac{1}{2i\pi}\mathsf{pv}\frac{H(y)\check{H}(x)H(x+y)}{y-x} \\ &= \frac{1}{2i\pi}\mathsf{pv}\frac{e^{t/2}e^{s/2}H(s-t)}{-e^t-e^s} + \frac{1}{2i\pi}\mathsf{pv}\frac{e^{t/2}e^{s/2}H(t-s)}{e^t+e^s} \end{split}$$

24

$$\begin{split} H(x)\check{H}(y) + \check{H}(x)H(y) \Big) H(x+y) &\frac{1}{2} \Big( \delta_0(y-x) + \frac{1}{i\pi} \mathsf{pv} \frac{1}{y-x} \Big) \\ &= \frac{1}{2i\pi} \mathsf{pv} \frac{H(x)\check{H}(y)H(x+y)}{y-x} + \frac{1}{2i\pi} \mathsf{pv} \frac{H(y)\check{H}(x)H(x+y)}{y-x} \\ &= \frac{1}{2i\pi} \mathsf{pv} \frac{e^{t/2}e^{s/2}H(s-t)}{-e^t-e^s} + \frac{1}{2i\pi} \mathsf{pv} \frac{e^{t/2}e^{s/2}H(t-s)}{e^t+e^s} \\ &= \frac{1}{2i\pi} \mathsf{pv} \frac{e^{\frac{t+s}{2}}\operatorname{sign}(t-s)}{e^t+e^s} = \frac{1}{2i\pi} \mathsf{pv} \frac{\operatorname{sign}(t-s)}{2\operatorname{cosh}(\frac{t-s}{2})}, \end{split}$$

$$\begin{split} & (H(x)\check{H}(y) + \check{H}(x)H(y))H(x+y)\frac{1}{2}\Big(\delta_0(y-x) + \frac{1}{i\pi}\mathsf{pv}\frac{1}{y-x}\Big) \\ &= \frac{1}{2i\pi}\mathsf{pv}\frac{H(x)\check{H}(y)H(x+y)}{y-x} + \frac{1}{2i\pi}\mathsf{pv}\frac{H(y)\check{H}(x)H(x+y)}{y-x} \\ &= \frac{1}{2i\pi}\mathsf{pv}\frac{e^{t/2}e^{s/2}H(s-t)}{-e^t-e^s} + \frac{1}{2i\pi}\mathsf{pv}\frac{e^{t/2}e^{s/2}H(t-s)}{e^t+e^s} \\ &= \frac{1}{2i\pi}\mathsf{pv}\frac{e^{\frac{t+2}{2}}\operatorname{sign}(t-s)}{e^t+e^s} = \frac{1}{2i\pi}\mathsf{pv}\frac{\operatorname{sign}(t-s)}{2\operatorname{cosh}(\frac{t-s}{2})}, \end{split}$$

which is another convolution.

24

Next step: study the explicit  $2 \times 2$  matrix multiplier. We work indeed on  $L^2(\mathbb{R}_t; \mathbb{C}^2)$  one space dimension (the *t* variable) but acting on vectors of  $\mathbb{C}^2$ . We have

$$\mathcal{N}( au) = egin{pmatrix} \mathsf{a}_{11}( au) & \mathsf{a}_{12}( au) \ \overline{\mathsf{a}_{12}( au)} & \mathsf{0} \end{pmatrix}$$

イロン イヨン イヨン イヨン

э

Next step: study the explicit 2 × 2 matrix multiplier. We work indeed on  $L^2(\mathbb{R}_t; \mathbb{C}^2)$  one space dimension (the *t* variable) but acting on vectors of  $\mathbb{C}^2$ . We have

$$\mathcal{N}( au) = egin{pmatrix} \mathsf{a}_{11}( au) & \mathsf{a}_{12}( au) \ \overline{\mathsf{a}_{12}( au)} & 0 \end{pmatrix}$$

The eigenvalues  $\lambda_{-} \leq \lambda_{+}$  of  $\mathcal{N}(\tau)$  are such that

25

Next step: study the explicit 2 × 2 matrix multiplier. We work indeed on  $L^2(\mathbb{R}_t; \mathbb{C}^2)$  one space dimension (the *t* variable) but acting on vectors of  $\mathbb{C}^2$ . We have

$$\mathcal{N}( au) = egin{pmatrix} \mathsf{a}_{11}( au) & \mathsf{a}_{12}( au) \ \overline{\mathsf{a}_{12}( au)} & 0 \end{pmatrix}$$

The eigenvalues  $\lambda_{-} \leq \lambda_{+}$  of  $\mathcal{N}(\tau)$  are such that

$$\lambda_{-}(\tau) < 0 < 1 < \lambda_{+}(\tau), \tag{14}$$

Next step: study the explicit 2 × 2 matrix multiplier. We work indeed on  $L^2(\mathbb{R}_t; \mathbb{C}^2)$  one space dimension (the *t* variable) but acting on vectors of  $\mathbb{C}^2$ . We have

$$\mathcal{N}( au) = egin{pmatrix} \mathsf{a}_{11}( au) & \mathsf{a}_{12}( au) \ \overline{\mathsf{a}_{12}( au)} & \mathsf{0} \end{pmatrix}$$

The eigenvalues  $\lambda_{-} \leq \lambda_{+}$  of  $\mathcal{N}(\tau)$  are such that

25

$$\lambda_{-}(\tau) < 0 < 1 < \lambda_{+}(\tau), \tag{14}$$

・ロン ・回 と ・ ヨ と ・ ヨ と

3

if and only if

$$a_{12}(\tau) \neq 0$$
 and  $|a_{12}(\tau)|^2 > 1 - a_{11}(\tau)$ . (15)

Next step: study the explicit 2 × 2 matrix multiplier. We work indeed on  $L^2(\mathbb{R}_t; \mathbb{C}^2)$  one space dimension (the *t* variable) but acting on vectors of  $\mathbb{C}^2$ . We have

$$\mathcal{N}( au) = egin{pmatrix} \mathsf{a}_{11}( au) & \mathsf{a}_{12}( au) \ \overline{\mathsf{a}_{12}( au)} & \mathsf{0} \end{pmatrix}$$

The eigenvalues  $\lambda_{-} \leq \lambda_{+}$  of  $\mathcal{N}(\tau)$  are such that

25

$$\lambda_{-}(\tau) < 0 < 1 < \lambda_{+}(\tau), \tag{14}$$

・ロト ・聞ト ・ヨト ・ヨト

3

if and only if

$$a_{12}(\tau) \neq 0$$
 and  $|a_{12}(\tau)|^2 > 1 - a_{11}(\tau).$  (15)

It is possible to translate the sought spectral property in terms of singularities of functions: the function  $g_0$  (involved in the symbol) defined by

Next step: study the explicit 2 × 2 matrix multiplier. We work indeed on  $L^2(\mathbb{R}_t; \mathbb{C}^2)$  one space dimension (the *t* variable) but acting on vectors of  $\mathbb{C}^2$ . We have

$$\mathcal{N}( au) = egin{pmatrix} \mathsf{a}_{11}( au) & \mathsf{a}_{12}( au) \ \overline{\mathsf{a}_{12}( au)} & \mathsf{0} \end{pmatrix}$$

The eigenvalues  $\lambda_{-} \leq \lambda_{+}$  of  $\mathcal{N}(\tau)$  are such that

25

$$\lambda_{-}(\tau) < 0 < 1 < \lambda_{+}(\tau), \tag{14}$$

if and only if

$$a_{12}(\tau) \neq 0$$
 and  $|a_{12}(\tau)|^2 > 1 - a_{11}(\tau)$ . (15)

It is possible to translate the sought spectral property in terms of singularities of functions: the function  $g_0$  (involved in the symbol) defined by

#### $g_0(t) = H(t) \operatorname{sech} t$

has a singularity at t = 0 so that its Fourier transform  $a_{12}(\tau)$  cannot go to 0 rapidly when  $\tau \to +\infty$ . On the other hand  $1 - a_{11}(\tau)$  belongs to the Schwartz space and decays rapidly when  $\tau \to +\infty$ .

The study of  $\mathbf{1}_{E}^{WEYL}$  for a subset *E* of the phase space is highly correlated to some particular set of special functions related to *E*:

・ロト ・回ト ・ヨト ・ヨト

26

The study of  $\mathbf{1}_{E}^{WEYL}$  for a subset *E* of the phase space is highly correlated to some particular set of special functions related to *E*:

• Hermite functions and Laguerre polynomials for ellipses,

The study of  $\mathbf{1}_{E}^{WEYL}$  for a subset *E* of the phase space is highly correlated to some particular set of special functions related to *E*:

- Hermite functions and Laguerre polynomials for ellipses,
- Airy functions for parabolas,

26

The study of  $\mathbf{1}_{E}^{WEYL}$  for a subset *E* of the phase space is highly correlated to some particular set of special functions related to *E*:

- Hermite functions and Laguerre polynomials for ellipses,
- Airy functions for parabolas,
- Homogeneous distributions for hyperbolas

The study of  $\mathbf{1}_{E}^{WEYL}$  for a subset *E* of the phase space is highly correlated to some particular set of special functions related to *E*:

- Hermite functions and Laguerre polynomials for ellipses,
- Airy functions for parabolas,

26

• Homogeneous distributions for hyperbolas and so on.

The study of  $\mathbf{1}_{E}^{WEYL}$  for a subset *E* of the phase space is highly correlated to some particular set of special functions related to *E*:

- Hermite functions and Laguerre polynomials for ellipses,
- Airy functions for parabolas,

26

• Homogeneous distributions for hyperbolas and so on.

It is quite likely that the "shape" of *E* will determine the type of special functions to be studied to getting a diagonalization of the operator  $\mathbf{1}_E^{\text{WEYL}}$ .

1. Introduction	Rethinking the whole business
2. Positive results, Examples and Counterexamples	Managing the quarter-plane
3. More results and comments	Final comments and questions

A couple of questions:

・ロン ・回 と ・ ヨ と ・ ヨ と

æ

A couple of questions:

27

• For a general convex polygon  $P_N$  with N vertices, it is possible to prove that

 $\mathbf{1}_{P_N}^{\mathrm{Weyl}} \leq \sigma_N,$ 

where  $\sigma_N$  does not depend on the area of the polygon but only on N. Is it true that  $\sup_N \sigma_N < +\infty$ ?

・ロン ・回と ・ヨン・

A couple of questions:

27

• For a general convex polygon  $P_N$  with N vertices, it is possible to prove that

 $\mathbf{1}_{P_N}^{\mathrm{Weyl}} \leq \sigma_N,$ 

where  $\sigma_N$  does not depend on the area of the polygon but only on N. Is it true that  $\sup_N \sigma_N < +\infty$ ?

• Does there exist  $\sigma > 1$  such that for all convex compact subsets K of the plane  $\mathbf{1}_{K}^{\text{WEYL}} \leq \sigma$ ?

• In 2n dimensions, defining

$$E(a_1,\ldots,a_n)=\{(x,\xi)\in\mathbb{R}^n\times\mathbb{R}^n, 2\pi\sum_{1\leq j\leq n}\frac{x_j^2+\xi_j^2}{a_j}\leq 1\}.$$

Find

$$\sup \operatorname{SPECTRUM}(\mathbf{1}_{E(a_1,\ldots,a_n)})^{\operatorname{WEYL}}$$

Done by E. LIEB & Y. OSTROVER for  $a_1 = \cdots = a_n$ , but in 2n dimensions with  $n \ge 2$ , ellipsoids are not symplectically equivalent to the Euclidean ball.

・ロン ・回と ・ヨン・

#### Maximum signal energy concentration in a time-frequency domain

P. Flandrin,

Proc. IEEE Int. Conf. Acoustics 4 (1988), no. 1.

イロト イポト イヨト イヨト 二日

Maximum signal energy concentration in a time-frequency domain

P. Flandrin,

Proc. IEEE Int. Conf. Acoustics 4 (1988), no. 1.

Bounds on integrals of the Wigner function: the hyperbolic case J.G. Wood, A.J. Bracken, Journal of Mathematical Physics, 46, 4, (2005).

Maximum signal energy concentration in a time-frequency domain P. Flandrin,

Proc. IEEE Int. Conf. Acoustics 4 (1988), no. 1.

29

Bounds on integrals of the Wigner function: the hyperbolic case J.G. Wood, A.J. Bracken, Journal of Mathematical Physics, 46, 4, (2005).

On Integrals over a Convex Set of the Wigner Distribution B. Delourme, T. Duyckaerts, N.L., Journal of Fourier Analysis and Applications, 26, (2020), 1.

イロン イヨン イヨン イヨン

Maximum signal energy concentration in a time-frequency domain P. Flandrin, Proc. IEEE Int. Conf. Acoustics 4 (1988), no. 1.

Bounds on integrals of the Wigner function: the hyperbolic case J.G. Wood, A.J. Bracken, Journal of Mathematical Physics, 46, 4, (2005).

On Integrals over a Convex Set of the Wigner Distribution B. Delourme, T. Duyckaerts, N.L., Journal of Fourier Analysis and Applications, 26, (2020), 1.

29

Integrating the Wigner Distribution on subsets of the phase space, a Survey
N.L.,
https://arxiv.org/abs/2102.08090

Maximum signal energy concentration in a time-frequency domain P. Flandrin, Proc. IEEE Int. Conf. Acoustics 4 (1988), no. 1.

Bounds on integrals of the Wigner function: the hyperbolic case J.G. Wood, A.J. Bracken, Journal of Mathematical Physics, 46, 4, (2005).

On Integrals over a Convex Set of the Wigner Distribution B. Delourme, T. Duyckaerts, N.L., Journal of Fourier Analysis and Applications, 26, (2020), 1.

Integrating the Wigner Distribution on subsets of the phase space, a Survey  ${\tt N.L.}$  ,

https://arxiv.org/abs/2102.08090

## Thank you for your attention

A D A A B A A B A A B A

Maximum signal energy concentration in a time-frequency domain P. Flandrin, Proc. IEEE Int. Conf. Acoustics 4 (1988), no. 1.

Bounds on integrals of the Wigner function: the hyperbolic case J.G. Wood, A.J. Bracken, Journal of Mathematical Physics, 46, 4, (2005).

On Integrals over a Convex Set of the Wigner Distribution B. Delourme, T. Duyckaerts, N.L., Journal of Fourier Analysis and Applications, 26, (2020), 1.

Integrating the Wigner Distribution on subsets of the phase space, a Survey  ${\tt N.L.}$  ,

https://arxiv.org/abs/2102.08090

# Thank you for your attention

# Best wishes to Jorge