

New examples of extremal domains for the first eigenvalue of the Laplace-Beltrami operator in a Riemannian manifold with boundary

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Abstract. We build new examples of extremal domains with small prescribed volume for the first eigenvalue of the Laplace-Beltrami operator in some Riemannian manifolds with boundary. These domains are close to half balls of small radius centered at a nondegenerate critical point of the mean curvature function of the boundary of the manifold, and their boundary intersects the boundary of the manifold orthogonally.

1 Introduction

New examples of domains with small prescribed volume that are critical points for the first eigenvalue of the Dirichlet Laplace-Beltrami operator are built in [28], under the hypothesis that the Riemannian manifold has at least one nondegenerate critical point of the scalar curvature function. In that case, such domains are given by small perturbations of geodesic balls of small radius centered at a nondegenerate critical point of the scalar curvature. This result has been generalized in [12] to all compact Riemannian manifolds by eliminating the hypothesis of the existence of a nondegenerate critical point of the scalar curvature.

Such examples of critical points for the Laplace-Beltrami operator are parallels to similar shape examples of critical points for the area functional, under the same assumptions, which lead to the construction of constant mean curvature small topological spheres, see [29, 35].

The aim of this paper is to give some new examples of domains Ω that are critical points for the first eigenvalue of the Laplace-Beltrami operator (i.e. extremal domains) in some Riemannian manifolds M with boundary. Such examples are new because the boundary of the domain is partially included in the boundary of the manifold. The domains we obtain are close to half-balls centered at a point of ∂M where the mean curvature of ∂M is critical and the criticality is not degenerate. In particular, in the simplest situation, M can be a domain of the Euclidean space, see Fig. 1. Again, we can make a parallel with the case of the area, for which a similar result has been proven in [17, 18].

Assume that we are given (M, g) an $(n + 1)$ -dimensional Riemannian manifold, $n \geq 1$, with boundary $\partial M \neq \emptyset$. The boundary ∂M is a smooth n -dimensional Riemannian manifold with the metric \tilde{g} induced by g . For a domain Ω contained in the interior of M , $\Omega \subset \overset{\circ}{M}$, the first eigenvalue of the Laplace-Beltrami operator with 0 Dirichlet boundary condition is then given by

$$\lambda_\Omega = \min_{u \in H_0^1(\Omega)} \frac{\int_\Omega |\nabla u|^2}{\int_\Omega u^2}.$$

If Ω is a boundary domain (i.e. a domain such that $\partial\Omega \cap \partial M \neq \emptyset$), we consider the first eigenvalue of the Laplace-

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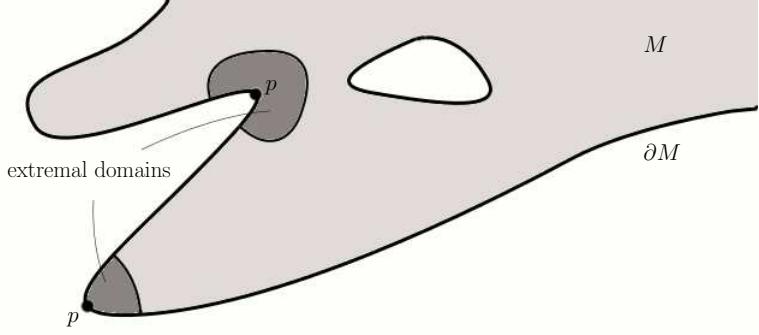


Figure 1: M can be a Euclidean domain (bounded or not). If p is a nondegenerate critical point for the mean curvature of ∂M , then it is possible to construct an extremal domain as a perturbation of a half-ball centered at p .

Beltrami operator given by

$$\lambda_\Omega = \min_{u \in \tilde{H}_0^1(\Omega)} \frac{\int_\Omega |\nabla u|^2}{\int_\Omega u^2} \quad (1)$$

where $\tilde{H}_0^1(\Omega)$ denotes the closure of the space

$$\{\varphi \in C^\infty(\Omega), \text{ Supp}(\varphi) \subset \Omega \cup \partial M\}$$

for the H^1 -norm. It is very classical that the optimization problem (1) admits a nonnegative solution if Ω has finite volume, and if Ω is connected such a solution is unique among nonnegative functions whose L^2 -norm is 1. This function is then called the first eigenfunction of Ω .

Under smoothness assumption (for example if Ω is a piecewise $C^{1,\alpha}$ -domain, see Section 2.1 for more detailed definitions, the space $\tilde{H}_0^1(\Omega)$ is equal to the space of functions in $H^1(\Omega)$ with 0 Dirichlet condition on $\partial\Omega \cap \mathring{M}$, and the function u solving (1) satisfies:

$$\begin{cases} \Delta_g u + \lambda_\Omega u = 0 & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \cap \mathring{M}, \\ g(\nabla u, \nu) = 0 & \text{on } \partial\Omega \cap \partial M \end{cases} \quad (2)$$

where ν denotes the outward normal vector to ∂M , which is well-defined as soon as Ω is included in a small enough ball, which will be the case in the whole paper. This will be referred to as a mixed eigenvalue problem over Ω . Moreover, it is also well-known that if there exists (u, λ) a nontrivial solution of (2) for a connected domain Ω such that u is nonnegative, then $\lambda = \lambda_\Omega$ is the first eigenvalue of Ω , and u is the first eigenfunction of Ω , up to a multiplicative constant.

Let us consider a boundary domain $\Omega_0 \subset M$. Ω_0 is said to be extremal if $\Omega \mapsto \lambda_\Omega$ is critical at Ω_0 with respect to variations of the domain Ω_0 which preserve its volume. In order to make this notion precise, we first introduce the definition of a *deformation* of Ω_0 .

Definition 1.1. We say that $(\Omega_t)_{t \in (-t_0, t_0)}$ is a deformation of Ω_0 , if there exists a vector field V on M , of class C^2 , such that its flow ξ , defined for $t \in (-t_0, t_0)$ by

$$\frac{d\xi}{dt}(t, p) = V(\xi(t, p)) \quad \text{and} \quad \xi(0, p) = p,$$

preserves the boundary of the manifold, i.e. $\xi(t, p) \in \partial M$ for all $(t, p) \in (-t_0, t_0) \times \partial M$, and for which

$$\Omega_t = \xi(t, \Omega_0).$$

The deformation is said to be volume preserving if the volume of Ω_t does not depend on t .

If $(\Omega_t)_{t \in (-t_0, t_0)}$ is a deformation of Ω_0 , we denote by λ_t the first eigenvalue of the Laplace-Beltrami operator $-\Delta_g$ on Ω_t . We prove in Section 2 that $t \mapsto \lambda_t$ is smooth in a neighborhood of $t = 0$. If $\Omega \subset \mathring{M}$ this fact is standard and follows from the implicit function theorem together with the fact that the first eigenvalue of the Laplace-Beltrami operator is simple, see for example [21]. When the boundary ∂M is invariant by the flow of the deformation, as required in Definition 1.1, a similar strategy still works when $\partial\Omega \cap \partial M \neq \emptyset$, but this is less classical since one needs to manage the singularities of the boundary domains under consideration, see Proposition 2.5. The derivative at 0 of $t \mapsto \lambda_t$ is then called the shape derivative of $\Omega \mapsto \lambda_\Omega$ at Ω_0 in the direction V .

This remark allows us to give the definition of an extremal domain.

Definition 1.2. A domain Ω_0 is an extremal domain for the first eigenvalue of $-\Delta_g$ if for any volume preserving deformation $\{\Omega_t\}_t$ of Ω_0 , we have

$$\frac{d\lambda_t}{dt}|_{t=0} = 0, \quad (3)$$

where $\lambda_t = \lambda_{\Omega_t}$ as defined in (1).

All along the paper, we will use a special system of coordinates, that we remind here: let $p \in \partial M$, and let N be the unit normal vector field on ∂M near p that points into M . We fix local geodesic normal coordinates $x = (x^1, \dots, x^n)$ in a neighborhood of $0 \in \mathbb{R}^n$ to parametrize U_p a neighborhood of p in ∂M by Φ . We consider the mapping

$$\Psi(x^0, x) = \text{Exp}_{\Phi(x)}(x^0 N(\Phi(x))) \quad (4)$$

which is a local diffeomorphism from a neighborhood of $0 \in \overline{\mathbb{R}_+^{n+1}}$ (where $\mathbb{R}_+^{n+1} = \{(x^0, x) \in \mathbb{R}^{n+1} : x^0 > 0\}$) into V_p a neighborhood of p in M . For all $\varepsilon > 0$ small enough, we denote $B_\varepsilon^+ \subset \mathbb{R}_+^{n+1}$ the half-ball given by the Euclidean ball of radius ε centered at the origin and restricted to $x^0 > 0$, and we denote $B_{g,\varepsilon}^+(p) = \Psi(B_\varepsilon^+) \subset \mathring{M}$.

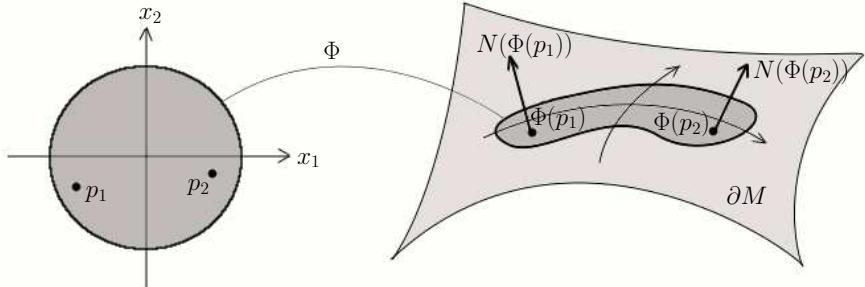


Figure 2: Our coordinates are defined as (x^0, x) , x being the normal geodesic coordinates on ∂M and x^0 the coordinate associated to the normal direction.

Now we can state the main result of the paper:

Theorem 1.3. Assume that $p_0 \in \partial M$ is a nondegenerate critical point of H , the mean curvature function of $(\partial M, \tilde{g})$. Then, for all $\varepsilon > 0$ small enough, say $\varepsilon \in (0, \varepsilon_0)$, there exists a boundary domain $\Omega_\varepsilon \subset M$ such that :

- (i) The volume of Ω_ε is equal to the Euclidean volume of B_ε^+ .
- (ii) The domain Ω_ε is extremal in the sense of Definition 1.2.
- (iii) The boundary $\overline{\partial\Omega_\varepsilon \cap \mathring{M}}$ intersects ∂M orthogonally,
- (iv) The boundary $\partial\Omega_\varepsilon \cap \mathring{M}$ is analytic if M is analytic.

Moreover, there exists $c > 0$ and, for all $\varepsilon \in (0, \varepsilon_0)$, there exists $p_\varepsilon \in \partial M$ such that $\overline{\partial\Omega_\varepsilon \cap \mathring{M}}$ is a normal graph over $\overline{\partial B_{g,\varepsilon}^+(p_\varepsilon) \cap \mathring{M}}$ for some function w_ε with

$$\|w_\varepsilon\|_{C^{2,\alpha}(\overline{\partial B_{g,\varepsilon}^+(p_\varepsilon) \cap \mathring{M}})} \leq c\varepsilon^3. \quad \text{and} \quad \text{dist}(p_\varepsilon, p_0) \leq c\varepsilon.$$

This result will be proven in Section 4.5. The strategy of the proof of this result is inspired by [28]. In order to give the outline of the paper, we recall here the strategy of the proof and insist on the main differences with [28]. The first step is to characterize the extremality of a domain Ω_0 with the Euler-Lagrange equation, that leads to:

$$g(\nabla u, \nu) = \text{constant} \quad \text{on } \partial\Omega_0 \cap \mathring{M}. \quad (5)$$

The difficulty here is to prove this characterization for domains that are only piecewise smooth (see Section 2.1 where we introduce the notion of boundary edge domain and analyze the regularity theory of mixed boundary value problem in such domains). In particular, we prove in Section 2.2 that in order to be extremal it is enough for a domain to satisfy (3) for deformations that preserve the contact angle on ∂M ; this important fact will be used in the rest of the paper for the construction of extremal domains. This is an interesting difference with the case of critical points of the area functional, as we explain in Section 2.3: condition (5) contains already the information that the contact angle between $\partial\Omega_0 \cap \mathring{M}$ and $\partial\Omega_0 \cap \partial M$ is constant and equal to $\pi/2$, see Corollary 2.6; this is due to the non-locality of the Euler-Lagrange equation for this problem. It also implies the analytic regularity of $\partial\Omega_0 \cap \mathring{M}$.

Then, thanks to a dilation of the metric and a control of the volume constraint, we reformulate in Section 3 the problem into solving for any small ε the equation

$$F(p, \varepsilon, \bar{v}) = 0 \quad (6)$$

where $p \in \partial M$, $\bar{v} \in C^{2,\alpha}(S_+^n)$ is a function that parametrize a perturbation of the half-geodesic ball $B_{g,\varepsilon}^+(p)$, and $F(p, \varepsilon, \bar{v})$ represents the difference between $g(\nabla u, \nu)$ and its mean value on the boundary of this perturbed half geodesic ball. We then want to solve this equation for $\varepsilon > 0$ by using the implicit function Theorem and therefore study the operator $\partial_{\bar{v}} F(p, 0, 0)$, which is basically related to the second order shape derivative of λ_1 at the Euclidian half-ball. This is the purpose of Sections 4.1 and 4.2, where we use a symmetrization argument to come down to the study of the same operator in the Euclidian ball, which has been done in [28]. As expected, that operator has a nontrivial kernel (because of the invariance of λ_1 by translation along $\partial\mathbb{R}_+^{n+1}$ in the Euclidian setting) and we are only able to solve

$$F(p, \varepsilon, \bar{v}(p, \varepsilon)) = k(p, \varepsilon)$$

where $k(p, \varepsilon)$ is a linear function induced by an element of $\partial\mathbb{R}_+^{n+1}$, see Proposition 4.4. Here comes the final step of the proof of Theorem 1.3, which takes into account the geometry of ∂M : by studying the expansion of F, \bar{v} with respect to ε , we prove in the end of Section 4 that close to a point p_0 which is a nondegenerate critical point of the mean curvature of ∂M , one can chose p_ε such that $k(p_\varepsilon, \varepsilon) = 0$ and conclude the proof. We insist on the fact that this step is more involved here than in [28]: indeed, the expansions in ε contain lower order term than in the case without boundary (see Lemma 4.3 and Propositions 4.4, 4.5). Nevertheless, thanks to the choice of our coordinates, the strategy still applies because these lower order terms are orthogonal to linear functions induced by elements of $\partial\mathbb{R}_+^{n+1}$.

2 Characterization of boundary extremal domains

In this section, we focus on an analytic characterization of extremal domains. The main difficulty here is to handle the shape derivative of $\Omega \mapsto \lambda_\Omega$ in a nonsmooth setting. Indeed, because of the presence of a boundary in M , we are naturally led to deal with domains that are only piecewise smooth. First, we will treat the regularity for the mixed problem (2) in some domains called boundary edge domains. We compute then the shape derivative of $\Omega \mapsto \lambda_\Omega$ in this setting. Since we have to deal with possibly nonsmooth eigenfunctions, one needs to carefully prove the differentiability of $\Omega \mapsto \lambda_\Omega$ and compute the shape derivative. We will also insist on some important aspects of the non-locality of the extremality condition for λ_1 , and compare it with the case of critical points for the area functional.

2.1 Boundary edge domains and regularity of the eigenfunction

Definition 2.1. Let Ω be a boundary domain of the manifold M , that is to say $\partial\Omega \cap \partial M \neq \emptyset$. We say that Ω is a boundary edge domain if it satisfies the following condition:

1. $\overline{\partial\Omega \cap \overset{\circ}{M}}$ and $\partial\Omega \cap \partial M$ are smooth n -dimensional submanifolds with boundary,
2. $\Gamma := \overline{\partial\Omega \cap \overset{\circ}{M}} \cap \partial M$ is a $(n-1)$ -dimensional smooth submanifold without boundary.

In that case, given $p \in \Gamma$ we can define $\omega(p)$ the angle between the normal vector to Γ tangent to ∂M and the normal vector to Γ tangent to $\partial\Omega \cap \overset{\circ}{M}$. The function $\omega : \Gamma \rightarrow [0, \pi]$ will be referred to as the contact angle of the domain Ω , see Fig. 3.

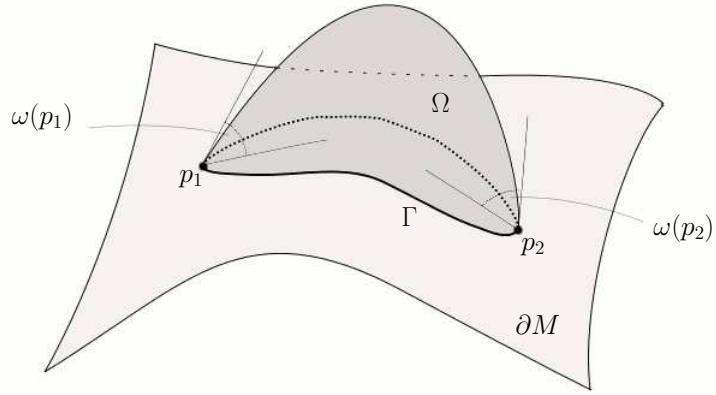


Figure 3: A boundary edge domain in M

Proposition 2.2. Let Ω be a connected boundary edge domain of finite volume such that the contact angle ω is strictly between 0 and π . Then there exists $\varepsilon > 0$ such that for any $f \in H^{-1/2+\varepsilon}(\Omega)$, the solution u of

$$\begin{cases} -\Delta_g u = f & \text{in } \Omega \\ u = 0 & \text{on } \partial\Omega \cap \overset{\circ}{M}, \\ g(\nabla u, \nu) = 0 & \text{on } \partial\Omega \cap \partial M \end{cases} \quad (7)$$

is in the space $H^{3/2+\varepsilon}(\Omega)$.

Remark 2.3. It is important for our purpose to work here with Sobolev regularity: if indeed we work with Hölder-regularity, we can only conclude that $u \in C^{0,1/2+\varepsilon}(\overline{\Omega})$, which does not suffice to justify the expression of the shape derivative, which uses the trace of the gradient on $\partial\Omega$, see Section 2.2, while from the fact that $u \in H^{3/2+\varepsilon}(\Omega)$, we can deduce that ∇u has a trace in $L^2(\partial\Omega)$ (we use here a trace theorem, valid since under our assumptions, Ω has a Lipschitz boundary).

Proof. Let $f \in H^s(\Omega)$ where $s \in (-1, 0)$. It is well-known from the variational formulation of the problem that there exists a unique $u \in H^1(\Omega)$ weak solution of (7). We wonder for which s we can state that $u \in H^{s+2}(\Omega)$. To that end, we work locally around a point $p \in \Gamma$: there exist special cylindrical coordinates (r, θ, y) such that Γ corresponds to $r = 0$, $y \in \Gamma$ parametrizes the edge (p corresponding to $y = 0$), and Ω corresponds to $0 < \theta < \omega(y)$; since Ω is a boundary edge domain, these coordinates are well-defined and C^∞ . From the literature on edge asymptotics, we know that u can be written around p as the sum of a singular function u_{sing} and a remainder term u_{reg} which is

more regular than u_{sing} ; more precisely, it is known (see for example [7, 8, 9, 13, 14]) that

$$\begin{aligned} \text{if } \omega(y) \in (0, \pi/2) &\quad \text{then } u_{sing}(r, \theta, y) = 0 \text{ and } u_{reg} \in H^{s+2}(\Omega) \\ \text{if } \omega(y) \in (\pi/2, \pi) &\quad \text{then } u_{sing}(r, \theta, y) = c(y) r^{\pi/2\omega(y)} \varphi(\theta, y), \\ \text{if } \omega(y) = \pi/2 \text{ in a neighborhood of } y = 0, &\quad \text{then } u_{sing}(r, \theta, y) = r \left(\sum_{q \geq 1} c_q(y) \ln^q(r) \varphi_q(y, \theta) \right), \end{aligned}$$

where $c, (c_q)_{q \in \mathbb{N}}$ (containing only a finite number of non-zero terms) and $\varphi, (\varphi_q)_{q \in \mathbb{N}}$ are smooth functions (we notice that when $n = 1$, the set Γ is made of two points, in that case the regularity on Γ is an empty condition). Let us conclude in the last two cases. In the second one, we know that

$$\text{if } \alpha > s' - \frac{n+1}{2}, \text{ then } r \mapsto r^\alpha \in H^{s'}(\mathbb{R}^{n+1}),$$

and therefore the regularity increases with small angles, and the worst regularity is obtained when the angle is close to π , but is always strictly better than $H^{3/2}$ which is the limit case when $\omega = \pi$ and $n+1=2$. In the last case, it is clear that $r \ln^q(r) = o(r^{1-\delta})$ for any small δ , so we obtain that the regularity is also better than $H^{3/2}$, therefore there exists s strictly above $-1/2$ such that $u \in H^{s+2}(\Omega)$.

It remains to understand the case where $\omega(0) = \pi/2$ but ω is not constant in a neighborhood of $y = 0$. In that case, the asymptotic development is more involved (phenomenon of crossing singularities), but it is explained in [8, 9] that up to an arbitrary small loss of regularity, we obtain the same range of regularity as in the case $\omega = \pi/2$, and therefore again u_{sing} is in $H^{3/2+\varepsilon}(\Omega)$. \square

In the previous proof, we have seen that the regularity is more or less monotone with respect to the contact angle: smaller is the angle, higher is the regularity, and for angles close to π , the regularity decreases up to the space $H^{3/2}$. However, it is also known that there exists some exceptional angles, for which the regularity is higher than expected (see for example [1] for a description of this phenomenon for the angle $\pi/4$ in dimension 2). We prove here that the angle $\pi/2$ is such an exceptional angle in our situation. More precisely we prove that when the angle is $\pi/2$ everywhere on the interface, the regularity is actually $C^{2,\alpha}$, whereas it was expected to be $C^{0,\alpha}$ for every α in the proof of the previous statement. This will be very useful in the proof of Theorem 1.3. This result is related to the fact that one can use a symmetrization argument to conclude that the first expected term in the asymptotic development of u vanishes.

Proposition 2.4. *Let Ω be a boundary edge domain, such that the angle ω defined on Γ is constant and equal to $\pi/2$. Then for every $\alpha < 1$ and any $f \in C^{0,\alpha}(\overline{\Omega})$, the solution u of (7) is in $C^{2,\alpha}(\overline{\Omega})$.*

Proof. We use the same setting as in the proof of Proposition 2.2, but now in the class of Hölder spaces, so we consider $f \in C^{0,\alpha}(\overline{\Omega})$. Around $p \in \Gamma$, from [8, 9, 13, 14], we know that the exponents in the asymptotic development for the mixed boundary problem are $(\pi/2\omega + k\pi/\omega)_{k \in \mathbb{N}}$, so for the angle $\pi/2$ the first terms are 1 and 3 and since $r \mapsto r^3 \ln^q(r)$ belongs to the space $C^{2,\alpha}(\overline{\Omega})$ for every α and any integer q , we conclude that

$$u(r, \theta, y) = r \left(\sum_{q \geq 1} c_q(y) \ln^q(r) \varphi_q(y, \theta) \right) + u_{reg}(r, \theta, y), \quad (8)$$

for y close to 0, r small, $\theta \in (0, \pi/2)$ and where functions (c_q, φ_q) are smooth and u_{reg} is in $C^{2,\alpha}$ locally around p .

The result will be proven if we prove that $c_q = 0$ for $q \geq 1$. To that end, we use a symmetrization procedure through ∂M , using around $p \in \Gamma$ the coordinates (x^0, x) described in (4). We define $U = \Psi^{-1}(\Omega \cap B_{g,r_0}^+(p)) \subset B_{r_0}^+$, so that $\partial U \cap (\{0\} \times \mathbb{R}^n) = \Psi^{-1}(\partial \Omega \cap \partial M \cap \overline{B_{g,r_0}^+(p)})$. With this choice of coordinates, U is again a boundary edge domain whose contact angle is constant and equal to $\pi/2$ on $\gamma = \Psi^{-1}(\Gamma)$.

We now define $W = \{(x^0, x) / (|x^0|, x) \in U\}$ and

$$\forall (x^0, x) \in W, \quad \dot{u}(x^0, x) = \begin{cases} u(x^0, x) & \text{if } x^0 > 0 \\ u(-x^0, x) & \text{if } x^0 < 0 \end{cases} \quad \text{and similarly we define } \dot{g} \text{ and } \dot{f}.$$

Since the contact angle is $\pi/2$, the symmetrized domain W is smooth around 0; using that u satisfies a Neumann boundary condition on $\partial\Omega \cap \partial M$, we deduce that \mathring{u} satisfies

$$\begin{cases} -\Delta_{\mathring{g}} \mathring{u} = \mathring{f} & \text{in } W \\ \mathring{u} = 0 & \text{on } \partial W \cap B_{r_0}. \end{cases}$$

and finally, the symmetrized metric \mathring{g} is no longer C^∞ but has Lipschitz coefficients, and \mathring{f} is again in $C^{0,\alpha}(\overline{W})$. Since the Laplace operator can be written in a divergence form

$$\Delta_{\mathring{g}} u = \frac{1}{\sqrt{|\mathring{g}|}} \partial_i \left(\sqrt{|\mathring{g}|} \mathring{g}^{ij} \partial_j \mathring{u} \right)$$

we can apply the regularity theory for elliptic PDE in divergence form in a smooth set, with Lipschitz coefficients: precisely, from [22, Theorem 8.34] we know that $\mathring{u} \in C^{1,\alpha}(\overline{W})$ and therefore $(c_q)_{q \geq 1}$ must be zero, and finally $u \in C^{2,\alpha}(\overline{\Omega})$. \square

2.2 Shape derivative in nonsmooth domains

Proposition 2.5. *Let Ω_0 be a connected boundary domain of finite volume. Assume that $(\Omega_t)_t$ is a deformation of Ω_0 induced by the vector field V , as defined in Definition 1.2. Then $t \mapsto \lambda_t$ is C^∞ around $t = 0$. If moreover Ω_0 is a boundary edge domain such that the contact angle is strictly between 0 and π , then $g(\nabla u_0, \nu_0) \in L^2(\partial\Omega_0)$ and*

$$\frac{d\lambda_t}{dt} |_{t=0} = - \int_{\partial\Omega_0 \cap \mathring{M}} (g(\nabla u_0, \nu_0))^2 g(V, \nu_0) \, d\text{vol}_g, \quad (9)$$

where $d\text{vol}_g$ is the volume element on $\partial\Omega_0 \cap \mathring{M}$ for the metric induced by g and ν_0 is the normal vector field on $\partial\Omega_0 \cap \mathring{M}$.

Before proving this result, we give some remarks and consequences. The differentiability of some similar shape functional for mixed boundary value problem is studied in [32, Section 3.9] in the case of a smooth domain, which corresponds to the case of a angle constant and equal to π . In that case formula (9) is not valid since the eigenfunction u is not smooth enough. Also in [3], the case of angles different from π is considered, but for a different shape functional, and restricted to the two-dimensional case.

Proposition 2.5 allows us to characterize extremal domains for the first eigenvalue of the Laplace-Beltrami operator under 0 mixed boundary conditions, and state the problem of finding extremal domains into the solvability of an over-determined elliptic problem. As a consequence of the previous result, we obtain indeed:

Corollary 2.6. *Let Ω_0 be a boundary edge domain. Then Ω_0 is extremal if and only if the first eigenfunction u_0 of Ω_0 satisfies*

$$g(\nabla u_0, \nu_0) = \text{constant on } \partial\Omega_0 \cap \mathring{M} \quad (10)$$

where ν_0 is the outward normal vector field on $\partial\Omega_0 \cap \mathring{M}$. In that case, $\partial\Omega_0 \cap \mathring{M}$ necessarily meets ∂M orthogonally, that is to say the contact angle function ω is equal to $\pi/2$ on Γ .

Proof of Corollary 2.6: Let Ω_0 be a boundary extremal domain for the first eigenvalue of the Laplace-Beltrami operator, with 0 Dirichlet boundary condition on $\partial\Omega_0 \cap \mathring{M}$ and 0 Neumann boundary condition on $\partial\Omega_0 \cap \partial M$. Using Proposition 2.5, we obtain

$$\int_{\partial\Omega_0 \cap \mathring{M}} (g(\nabla u_0, \nu_0))^2 g(V, \nu_0) \, d\text{vol}_g = 0$$

for all field V preserving the volume of the domain, i.e. such that

$$\int_{\partial\Omega_0 \cap \mathring{M}} g(V, \nu_0) \, d\text{vol}_g = 0. \quad (11)$$

This means that $g(\nabla u_0, \nu_0)$ is constant. On the other hand, if $g(\nabla u_0, \nu_0)$ is constant, by the previous proposition we have that Ω_0 is extremal, because V satisfy (11).

It remains to investigate the angle between $\partial\Omega_0 \cap \overset{\circ}{M}$ and $\partial\Omega_0 \cap \partial M$, when (10) is satisfied. Let's assume that $y \mapsto \omega(y)$ is not constantly equal to $\pi/2$; then there exists a neighborhood in $\mathcal{Y} \subset \Gamma = \partial\Omega \cap \overset{\circ}{M} \cap \partial M$ where ω is different from $\pi/2$. We work locally around a point $y_0 \in \mathcal{Y}$. We need now a more explicit version of the asymptotic development written in the proof of Proposition 2.2. To that end, we use the results of [10, 11, 9] which asserts that since the principal part of our operator is the Euclidian Laplacian, we have, up to a smooth change of coordinates, that $u_0(r, \theta, y)$ can be written $u_{reg}(r, \theta, y) + u_{sing}(r, \theta, y)$ with:

if $\omega(y) \in (0, \pi/2)$ in \mathcal{Y} , then $u_{sing} = 0$

and $u_{reg} \in H^{s+2}(\Omega)$ is flat at order 2, which means $u_{reg} = \mathcal{O}(r^2)$ and $\nabla u_{reg} = \mathcal{O}(r)$,

if $\omega(y) \in (\pi/2, \pi)$ in \mathcal{Y} , then $u_{sing}(r, \theta, y) = c(y)r^{\pi/2\omega(y)} \cos\left(\frac{\pi}{2\omega(y)}\theta\right)$,

and u_{reg} is more flat than u_{sing} , meaning $u_{reg} = o(r)$ and $\nabla u_{reg} = o(1)$,

(note that here, with the terminology of [8, 9], there is no crossing singularities, since $\omega(y) \neq \pi/2$ on \mathcal{Y} and we are only interested in the first term of the asymptotic). Therefore in the first case $g(\nabla u_0, \nu_0) = \mathcal{O}(r)$ and in the second case $g(\nabla u_0, \nu_0)$ behaves like $-\frac{\pi}{2\omega(y)}c(y)r^{\pi/2\omega(y)-1} \sin\left(\frac{\pi}{2\omega(y)}\theta\right)$, and therefore, in both cases, cannot be a nonzero constant on $\partial\Omega \cap \overset{\circ}{M} = \{\theta = \omega(y)\}$. This is a contradiction (remind that from maximum principle, the constant $g(\nabla u_0, \nu_0)$ cannot be a zero), and one concludes that $\omega(y) = \pi/2$ everywhere on Γ . \square

Proof of Proposition 2.5: Let Ω_0 be a boundary domain, connected and of finite volume. We denote by $\xi_t = \xi(t, \cdot)$ the flow associated to V , ν_t the outward unit normal vector field to $\partial\Omega_t$. We first remind that, since Ω_t is connected, for t small enough λ_t the first eigenvalue of Ω_t with mixed boundary condition is simple, so one can define $t \mapsto u_t \in \tilde{H}_0^1(\Omega_t)$ the one-parameter family of first eigenfunctions of the Laplace-Beltrami operator, normalized to be positive and have $L^2(\Omega_t)$ -norm equal to 1. As usual in the computation of a shape derivative, we consider $\hat{u}_t = u_t \circ \xi(t, \cdot) \in \tilde{H}_0^1(\Omega_0)$.

Step 1: $\exists t_0 > 0$ such that $t \in (-t_0, t_0) \mapsto (\hat{u}_t, \lambda_t) \in \tilde{H}_0^1(\Omega_0) \times \mathbb{R}$ is C^∞ .

The variational formulation of the equation satisfied by u_t is:

$$\int_{\Omega_t} g(\nabla u_t, \nabla \varphi) = \lambda_t \int_{\Omega_t} u_t \varphi, \quad \forall \varphi \in \tilde{H}_0^1(\Omega_t).$$

We are going to transport that formulation on the fixed domain Ω_0 , in order to obtain the variational formulation satisfied by $\hat{u}_t \in \tilde{H}_0^1(\Omega_0)$. To that aim, we use the following equality, which relies on the fact that

$$\xi_t(\partial\Omega_0 \cap \partial M) = \partial\Omega_t \cap \partial M$$

and is a consequence of the hypothesis $\xi_t(\partial M) \subset \partial M$:

$$\tilde{H}_0^1(\Omega_0) = \{\varphi \circ \xi_t, \varphi \in \tilde{H}_0^1(\Omega_t)\}.$$

With this equality and a change of variable (see for example [21] for details), we obtain:

$$\int_{\Omega_0} g(A(t) \nabla \hat{u}_t, \nabla \varphi) = \lambda_t \int_{\Omega_0} \hat{u}_t \varphi J_t, \quad \forall \varphi \in \tilde{H}_0^1(\Omega_0),$$

where

$$J_t = \det(D\xi_t), \quad \text{and } A(t) := J_t D\xi_t^{-1} (D\xi_t^{-1})^T.$$

We then define

$$\begin{aligned} G : (-t_0, t_0) \times \tilde{H}_0^1(\Omega_0) \times \mathbb{R} &\longrightarrow \tilde{H}_0^1(\Omega_0)' \times \mathbb{R} \\ (t, v, \mu) &\longmapsto \left(-\operatorname{div}_g(A(t) \nabla v) - \mu v J_t, \int_{\Omega_0} v^2 J_t - 1 \right) \end{aligned}$$

where $\tilde{H}_0^1(\Omega_0)'$ is the dual space of $\tilde{H}_0^1(\Omega_0)$, and $-\operatorname{div}_g(A(t)\nabla v)$ has to be understood in the weak sense:

$$\langle -\operatorname{div}_g(A(t)\nabla v), \varphi \rangle_{\tilde{H}_0^1(\Omega_0)'\times\tilde{H}_0^1(\Omega_0)} = \int_{\Omega_0} g(A(t)\nabla v, \nabla \varphi).$$

It is easy to check that G is C^∞ , see again [21] for more details. In order to apply the implicit function theorem for the equation $G(t, \hat{u}_t, \lambda_t) = 0$, we focus on the differential of G at $(0, u_0, \lambda_0)$ with respect to the couple (v, μ) :

$$\partial_{(v,\mu)} G(0, u_0, \lambda_0)(w, \nu) = \left(-\Delta_g w - \nu u_0 - \lambda_0 w, 2 \int_{\Omega_0} u_0 w \right), \quad \forall (w, \nu) \in \tilde{H}_0^1(\Omega_t) \times \mathbb{R}.$$

Because of the Banach isomorphism Theorem, in order to prove that such differential is an isomorphism, it is enough to prove that given $(f, \Lambda) \in \tilde{H}_0^1(\Omega_0)' \times \mathbb{R}$, the equation

$$\left(-\Delta_g w - \nu u_0 - \lambda_0 w, 2 \int_{\Omega_0} u_0 w \right) = (f, \Lambda)$$

admits a unique solution $(w, \nu) \in \tilde{H}_0^1(\Omega_0) \times \mathbb{R}$. The operator $-\Delta_g - \lambda_0 \mathbb{1}$ has a one-dimensional kernel, spanned by u_0 . Therefore $f + \nu u_0$ is in the range of $-\Delta_g - \lambda_0 \mathbb{1}$ if and only if it is orthogonal to u_0 (in the sense of the duality $\tilde{H}_0^1(\Omega_0)' \times \tilde{H}_0^1(\Omega_0)$). This leads to the unique value $\nu = -\langle f, u_0 \rangle$.

Moreover, one knows that the solutions w of $(-\Delta_g - \lambda_0 \mathbb{1})w = f + \nu u_0$ form a one-dimensional affine space $v_0 + \text{Span}(u_0)$, so $w = v_0 + \alpha u_0$ for some $\alpha \in \mathbb{R}$. The equation $2 \int_{\Omega_0} u_0 w = \Lambda$ uniquely determines α and so w . We can conclude that $\partial_{(v,\mu)} F(0, u_0, \lambda_0)$ is an isomorphism, and therefore $t \mapsto (\hat{u}_t, \lambda_t)$ is C^∞ .

Now and for the rest of the proof, Ω_0 is assumed to be a boundary edge domain whose contact angle is always strictly between 0 and π .

Step 2: Generalized Green formula: we prove in this step that given $\varepsilon \in (0, 1/2)$ and Ω a Lipschitz domain, denoting $H^s(\Delta_g, \Omega) := \{\varphi \in H^s(\Omega), \Delta_g \varphi \in L^2(\Omega)\}$ for $s \in (1/2, 3/2)$ we have:

$$\begin{aligned} \forall u \in H^{3/2-\varepsilon}(\Delta_g, \Omega), \forall v \in H^{1/2+\varepsilon}(\Delta_g, \Omega), \\ \int_{\Omega} (v \Delta_g u - u \Delta_g v) = \langle g(\nabla u, \nu_0), v \rangle_{H^{-\varepsilon}(\partial\Omega) \times H^{\varepsilon}(\partial\Omega)} - \langle u, g(\nabla v, \nu_0) \rangle_{H^{1-\varepsilon}(\partial\Omega) \times H^{-1+\varepsilon}(\partial\Omega)} \end{aligned} \quad (12)$$

When u, v are smooth, this equality is just the classical Green formula. The above generalization is easily obtained by a density argument, using the following result from [6, Lemma 2 and 3]:

$$\begin{aligned} H^{3/2-\varepsilon}(\Delta_g, \Omega) &= \{\varphi \in H^1(\Omega), \Delta_g \varphi \in L^2(\Omega) \text{ and } \varphi|_{\partial\Omega} \in H^{1-\varepsilon}(\Omega)\}, \\ \text{and } H^{1/2+\varepsilon}(\Delta_g, \Omega) &= \{\varphi \in H^{\varepsilon}(\Omega), \Delta_g \varphi \in L^2(\Omega) \text{ and } g(\nabla \varphi, \nu_0)|_{\partial\Omega} \in H^{-1+\varepsilon}(\Omega)\} \end{aligned} \quad (13)$$

and that $C^\infty(\overline{\Omega})$ is dense in $H^{3/2-\varepsilon}(\Delta_g, \Omega)$.

Step 3: Computation of $\frac{d}{dt}u_t$: From $u_t = \hat{u}_t \circ \xi_t^{-1}$, we obtain that $u' = \frac{d}{dt}|_{t=0} u_t$ is well-defined in Ω_0 and that

$$u' = \hat{u}' - g(\nabla u, V), \quad (14)$$

where $\hat{u}' = \frac{d}{dt}|_{t=0} \hat{u}_t \in \tilde{H}_0^1(\Omega_0)$, well-defined from Step 1. Using that $u \in H^{3/2+\varepsilon}(\Omega_0)$ and that $\hat{u}' \in H^1(\Omega_0)$, we know from (14) that $u' \in H^{1/2+\varepsilon}(\Omega_0)$. We also know that, the domain Ω_0 being piecewise C^∞ , the functions u and u' are locally C^∞ on $\overline{\Omega_0} \setminus \Gamma$. With these regularities, we can compute the equation and the boundary conditions satisfied by u' : first, we differentiate with respect to t the identity

$$\Delta_g u_t + \lambda_t u_t = 0. \quad (15)$$

and evaluate the result at $t = 0$ to obtain

$$\Delta_g u'_0 + \lambda_0 u'_0 = -\lambda'_0 u_0, \text{ in } \Omega_0. \quad (16)$$

Moreover, using again (14), we obtain that

$$u' = -g(\nabla u, V) \text{ on } \partial\Omega \cap \mathring{M}.$$

and since $u_0 = 0$ on $\partial\Omega_0 \cap \mathring{M}$, only the normal component of V plays a rôle in the previous formula. Therefore, we have, again since $\xi(t, \partial\Omega_0 \cap \partial M) = \partial\Omega_t \cap \partial M$:

$$u' = -g(\nabla u_0, \nu_0) g(V, \nu_0), \quad \text{on } \partial\Omega_0 \cap \mathring{M} \quad (17)$$

About the Neumann part of the boundary, we have:

$$\text{for all } p \in \partial\Omega_0 \cap \partial M, \quad g(\nabla u_t(\xi(t, p)), \nu_t) = 0.$$

Since V is tangential on ∂M , using the normal geodesic coordinates we have $\nu_t = -\partial_{x^0}$ on $\partial\Omega_t \cap \partial M$, and in particular it does not depend on t and

$$g(\nabla u_t(\xi(t, p)), \nu_t) = -\partial_{x^0} u_t(\xi(t, p)) = 0. \quad (18)$$

So, differentiating (18) with respect to t and evaluating the result at $t = 0$ we obtain

$$0 = -\partial_{x^0} \partial_t u_0 - g(\nabla \partial_{x^0} u_0, V) = -\partial_{x^0} \partial_t u_0 = g(\nabla \partial_t u_0, \nu_0) \quad (19)$$

on $\partial\Omega_0 \cap \partial M$, where we used the facts that $\partial_{x^0} u_0 = 0$ on $\partial\Omega_0 \cap \partial M$ and that $g(V, \nu_0) = 0$ in $\partial\Omega_0 \cap \partial M$.

Step 5: Computation of $\frac{d}{dt}|_{t=0} \lambda_t$: From (16), multiplying by u and integrating over Ω , we obtain, using the generalized Green formula together with the regularity we have proven on u and u' :

$$\lambda'_0 = \int_{\Omega} (-\Delta_g u' - \lambda_0 u') u = \int_{\Omega} (-\Delta_g u - \lambda u) u' + \langle u', g(\nabla u, \nu_0) \rangle_{H^{-\varepsilon}(\partial\Omega) \times H^{\varepsilon}(\partial\Omega)} - \langle u, g(\nabla u', \nu_0) \rangle_{H^{1-\varepsilon}(\partial\Omega) \times H^{-1+\varepsilon}(\partial\Omega)}.$$

Since $u = 0$ on $\partial\Omega \cap \mathring{M}$ and $g(\nabla u', \nu_0) = 0$ on $\partial\Omega \cap \partial M$, we have $\langle u, g(\nabla u', \nu_0) \rangle_{H^{1-\varepsilon}(\partial\Omega) \times H^{-1+\varepsilon}(\partial\Omega)} = 0$. Finally, since u and u' are smooth enough so that $(g(\nabla u, \nu_0)|_{\partial\Omega}, u'|_{\partial\Omega}) \in L^2(\partial\Omega)$, we can write

$$\langle u', g(\nabla u, \nu_0) \rangle_{H^{-\varepsilon}(\partial\Omega) \times H^{\varepsilon}(\partial\Omega)} = \int_{\partial\Omega} u' g(\nabla u, \nu_0) = - \int_{\partial\Omega \cap \mathring{M}} (g(\nabla u, \nu_0))^2 g(V, \nu),$$

and we finally obtain

$$\lambda' = - \int_{\partial\Omega \cap \mathring{M}} (g(\nabla u, \nu_0))^2 g(V, \nu).$$

□

2.3 Extremal domains versus the isoperimetric problem

As we said, extremal domains are the critical points of the functional

$$\Omega \rightarrow \lambda_{\Omega}$$

under a volume constraint $\text{Vol}_g \Omega = \kappa$. The problem of finding extremal domains for the first eigenvalue of the Laplace-Beltrami operator is considered, by the mathematical community, very close to the isoperimetric problem.

Given a compact Riemannian manifold M and a positive number $\kappa < \text{Vol}_g(M)$, where $\text{Vol}_g(M)$ denotes the volume of the manifold M , the isoperimetric problem consists in studying, among the compact hypersurfaces $\Sigma \subset M$ enclosing a region Ω of volume κ , those which minimize the area functional

$$\Omega \rightarrow \text{Vol}_g(\partial\Omega \cap \mathring{M})$$

(note that we do not take in account the area of $\partial\Omega$ coming from the boundary of M). The solutions of the isoperimetric problem are (where they are smooth enough) constant mean curvature hypersurfaces and intersect

the boundary of the manifold orthogonally (see for example [31]). In fact, constant mean curvature hypersurfaces intersecting ∂M orthogonally are exactly the critical points of the area functional

$$\Omega \rightarrow \text{Vol}_g(\partial\Omega \cap \mathring{M})$$

under a volume constraint $\text{Vol}_g \Omega = \kappa$.

In the case of a manifold M without boundary, it is well known that the determination of the isoperimetric profile

$$I_\kappa := \inf_{\Omega \subset M : \text{Vol}_g \Omega = \kappa} \text{Vol}_g \partial\Omega$$

is related to the Faber-Krahn profile, where one looks for the least value of the first eigenvalue of the Laplace-Beltrami operator amongst domains with prescribed volume

$$FK_\kappa := \inf_{\Omega \subset M : \text{Vol}_g \Omega = \kappa} \lambda_\Omega$$

(see [4]). For this reason it is natural to expect that the solutions to the isoperimetric problem for small volumes are close in some sense to the solutions of the Faber-Krahn minimization problem. And such closeness can be expected also for the corresponding critical points.

The results known up to now about extremal domains underline such expectations. In the case of a manifold without boundary, the constructions of extremal domains in [28, 12] are the parallel of the constructions of constant mean curvature topological spheres in a Riemannian manifold M done in [35, 29]. And in the case of a manifold with boundary, our construction is the parallel of the constructions of constant mean curvature topological half-spheres in a Riemannian manifold M done in [17, 18].

For small isoperimetric solutions and small constant mean curvature topological spheres the following results are also known:

- small solutions of the isoperimetric problem are close to geodesic spheres of small radius centered at a point where the scalar curvature function on (M, g) is maximal, [26, 15, 27];
- small solutions of the isoperimetric problem in a domain are close to geodesic topological half-spheres of small radius centered at a point where the mean curvature of the boundary of the domain is maximal, [18];
- in dimension 3 simply connected constant mean curvature surfaces with small diameter necessarily concentrate at a critical point of the scalar curvature, [25];
- in dimension 3 simply connected constant mean curvature surfaces with bounded area and free boundary in a domain necessarily concentrate at a critical point of the mean curvature of the boundary of the domain, [24].

More precisely, in [15] the author provide an asymptotic expansion of the isoperimetric profile for small volume; the parallel for the Faber-Krahn profile has been proved in [16]. Therefore it is expected that Faber-Krahn minimizers of small volume exist and are also close to geodesic balls of small radius centered at a point where the scalar curvature function on (M, g) is maximal, though as far as we know, such results are open. It is natural to expect that the parallel results to [18, 24, 25] for extremal domains also hold.

Nevertheless, Proposition 2.5 and Corollary 2.6 show a very interesting difference between extremal domains and critical points of the area functional, based on the following:

Remark 2.7. A significant fact contained in the statement of Proposition 2.5 is that the shape derivative for the first eigenvalue of the Laplace-Beltrami operator with mixed boundary condition in the boundary edge domain Ω_0 does not contain a singular term supported by the “corner part” of the boundary $\partial\Omega_0$, as it is the case for the area functional, see (21).

In order to understand the consequences of this remark, let's compare the Euler-Lagrange equations of the two problems: criticality for λ_1 is written

$$\frac{d\lambda_t}{dt}|_{t=0} = \int_{\partial\Omega_0 \cap \mathring{M}} (g(\nabla u_0, \nu_0))^2 g(V, \nu_0) \, d\text{vol}_g = 0 \quad (20)$$

whereas for the area functional we have

$$\frac{d}{dt} \text{Vol}_g(\partial\Omega_t \cap \mathring{M})|_{t=0} = \int_{\partial\Omega_0 \cap \mathring{M}} H_0 g(V, \nu_0) + \int_{\Gamma} g(V, \tau_0) = 0, \quad (21)$$

where $(\Omega_t)_t$ is a volume preserving deformation of Ω_0 given by the vector field V , H_0 is the mean curvature of $\partial\Omega_0 \cap \mathring{M}$, ν_0 is the normal vector on $\partial\Omega_0 \cap \mathring{M}$, and τ_0 is the normal vector to Γ tangent to $\partial\Omega_0 \cap \mathring{M}$. For the area functional, the consequence of (21) is that in order to be critical Ω_0 must satisfy, denoting ν_1 the normal vector to Γ tangent to ∂M :

$$H_0 \equiv \text{constant}, \quad \left[\text{and } g(\tau_0, \nu_1) = 0 \text{ or equivalently } \omega \equiv \pi/2 \text{ on } \Gamma \right],$$

the first condition being obtained with vector fields V supported in \mathring{M} whereas the second condition is obtained thanks to vector fields V that are supported in a neighborhood of Γ . For λ_1 , only using vector fields V that are supported in \mathring{M} we obtain as a consequence of (20) that in order to be critical Ω_0 must satisfy:

$$g(\nabla u_0, \nu_0) = \text{constant} \quad \text{on } \partial\Omega_0 \cap \mathring{M}. \quad (22)$$

The fact that the contact angle is $\pi/2$ on Γ is already contained in the above equation (see Corollary 2.6), and therefore domains that are critical domains for λ_1 in the sense of Definition 1.2 (i.e. for any vector field V tangent on ∂M) are the same as critical domains for λ_1 restricted to vector fields supported in \mathring{M} , which is not the case for the area functional.

In other words, one can easily build surfaces that have a constant mean curvature but intersects the boundary ∂M with an angle different from $\pi/2$ (and therefore are not extremal sets for the relative perimeter under volume constraint), whereas every set satisfying (22) intersects the boundary ∂M with angle equal to $\pi/2$.

These properties lie on the fact that the operator given by the mean curvature is local while the Dirichlet to Neumann operator is nonlocal.

3 Analysis of the problem

3.1 Notations and formulation of the problem

Euclidean notations. We define the following notations:

$$\mathbb{R}_+^{n+1} = \{x = (x^0, x') = (x^0, x^1, \dots, x^n) \in \mathbb{R}^{n+1} : x^0 > 0\}$$

will be the upper Euclidean half-space,

$$B_1^+ = B_1 \cap \mathbb{R}_+^{n+1}$$

will be the upper Euclidean unit half-ball and

$$S_+^n = \{x \in S^n : x^0 > 0\}$$

will be the upper Euclidean unit hemisphere. Given a continuous function $f : \overline{S_+^n} \mapsto (0, \infty)$, we also denote

$$B_f^+ := \{x \in \mathbb{R}_+^{n+1} : 0 < |x| < f(x/|x|)\}.$$

Riemannian notations in (M, g) . Let p a point of ∂M . We denote by E_1, \dots, E_n the orthonormal base of $T_p \partial M$ associated to the geodesic normal coordinates x^1, \dots, x^n in ∂M around p . If the point $q \in \partial M$ has coordinates $x' \in \mathbb{R}^n$, we set

$$\Theta(x') := \sum_{i=1}^n x^i E_i \in T_p \partial M. \quad (23)$$

The point $q \in \partial M$ whose geodesic coordinates are given by x' is

$$q = \Phi(x') = \text{Exp}_p^{\partial M}(\Theta(x')).$$

Given a continuous function $f : \overline{S_+^n} \mapsto (0, \infty)$ whose L^∞ norm is small (say less than the cut locus of p) we define

$$B_{g,f}^+(p) := \left\{ \text{Exp}_{\Phi(x')}^M(x^0 N(\Phi(x'))) \quad : \quad x \in \mathbb{R}_+^{n+1} \quad 0 < |x| < f(x/|x|) \right\}.$$

The subscript g is meant to remind the reader that this definition depends on the metric.

Formulation of the problem. Our aim is to show that, for all $\varepsilon > 0$ small enough, we can find a point $p_\varepsilon \in \partial M$ and a (smooth) function $v = v(p_\varepsilon, \varepsilon) : \overline{S_+^n} \rightarrow \mathbb{R}$ with 0 Neumann condition at the boundary of S_+^n such that

$$\text{Vol } B_{g,\varepsilon(1+v)}^+(p) = \varepsilon^n \text{Vol } B_1^+ \tag{24}$$

and the over-determined elliptic problem

$$\begin{cases} \Delta_g \phi + \lambda \phi = 0 & \text{in } B_{g,\varepsilon(1+v)}^+(p) \\ \phi = 0 & \text{on } \partial B_{g,\varepsilon(1+v)}^+(p) \cap \mathring{M} \\ g(\nabla \phi, \nu) = 0 & \text{on } \partial B_{g,\varepsilon(1+v)}^+(p) \cap \partial M \\ g(\nabla \phi, \nu) = \text{constant} & \text{on } \partial B_{g,\varepsilon(1+v)}^+(p) \cap \mathring{M} \end{cases} \tag{25}$$

has a nontrivial positive solution, where ν is the normal vector on $\partial B_{g,\varepsilon(1+v)}^+(p)$. Notice that the 0 Neumann boundary condition on v is justified by Corollary 2.6. Indeed, the half ball $B_{g,\varepsilon}^+(p)$ intersects ∂M orthogonally, and then, since an extremal domain also intersects ∂M orthogonally, the deformation v should satisfy a fortiori a 0 Neumann boundary condition.

3.2 Dilation of the metric

We follow the strategy of [28], paying attention to the fact that we are working in a more general situation because our domains are boundary edge domains. Our first aim is to give a sense to the problem when $\varepsilon = 0$. Observe that, considering the dilated metric $\bar{g} := \varepsilon^{-2} g$, Problem (24)-(25) is equivalent to finding a point $p \in \partial M$ and a function $v : \overline{S_+^n} \rightarrow \mathbb{R}$ with 0 Neumann condition at the boundary of S_+^n such that

$$\text{Vol } B_{\bar{g},1+v}^+(p) = \text{Vol } B_1^+ \tag{26}$$

and for which the over-determined elliptic problem

$$\begin{cases} \Delta_{\bar{g}} \bar{\phi} + \bar{\lambda} \bar{\phi} = 0 & \text{in } B_{\bar{g},1+v}^+(p) \\ \bar{\phi} = 0 & \text{on } \partial B_{\bar{g},1+v}^+(p) \cap \mathring{M} \\ \bar{g}(\nabla \bar{\phi}, \bar{\nu}) = 0 & \text{on } \partial B_{\bar{g},1+v}^+(p) \cap \partial M \\ \bar{g}(\nabla \bar{\phi}, \bar{\nu}) = \text{constant} & \text{on } \partial B_{\bar{g},1+v}^+(p) \cap \mathring{M} \end{cases} \tag{27}$$

has a nontrivial positive solution, where $\bar{\nu}$ is the normal vector on $\partial B_{\bar{g},1+v}^+(p)$. The relation between the solutions of the two problems is simply given by

$$\phi = \varepsilon^{-n/2} \bar{\phi} \quad \text{and} \quad \lambda = \varepsilon^{-2} \bar{\lambda}.$$

Let us define the coordinates $y = (y^0, y') = (y^0, y^1, \dots, y^n) \in B_1^+$ by

$$\bar{\Psi}(y) := \text{Exp}_{\Phi(y')}^M(\varepsilon y^0 \bar{N}(\bar{\Phi}(y')))$$

where

$$\bar{\Phi}(y') := \text{Exp}_p^{\partial M} \left(\varepsilon \sum_{i=1}^n y^i E_i \right)$$

for $p \in \partial M$, and \bar{N} is the unit normal vector about ∂M for the metric \bar{g} pointing into M . Using Proposition 5.1 of the Appendix, in the new coordinates y the metric \bar{g} can be written as

$$\begin{aligned}\bar{g}_{00} &= 1 \\ \bar{g}_{0j} &= 0 \\ \bar{g}_{ij} &= \delta_{ij} + 2\varepsilon g(\nabla_{E_i} N, E_j) y^0 + \varepsilon^2 R_{0i0j}(y^0)^2 + \varepsilon^2 g(\nabla_{E_i} N, \nabla_{E_j} N)(y^0)^2 \\ &\quad + 2\varepsilon^2 \sum_k R_{k0ij} y^k y^0 + \frac{1}{3}\varepsilon^2 \sum_{k,\ell} \tilde{R}_{ikjl} y^k y^\ell + \mathcal{O}(\varepsilon^3)\end{aligned}\tag{28}$$

for $i, j, k, l = 1, \dots, n$, where R and \tilde{R} are respectively the curvature tensors of M and ∂M , and

$$\begin{aligned}R_{0i0j} &= g(R(N, E_i) N, E_j) \\ R_{k0ij} &= g(R(E_k, N) E_i, E_j) \\ \tilde{R}_{ijkl} &= \tilde{g}(\tilde{R}(E_i, E_k) E_j, E_\ell).\end{aligned}$$

In the coordinates y and the metric \bar{g} , the problem can be continuously extended for $\varepsilon = 0$ and in this case it becomes

$$\left\{ \begin{array}{lcl} \Delta \bar{\phi} + \bar{\lambda} \bar{\phi} & = & 0 \quad \text{in } B_{1+v}^+ \\ \bar{\phi} & = & 0 \quad \text{on } \partial B_{1+v}^+ \cap \mathbb{R}_+^{n+1} \\ \langle \nabla \bar{\phi}, \bar{\nu} \rangle & = & 0 \quad \text{on } \partial B_{1+v}^+ \cap \partial \mathbb{R}_+^{n+1} \end{array} \right. \tag{29}$$

where Δ denotes the usual Laplacian in \mathbb{R}^{n+1} and $\langle \cdot, \cdot \rangle$ the usual scalar product in \mathbb{R}^{n+1} , with the normalization

$$\int_{B_{1+v}^+} \bar{\phi}^2 = 1 \tag{30}$$

and the volume constraint

$$\text{Vol}(B_{1+v}^+) = \text{Vol}(B_1^+).$$

In particular, when $v = 0$ we have

$$\left\{ \begin{array}{lcl} \Delta \phi_1 + \lambda_1 \phi_1 & = & 0 \quad \text{in } B_1^+ \\ \phi_1 & = & 0 \quad \text{on } \partial B_1^+ \cap \mathbb{R}_+^{n+1} \\ \langle \nabla \phi_1, \nu \rangle & = & 0 \quad \text{on } \partial B_1^+ \cap \partial \mathbb{R}_+^{n+1} \end{array} \right. \tag{31}$$

where λ_1 is the first eigenvalue of the unit Euclidean ball and ϕ_1 is the restriction to B_1^+ of the solution to

$$\left\{ \begin{array}{lcl} \Delta \tilde{\phi}_1 + \lambda_1 \tilde{\phi}_1 & = & 0 \quad \text{in } B_1 \\ \tilde{\phi}_1 & = & 0 \quad \text{on } \partial B_1 \end{array} \right..$$

chosen in order to be positive and have $L^2(B_1)$ norm equal to 2.

3.3 Volume constraint and differentiability with respect to (ε, \bar{v})

In this section, we deal with the volume condition (which leads to replace the variable v by \bar{v} subject to the condition of having a zero mean), and prove the differentiability of $(\bar{\lambda}, \bar{\phi})$ with respect to (ε, \bar{v}) . The result is similar to Proposition 3.2 in [28], and we use the same strategy, though we have to pay attention to the singularities at the boundary of our domain. Let us define the space

$$C_{m,NC}^{2,\alpha}(\overline{S_+^n}) := \left\{ v \in C^{2,\alpha}(\overline{S_+^n}), \int_{S_+^n} \bar{v} = 0, \partial_N v = 0 \text{ on } \partial S_+^n \right\},$$

where $\partial_N v = 0$ denotes the 0 Neumann condition at the boundary of S_+^n .

Proposition 3.1. Given a point $p \in \partial M$, there exists $\varepsilon_0 > 0$, locally uniform in p , such that for all $\varepsilon \in (0, \varepsilon_0)$ and all function $\bar{v} \in C_{m, NC}^{2,\alpha}(\overline{S_+^n})$ such that $\|\bar{v}\|_{C^{2,\alpha}(\overline{S_+^n})} \leq \varepsilon_0$, there exists a unique positive function $\bar{\phi} = \bar{\phi}(p, \varepsilon, \bar{v}) \in C^{2,\alpha}(\overline{B_{\bar{g},1+v}^+(p)})$, a constant $\bar{\lambda} = \bar{\lambda}(p, \varepsilon, \bar{v}) \in \mathbb{R}$ and a constant $v_0 = v_0(p, \varepsilon, \bar{v}) \in \mathbb{R}$ such that

$$\text{Vol}_{\bar{g}}(B_{\bar{g},1+v}^+(p)) = \text{Vol}(B_1^+) \quad (32)$$

where $v := v_0 + \bar{v}$ and $\bar{\phi}$ is a solution to the problem

$$\begin{cases} \Delta_{\bar{g}} \bar{\phi} + \bar{\lambda} \bar{\phi} = 0 & \text{in } B_{\bar{g},1+v}^+(p) \\ \bar{\phi} = 0 & \text{on } \partial B_{\bar{g},1+v}^+(p) \cap \mathring{M} \\ \bar{g}(\nabla \bar{\phi}, \bar{v}) = 0 & \text{on } \partial B_{\bar{g},1+v}^+(p) \cap \partial M \end{cases} \quad (33)$$

which is normalized by

$$\int_{B_{\bar{g},1+v}^+(p)} \bar{\phi}^2 \, d\text{vol}_{\bar{g}} = 1. \quad (34)$$

In addition $\bar{\phi}$, $\bar{\lambda}$ and v_0 depend smoothly on the function \bar{v} and the parameter ε , can be extended smoothly to $\varepsilon = 0$ by (29), and in particular $(\bar{\phi}, \bar{\lambda}, v_0) = (\phi_1, \lambda_1, 0)$ when $(\varepsilon, \bar{v}) = (0, 0)$.

Proof. The proof of this result is similar to the proof of Proposition 3.2 in [28], basically based on the implicit function Theorem. Therefore we only describe the differences from [28], which are the choice of coordinates and the regularity theory for the Laplace-Beltrami operator in domains with singularities.

For the choice of coordinates we use the following coordinates: given $(v_0, \bar{v}) \in \mathbb{R} \times C_{m, NC}^{2,\alpha}(\overline{S_+^n})$ and $v = v_0 + \bar{v}$, we consider the parameterization of $B_{\bar{g},1+v}^+(p) = B_{g,\varepsilon(1+v)}^+(p)$ given by

$$\hat{\Psi}(y) := \text{Exp}_{\hat{\Phi}(y')}^M \left(\left(1 + v_0 + \chi(y) \bar{v} \left(\frac{y}{|y|} \right) \right) y^0 N(\hat{\Phi}(y')) \right)$$

where

$$\hat{\Phi}(y') = \text{Exp}_p^{\partial M} \left(\left(1 + v_0 + \chi(y) \bar{v} \left(\frac{y}{|y|} \right) \right) \sum_{i=1}^n y^i E_i \right).$$

Here $y = (y^0, y') \in B_1^+$, χ is a cutoff function identically equal to 0 when $|y| \leq 1/2$ and identically equal to 1 when $|y| \geq 3/4$, introduced to avoid the singularity at the origin of the polar coordinates. In these coordinates the metric

$$\hat{g} := \hat{\Psi}^* \bar{g} \quad (35)$$

can be written as

$$\hat{g} = (1 + v_0)^2 \sum_{i,j} (\delta_{ij} + C^{ij}) dy_i dy_j,$$

where the coefficients $C^{ij} = C_{\varepsilon,v}^{ij} \in C^{1,\alpha}(\overline{B_1^+})$ are functions of y depending on ε , $v = v_0 + \bar{v}$ and the first partial derivatives of v . It is important here to notice that

$$(\varepsilon, v_0, \bar{v}) \mapsto C_{\varepsilon,v}^{ij} \in C^{1,\alpha}(\overline{B_1^+})$$

are smooth maps, as in [28].

Now for all $\psi \in C^{2,\alpha}(B_1^+)$ such that

$$\int_{B_1^+} \psi \phi_1 = 0$$

we define

$$N(\varepsilon, \bar{v}, \psi, v_0) := (\Delta \psi + \lambda_1 \psi + (\Delta_{\hat{g}} - \Delta + \mu)(\phi_1 + \psi), \text{Vol}_{\hat{g}}(B_1^+) - \text{Vol}(B_1^+))$$

where μ is given by

$$\mu = - \int_{B_1^+} \phi_1 (\Delta_{\hat{g}} - \Delta)(\phi_1 + \psi),$$

so that the first entry of N is $L^2(B_1^+)$ -orthogonal to ϕ_1 (for the Euclidean metric). Thanks to the choice of coordinates, the mapping N is a smooth map from a neighborhood of $(0, 0, 0, 0)$ in $[0, \infty) \times C_{m, NC}^{2, \alpha}(\overline{S_+^n}) \times C_{\perp, 0}^{2, \alpha}(\overline{B_1^+}) \times \mathbb{R}$ into a neighborhood of $(0, 0)$ in $C_{\perp}^{0, \alpha}(B_1^+) \times \mathbb{R}$. Here the subscript \perp indicates that the functions in the corresponding space are $L^2(B_1^+)$ -orthogonal to ϕ_1 and the subscript 0 indicates that the functions satisfy the mixed condition at the boundary of B_1^+ . The differential of N with respect to (ψ, v_0) , computed at $(0, 0, 0, 0)$, given by

$$\partial_{(\psi, v_0)} N(0, 0, 0, 0) = (\Delta + \lambda_1, n \operatorname{Vol}(B_1^+))$$

is invertible from $C_{\perp, 0}^{2, \alpha}(\overline{B_1^+}) \times \mathbb{R}$ into $C_{\perp}^{0, \alpha}(\overline{B_1^+}) \times \mathbb{R}$, by Proposition 2.4. Then the implicit function theorem applies as in [28] and completes the proof of the result. \square

3.4 Strategy for the proof of Theorem 1.3

We define the operator

$$F(p, \varepsilon, \bar{v}) = \bar{g}(\nabla \bar{\phi}, \bar{\nu})|_{\partial B_{\bar{g}, 1+v}^+(p) \cap \mathring{M}} - \frac{1}{\operatorname{Vol}_{\bar{g}}(\partial B_{\bar{g}, 1+v}^+(p) \cap \mathring{M})} \int_{\partial B_{\bar{g}, 1+v}^+(p) \cap \mathring{M}} \bar{g}(\nabla \bar{\phi}, \bar{\nu}) \, d\operatorname{vol}_{\bar{g}},$$

where $\bar{\nu}$ denotes the unit normal vector field to $\partial B_{\bar{g}, 1+v}^+(p) \cap \mathring{M}$, $(\bar{\phi}, v_0)$ is the solution of (32)-(33)-(34). Recall that $v = v_0 + \bar{v}$. The operator F is locally well defined in a neighborhood of $(p, 0, 0)$ in $\partial M \times [0, \infty) \times C_{m, NC}^{2, \alpha}(S_+^n)$, and after canonical identification of $\partial B_{\bar{g}, 1+v}^+(p) \cap \mathring{M}$ with S_+^n we can consider that it takes its values in $C^{1, \alpha}(S_+^n)$. Moreover, it is easy to see that the zero mean condition is preserved, and then we will write that F takes its values in $C_m^{1, \alpha}(S_+^n)$.

Our aim is to find $(p, \varepsilon, \bar{v})$ such that $F(p, \varepsilon, \bar{v}) = 0$. Observe that, with this condition, $\bar{\phi} = \bar{\phi}(\varepsilon, \bar{v})$ will be the solution to the problem (27).

Following the proof of the previous result, we have the alternative expression for F :

$$F(p, \varepsilon, \bar{v}) = \hat{g}(\nabla \hat{\phi}, \hat{\nu})|_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}} - \frac{1}{\operatorname{Vol}_{\hat{g}}(\partial B_1^+ \cap \mathbb{R}_+^{n+1})} \int_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}} \hat{g}(\nabla \hat{\phi}, \hat{\nu}) \, d\operatorname{vol}_{\hat{g}},$$

where this time $\hat{\nu}$ is the the unit normal vector field to ∂B_1^+ using the metric \hat{g} defined by (35).

Our aim is to solve the equation

$$F(p, \varepsilon, \bar{v}) = 0$$

for some $(p, \varepsilon, \bar{v})$. The first question we should consider is the following: if we fix a point $p \in \partial M$, can we find for all ε small enough a function $\bar{v} = \bar{v}(\varepsilon)$ in order that

$$F(p, \varepsilon, \bar{v}(\varepsilon)) = 0 ?$$

The answer will be negative, because we will see that the kernel K of

$$\partial_{\bar{v}} F(p, 0, 0) : C_{m, NC}^{2, \alpha}(\overline{S_+^n}) \rightarrow C_m^{1, \alpha}(\overline{S_+^n})$$

is nontrivial. Nevertheless, we will obtain a characterization of K proving that it is given by the space of linear functions (restraint to the half-sphere) depending only on the coordinates y^1, \dots, y^n , i.e. functions

$$\begin{aligned} \overline{S_+^n} &\rightarrow \mathbb{R} \\ y &\rightarrow \langle a, y \rangle \end{aligned}$$

for some $a = (a^0, a) \in \mathbb{R}^{n+1}$ with $a^0 = 0$. Moreover we will prove that $\partial_{\bar{v}} F(p, 0, 0)$ is an isomorphism from K^\perp to the image of $\partial_{\bar{v}} F(p, 0, 0)$, and then the implicit function theorem will give the following result: for all ε small enough there exist an element $k(\varepsilon) \in K$ and a function $\bar{v}(\varepsilon)$ such that

$$F(p, \varepsilon, \bar{v}(\varepsilon)) = k(\varepsilon).$$

Clearly, since we fixed the point p , the function \bar{v} and the element k depend also on p , and in fact we have to write

$$F(p, \epsilon, \bar{v}(p, \epsilon)) = k(p, \epsilon).$$

In the last section we will show that it is possible to apply the implicit function theorem to the equation

$$k(p, \epsilon) = 0$$

obtaining that: for all ϵ small enough, there exists a point p_ϵ such that

$$k(p_\epsilon, \epsilon) = 0.$$

and this will complete the proof of the result.

4 Solving the problem

4.1 Computation of the linearization of F with respect to \bar{v} at $(p, \epsilon, \bar{v}) = (p, 0, 0)$

In Section 3.3 we established the existence of a unique positive function $\bar{\phi} \in C^{2,\alpha}(B_{1+v}^+)$ (close to ϕ_1), a constant $\bar{\lambda} \in \mathbb{R}$ (close to λ_1) and a constant $v_0 \in \mathbb{R}$ (close to 0), solutions to (32)-(33)-(34). Recall that λ_1 is the first eigenvalue of $-\Delta$ in the half ball B_1^+ with 0 mixed boundary condition and ϕ_1 is the associated eigenfunction which is normalized to be positive and have $L^2(B_1^+)$ norm equal to 1. For all $\bar{v} \in C_{m,NC}^{2,\alpha}(\overline{S_+^n})$ let ψ be the (unique) solution of

$$\begin{cases} \Delta\psi + \lambda_1\psi = 0 & \text{in } B_1^+ \\ \psi = -\partial_r\phi_1 \bar{v} & \text{on } \partial B_1^+ \cap \mathbb{R}_+^{n+1} \\ \langle \nabla\psi, \nu \rangle = 0 & \text{on } \partial B_1^+ \cap \partial\mathbb{R}_+^{n+1} \end{cases} \quad (36)$$

which is $L^2(B_1^+)$ -orthogonal to ϕ_1 . We define

$$L_0(\bar{v}) := (\partial_r\psi + \partial_r^2\phi_1 \bar{v})|_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}} \quad (37)$$

Clearly we have

$$L_0 : C_{m,NC}^{2,\alpha}(\overline{S_+^n}) \rightarrow C_m^{1,\alpha}(\overline{S_+^n}).$$

Proposition 4.1. *The linearization the operator F with respect to \bar{v} computed at $(p, 0, 0)$, i.e.*

$$\partial_{\bar{v}} F(p, 0, 0),$$

is equal to L_0 .

Proof. When $\epsilon = 0$ we have already seen that \bar{g} in the coordinates y is the Euclidean metric. If $v \in C_m^{2,\alpha}(S^n)$ we can define the operator F :

$$\tilde{F}(v) = \langle \nabla \tilde{\phi}, \tilde{\nu} \rangle|_{\partial B_{1+v}} - \frac{1}{\text{Vol}(\partial B_{1+v})} \int_{\partial B_{1+v}} \langle \nabla \tilde{\phi}, \tilde{\nu} \rangle,$$

where $\tilde{\nu}$ denotes the unit normal vector field to ∂B_{1+v} and $\tilde{\phi}$ is the solution, with L^2 -norm equal to 2, of

$$\begin{cases} \Delta \tilde{\phi} + \bar{\lambda} \tilde{\phi} = 0 & \text{in } B_{1+v} \\ \tilde{\phi} = 0 & \text{on } \partial B_{1+v} \end{cases}. \quad (38)$$

After identification of ∂B_{1+v} with S^n we can considered the operator \tilde{F} well defined from $C_m^{2,\alpha}(S^n)$ into $C_m^{1,\alpha}(S^n)$. In the proof of Proposition 4.3 in [28] it is proved that the linearization of \tilde{F} with respect to v at $v = 0$ is given by the operator

$$\begin{aligned} \tilde{L}_0 & : C_m^{2,\alpha}(S^n) & \longrightarrow & C_m^{1,\alpha}(S^n) \\ v & \mapsto & \left(\partial_r \tilde{\psi} + \partial_r^2 \tilde{\phi}_1 v \right)|_{\partial B_1} \end{aligned} \quad (39)$$

where $\tilde{\phi}_1$ is the first eigenfunction of $-\Delta$ in B_1 with 0 Dirichlet boundary condition and normalized to have L^2 -norm equal to 2, and $\tilde{\psi}$ is the (unique) solution of

$$\begin{cases} \Delta \tilde{\psi} + \lambda_1 \tilde{\psi} = 0 & \text{in } B_1 \\ \tilde{\psi} = -\partial_r \tilde{\phi}_1 v & \text{on } \partial B_1 \end{cases} \quad (40)$$

which is $L^2(B_1)$ -orthogonal to $\tilde{\phi}_1$. Notice that ϕ_1 and ψ are then the restrictions of $\tilde{\phi}_1$ and $\tilde{\psi}$ to the half-ball B_1^+ .

Let w be a function in $C_{m,NC}^{2,\alpha}(\overline{S_+^n})$. We extend the function w to a function \tilde{w} over all S^n in this way: for $(y^0, y^1, \dots, y^n) \in S_+^n$ we set

$$\tilde{w}(-y^0, y^1, \dots, y^n) = w(y^0, y^1, \dots, y^n).$$

Observe that $\tilde{w} \in C^{2,\alpha}(S^n)$ because the function w satisfies the Neumann condition at the boundary of S_+^n , and his mean is 0 because are 0 the means over S_+^n and over the complement of S_+^n . We conclude that $\tilde{w} \in C_{m,Sym}^{2,\alpha}(S^n)$, where the subscript *Sym* means that the function is symmetric with respect to the hyperplane $\{x_0 = 0\}$, and m means as usual that the function has mean 0. We have defined the mapping

$$\begin{aligned} \alpha : C_{m,NC}^{2,\alpha}(\overline{S_+^n}) &\longrightarrow C_{m,Sym}^{2,\alpha}(S^n) \\ w &\mapsto \tilde{w} \end{aligned}, \quad (41)$$

and it is easy to see that this mapping is an isomorphism.

If we consider the operator \tilde{F} defined only in $C_{m,Sym}^{2,\alpha}(S^n)$, it is natural that its linearization with respect to v at $v = 0$ is given by the operator \tilde{L}_0 restricted to $C_{m,Sym}^{2,\alpha}(S^n)$ with image in $C_{m,Sym}^{1,\alpha}(S^n)$. We observe that if $v \in C_{m,Sym}^{2,\alpha}(S^n)$, then the solution of (38) is symmetric with respect to the hyperplain $\{x_0 = 0\}$ and the normal derivative with respect to x_0 computed at $\{x_0 = 0\}$ is 0. Then from the definitions of F and \tilde{F} we conclude that

$$F(p, 0, \bar{v}) = \tilde{F}(\alpha(\bar{v}))|_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}}$$

where α is the isomorphism defined in (41). We define also the mapping

$$\begin{aligned} \beta : C_{m,Sym}^{1,\alpha}(S^n) &\longrightarrow C_{m,NC}^{1,\alpha}(S_+^n) \\ v &\mapsto v|_{S_+^n} \end{aligned}$$

and we observe that it is an isomorphism. We claim that

$$L_0 = \beta \circ \tilde{L}_0 \circ \alpha.$$

We remark that the operator $\beta \circ \tilde{L}_0 \circ \alpha$ is defined on $C_{m,NC}^{2,\alpha}(\overline{S_+^n})$ and his image is contained in $C_{m,NC}^{1,\alpha}(S_+^n)$. We have to prove that

$$L_0(w) = \tilde{L}_0(\tilde{w})|_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}}$$

By the symmetry of the function \tilde{w} with respect to the hyperplane $\{x_0 = 0\}$, we conclude that the solution of (40) with $v = \tilde{w}$ is symmetric with respect to the hyperplane $\{x_0 = 0\}$, then $\partial_{x_0} \tilde{\psi}|_{\{x_0=0\}} = 0$ and $\tilde{L}_0(\tilde{w})$ is symmetric with respect to the hyperplane $\{x_0 = 0\}$. So the restriction of $\tilde{\psi}$ to the half-ball B_1^+ is the solution of (36), where $\bar{v} = w$, and $L_0(w)$ is exactly the restriction of $\tilde{L}_0(\tilde{w})$ to $\partial B_1^+ \cap \mathbb{R}_+^{n+1}$. This completes the proof of the claim. Using this relation we conclude that that

$$L_0(w) = \tilde{L}_0(\alpha(w))|_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}}.$$

This completes the proof of the proposition. \square

4.2 Study of the operator L_0

Proposition 4.2. *The operator*

$$L_0 : C_{m,NC}^{2,\alpha}(\overline{S_+^n}) \longrightarrow C_{m,NC}^{1,\alpha}(S_+^n),$$

is a self adjoint, first order elliptic operator. Its kernel K is given by the space of linear functions depending only on the coordinates y^1, \dots, y^n , i.e. functions

$$\begin{array}{ccc} \overline{S_+^n} & \rightarrow & \mathbb{R} \\ y & \rightarrow & \langle a, y \rangle \end{array}$$

for some $a = (a^0, a') \in \mathbb{R}^{n+1}$ with $a^0 = 0$. Moreover, L_0 has closed range and is an isomorphism from K^\perp to $\text{Im}(L_0)$, where K^\perp is the space L^2 -orthogonal to K in $C_{m,NC}^{2,\alpha}(\overline{S_+^n})$ and $\text{Im}(L_0)$ denotes the range of L_0 in $C_{m,NC}^{1,\alpha}(S_+^n)$.

Proof. Let \tilde{L}_0 the operator defined in (39) and α the isomorphism defined in (41). In Proposition 4.2 of [28] it is proved that:

- \tilde{L}_0 is a self adjoint, first order elliptic operator,
- its kernel is given by the space of linear functions restraint to S^n , and
- there exists a constant $c > 0$ such that

$$\|v\|_{C^{2,\alpha}(S^n)} \leq c \|\tilde{L}_0(v)\|_{C^{1,\alpha}(S^n)}, \quad (42)$$

provided that v is $L^2(S^n)$ -orthogonal to the kernel of \tilde{L}_0 .

The last elliptic estimate implies that the operator \tilde{L}_0 has closed range, and using the other two properties we have that \tilde{L}_0 is an isomorphism from the space L^2 -orthogonal to its kernel and its range.

We are interested in considering the operator \tilde{L}_0 defined only in the domain $C_{m,Sym}^{2,\alpha}(S^n)$ and from now on \tilde{L}_0 will be defined only in $C_{m,Sym}^{2,\alpha}(S^n)$. The image of \tilde{L}_0 is naturally given by functions that are symmetric with respect to the hyperplane $\{x_0 = 0\}$, then we have

$$\tilde{L}_0 : C_{m,Sym}^{2,\alpha}(S^n) \longrightarrow C_{m,Sym}^{1,\alpha}(S^n)$$

We can conclude that the new operator \tilde{L}_0 is a self-adjoint, first order elliptic operator, with kernel \tilde{K} given by the space of linear functions which are symmetric with respect to the hyperplane $\{x_0 = 0\}$, i.e. functions

$$\begin{array}{ccc} \overline{S_+^n} & \rightarrow & \mathbb{R} \\ y & \rightarrow & \langle a, y \rangle \end{array}$$

for some $a = (a^0, a') \in \mathbb{R}^{n+1}$ with $a^0 = 0$. Inequality (42) holds naturally also for the new operator \tilde{L}_0 , provided v is $L^2(S^n)$ -orthogonal to \tilde{K} .

From the proof of Proposition 4.1 we have

$$L_0 = \beta \circ \tilde{L}_0 \circ \alpha.$$

With this characterization of the operator L_0 and the properties of \tilde{L}_0 , we deduce that the kernel of L_0 is given by the space K of functions

$$\begin{array}{ccc} \overline{S_+^n} & \rightarrow & \mathbb{R} \\ y & \rightarrow & \langle a, y \rangle \end{array}$$

for some $a = (a^0, a') \in \mathbb{R}^{n+1}$ with $a^0 = 0$, and that L_0 has closed range and is an isomorphism from K^\perp to $\text{Im}(L_0)$, where K^\perp is the space L^2 -orthogonal to K in $C_{m,NC}^{2,\alpha}(\overline{S_+^n})$ and $\text{Im}(L_0)$ denotes the range of L_0 in $C_{m,NC}^{1,\alpha}(S_+^n)$. \square

4.3 Solving the problem on the space orthogonal to the kernel of L_0

Lemma 4.3. *Let $p \in \partial M$. There exists a function $f_p \in C^{1,\alpha}([0, 1])$ such that*

$$F(p, \varepsilon, 0)(y^0, y') = \varepsilon f_p(y^0) + \mathcal{O}(\varepsilon^2)$$

for all ε small enough.

Proof. We keep the notations of the proof of the Proposition 3.1 with $\bar{v} \equiv 0$. Since $\bar{v} \equiv 0$, we have

$$N(\varepsilon, 0, 0, 0) = ((\Delta_{\hat{g}} - \Delta + \mu) \phi_1, \text{Vol}_{\hat{g}}(B_1^+) - \text{Vol}(B_1^+)) ,$$

and

$$\mu = - \int_{B_1^+} \phi_1 (\Delta_{\hat{g}} - \Delta) \phi_1 .$$

If in addition $v_0 = 0$, we can estimate

$$\hat{g}_{ij} = \delta_{ij} + \hat{G}_{ij} \varepsilon y^0 + \mathcal{O}(\varepsilon^2) ,$$

where \hat{G}_{ij} are real constants. Hence, by the symmetry of the problem,

$$N(\varepsilon, 0, 0, 0)(y^0, y') = \varepsilon (\varphi(y^0, |y'|), V) + \mathcal{O}(\varepsilon^2) ,$$

where the first component of $\varphi \in C^{0,\alpha}([0, 1]^2)$ and V is a real number. The implicit function theorem immediately implies that the solution of

$$N(\varepsilon, v_0, 0, \psi) = 0$$

satisfies

$$\|\psi(\varepsilon, p, 0)\|_{C^{2,\alpha}} + |v_0(\varepsilon, p, 0)| \leq c \varepsilon$$

but in addition there exist a function $\tilde{\psi}_p \in C^{2,\alpha}([0, 1]^2)$ such that

$$\psi(\varepsilon, p, 0)(y^0, y') = \varepsilon \tilde{\psi}_p(y^0, |y'|) + \mathcal{O}(\varepsilon^2) .$$

To complete the proof, observe that $\hat{\nu} = (1 + v_0)^{-1} \partial_r$ on $\partial B_1^+ \cap \mathbb{R}_+^{n+1}$ when $\bar{v} \equiv 0$. Therefore there exist a function $\hat{f}_p \in C^{2,\alpha}([0, 1]^2)$ such that

$$\hat{g}(\nabla \hat{\phi}, \hat{\nu})(y^0, y') = \partial_r \phi_1 + \varepsilon \hat{f}_p(y^0, |y'|) + \mathcal{O}(\varepsilon^2) .$$

(be careful that \hat{g} is defined with $v_0 = v_0(\varepsilon, p, 0)$ and $\bar{v} \equiv 0$). Since $\partial_r \phi_1$ is constant along $\partial B_1^+ \cap \mathbb{R}_+^{n+1}$, we conclude that there exist a function $f_p \in C^{2,\alpha}([0, 1])$ such that

$$F(p, \varepsilon, 0)(y^0, y') = \varepsilon f_p(y^0) + \mathcal{O}(\varepsilon^2) .$$

This completes the proof of the Lemma. \square

Proposition 4.4. *There exists $\varepsilon_0 > 0$ such that, for all $\varepsilon \in [0, \varepsilon_0]$ and for all p in a compact subset of ∂M , there exists a unique function $\bar{v} = \bar{v}(p, \varepsilon) \in K^\perp$ such that*

$$F(p, \varepsilon, \bar{v}(p, \varepsilon)) \in K .$$

The function $\bar{v}(p, \varepsilon)$ depends smoothly on p and ε and

$$\bar{v}(p, \varepsilon)(y^0, y') = \varepsilon \tilde{v}_p(y^0) + \mathcal{O}(\varepsilon^2)$$

for a suitable function $\tilde{v}_p \in C^{2,\alpha}([0, 1])$.

Proof. We fix p in a compact subset of ∂M and define

$$\bar{F}(p, \varepsilon, \bar{v}, a) := F(p, \varepsilon, \bar{v}) + \langle a, \cdot \rangle$$

By Proposition 3.1, \bar{F} is a C^1 map from a neighborhood of $(p, 0, 0, 0)$ in $M \times [0, \infty) \times K^\perp \times \partial \mathbb{R}_+^{n+1}$ into a neighborhood of 0 in $C^{1,\alpha}(S^n_+)$. Moreover we have

- $\bar{F}(p, 0, 0, 0) = 0$,
- the differential of \bar{F} with respect to \bar{v} computed at $(p, 0, 0, 0)$ is given by L_0 restricted to K^\perp , and
- the image of the linear map $a \mapsto \langle a, \cdot \rangle$, $a = (a_0, a')$ with $a_0 = 0$ coincides with K .

Thanks to the result of Proposition 4.2, the implicit function theorem can be applied to the equation

$$\bar{F}(p, \varepsilon, \bar{v}, a) = 0$$

at $(p, 0, 0, 0)$ with respect to the variable ε . We obtain the existence of $\bar{v}(p, \varepsilon) \in C_{m, NC}^{2,\alpha}(S_+^n)$ and $a(p, \varepsilon) \in \partial\mathbb{R}_+^{n+1}$, smoothly depending on ε such that

$$\bar{F}(p, \varepsilon, \bar{v}(p, \varepsilon), a(p, \varepsilon)) = 0,$$

that means, by the definition of \bar{F} ,

$$F(p, \varepsilon, \bar{v}(p, \varepsilon)) \in K.$$

The fact that \bar{v} depends smoothly on p and ε is standard. The ε -expansion of \bar{v} follow at once from Lemma 4.3. \square

4.4 Projecting over the kernel of L_0 : appearance of the mean curvature of ∂M

Thanks to Proposition 4.4 we are able to build, for all p in a compact subset of ∂M and ε small enough, a function $\bar{v}(p, \varepsilon)$ in K^\perp such that

$$F(p, \varepsilon, \bar{v}(p, \varepsilon)) \in K.$$

Now, as natural, we project the operator F over its K and we then we have to find, for each ε , the good point p_ε in order that such the projection of F over K is equal to 0. In other words, for all ε small enough we want to find a point $p_\varepsilon \in \partial M$ such that

$$\int_{S_+^n} F(p_\varepsilon, \varepsilon, \bar{v}(p_\varepsilon, \varepsilon)) \langle b, \cdot \rangle = 0$$

for all $b \in \partial\mathbb{R}_+^{n+1}$. The main result of this section is the following:

Proposition 4.5. *For all $p \in \partial M$ and all $b = (0, b') \in \partial\mathbb{R}_+^{n+1}$ with $|b| = 1$, we have the following ε -expansion:*

$$\int_{S_+^n} F(p, \varepsilon, \bar{v}(p, \varepsilon)) \langle b, \cdot \rangle = C \varepsilon^2 \tilde{g}(\nabla^{\tilde{g}} H(p), \Theta(b')) + \mathcal{O}(\varepsilon^3).$$

where C is a real constant, H is the mean curvature of ∂M , \tilde{g} is the metric of ∂M induced by g and Θ has been defined in (23).

Proof. Take $p \in \partial M$, ε small enough, $\bar{v} \in C_{m, NC}^{2,\alpha}$ with small norm, and $b \in \partial\mathbb{R}_+^{n+1}$. We denote by L_ε the linearization of F with respect to \bar{v} , and by L_ε^2 the second derivative of F with respect to \bar{v} , both computed at the point $(p, \varepsilon, 0)$:

$$L_\varepsilon = \partial_{\bar{v}} F(p, \varepsilon, 0) \quad \text{and} \quad L_\varepsilon^2 = \partial_{\bar{v}}^2 F(p, \varepsilon, 0).$$

We have

$$\int_{S_+^n} F(p, \varepsilon, \bar{v}) \langle b, \cdot \rangle = \int_{S_+^n} (F(p, \varepsilon, 0) + L_0 \bar{v}) \langle b, \cdot \rangle + \int_{S_+^n} (F(p, \varepsilon, \bar{v}) - F(p, \varepsilon, 0) - L_\varepsilon \bar{v}) \langle b, \cdot \rangle + \int_{S_+^n} (L_\varepsilon - L_0) \bar{v} \langle b, \cdot \rangle$$

Now we apply this formula for our function $\bar{v} = \bar{v}(p, \varepsilon)$ given by Proposition 4.4. We have $\bar{v} \in K^\perp$, so $L_0 \bar{v} \in K^\perp$, and then

$$\int_{S_+^n} L_0 \bar{v} \langle b, \cdot \rangle = 0.$$

We obtain that

$$\int_{S_+^n} F(p, \varepsilon, \bar{v}) \langle b, \cdot \rangle = \int_{S_+^n} F(p, \varepsilon, 0) \langle b, \cdot \rangle + \int_{S_+^n} (F(p, \varepsilon, \bar{v}) - F(p, \varepsilon, 0) - L_\varepsilon \bar{v}) \langle b, \cdot \rangle + \int_{S_+^n} (L_\varepsilon - L_0) \bar{v} \langle b, \cdot \rangle \quad (43)$$

where $\bar{v} = \bar{v}(p, \varepsilon)$ is the function given by Proposition 4.4. We need now two intermediate lemmas.

Lemma 4.6. For all $p \in \partial M$, for all $b = (0, b') \in \partial \mathbb{R}_+^{n+1}$ we have the following ε -expansion:

$$\int_{S_+^n} F(p, \varepsilon, 0) \langle b, \cdot \rangle = C \varepsilon^2 \tilde{g}(\nabla^{\tilde{g}} H(p), \Theta(b')) + |b| \mathcal{O}(\varepsilon^3),$$

where Θ is defined in (23) and

$$C = -2 \left(\int_{S_+^n} y^0 (y^1)^2 \right) \frac{1}{\partial_r \phi_1(1)} \int_{B_1^+} r |\partial_r \phi_1|^2$$

where $r = |y|$.

Proof. We recall that

$$F(p, \varepsilon, \bar{v}) = \hat{g}(\nabla \hat{\phi}, \hat{\nu})|_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}} - \frac{1}{\text{Vol}_{\hat{g}}(\partial B_1^+ \cap \mathbb{R}_+^{n+1})} \int_{\partial B_1^+ \cap \mathbb{R}_+^{n+1}} \hat{g}(\nabla \hat{\phi}, \hat{\nu}) \, d\text{vol}_{\hat{g}},$$

where the metric \hat{g} has been defined in (35) for the coordinates y . Then

$$\int_{S_+^n} F(p, \varepsilon, \bar{v}) \langle b, \cdot \rangle = \int_{S_+^n} \hat{g}(\nabla \hat{\phi}, \hat{\nu}) \langle b, \cdot \rangle.$$

When $\bar{v} = 0$ we have $\hat{\nu} = (1 + v_0) \partial_r$ on $\partial B_1^+ \cap \mathbb{R}_+^{n+1}$, where $r = |y|$. Then

$$\int_{S_+^n} F(p, \varepsilon, \bar{v}) \langle b, \cdot \rangle = (1 + v_0) \int_{S_+^n} \frac{\partial \hat{\phi}}{\partial r} \langle b, \cdot \rangle = \frac{1 + v_0}{\partial_r \phi_1(1)} \int_{S_+^n} \frac{\partial \hat{\phi}}{\partial r} \langle \nabla \phi_1, b \rangle \quad (44)$$

where we used the fact that ϕ_1 is a radial function. Using this last property and the Green's identities we have:

$$\begin{aligned} \int_{S_+^n} \frac{\partial \hat{\phi}}{\partial r} \langle \nabla \phi_1, b \rangle &= \int_{B_1^+} (\Delta + \lambda_1) \hat{\phi} \langle \nabla \phi_1, b \rangle - \int_{B_1^+} \hat{\phi} (\Delta + \lambda_1) \langle \nabla \phi_1, b \rangle \\ &= \int_{B_1^+} (\Delta + \lambda_1) \hat{\phi} \langle \nabla \phi_1, b \rangle \\ &= \int_{B_1^+} (\Delta - \Delta_{\hat{g}}) \hat{\phi} \langle \nabla \phi_1, b \rangle + (\lambda_1 - \hat{\lambda}) \int_{B_1^+} \hat{\phi} \langle \nabla \phi_1, b \rangle \\ &= \int_{B_1^+} (\Delta - \Delta_{\hat{g}}) \phi_1 \langle \nabla \phi_1, b \rangle + \int_{B_1^+} (\Delta - \Delta_{\hat{g}}) (\hat{\phi} - \phi_1) \langle \nabla \phi_1, b \rangle + (\lambda_1 - \hat{\lambda}) \int_{B_1^+} (\hat{\phi} - \phi_1) \langle \nabla \phi_1, b \rangle \end{aligned}$$

Let compute the first term. Recall that

$$\Delta_{\hat{g}} := \sum_{i,j=0}^n \hat{g}^{ij} \partial_{y_i} \partial_{y_j} + \sum_{i,j=0}^n \partial_{y_i} \hat{g}^{ij} \partial_{y_j} + \frac{1}{2} \sum_{i,j=0}^n \hat{g}^{ij} \partial_{y_i} \log |\hat{g}| \partial_{y_j}.$$

From (28) we have that the coefficients of the metric \hat{g} can be expanded, for $i, k, j, \ell = 1, \dots, n$, as

$$\begin{aligned} \hat{g}_{00}(y) &= (1 + v_0)^2 \\ \hat{g}_{0j}(y) &= 0 \\ \hat{g}_{ij}(y) &= (1 + v_0)^2 \left(\delta_{ij} + 2(1 + v_0) \varepsilon g(\nabla_{E_i} N, E_j) y^0 + R_{0i0j} (1 + v_0)^2 \varepsilon^2 (y^0)^2 \right. \\ &\quad \left. + (1 + v_0)^2 \varepsilon^2 g(\nabla_{E_i} N, \nabla_{E_j} N) (y^0)^2 + 2(1 + v_0)^2 \varepsilon^2 \sum_k R_{k0ij} y^k y^0 \right. \\ &\quad \left. + \frac{1}{3} (1 + v_0)^2 \varepsilon^2 \sum_{k,\ell} \tilde{R}_{ikj\ell} y^k y^\ell + \mathcal{O}(\varepsilon^3) \right) \end{aligned}$$

Keeping in mind that $v_0 = v_0(p, \varepsilon) = \mathcal{O}(\varepsilon)$, the third equality simplifies slightly obtaining

$$\begin{aligned}\hat{g}^{00}(y) &= (1 + v_0)^{-2} \\ \hat{g}^{0j}(y) &= 0 \\ \hat{g}^{ij}(y) &= (1 + v_0)^{-2} \left(\delta_{ij} - 2(1 + v_0) \varepsilon g(\nabla_{E_i} N, E_j) y^0 - R_{0i0j} \varepsilon^2 (y^0)^2 \right. \\ &\quad \left. - \varepsilon^2 g(\nabla_{E_i} N, \nabla_{E_j} N) (y^0)^2 - 2\varepsilon^2 \sum_k R_{k0ij} y^k y^0 - \frac{1}{3} \varepsilon^2 \sum_{k,\ell} \tilde{R}_{ikj\ell} y^k y^\ell \right) + \mathcal{O}(\varepsilon^3).\end{aligned}$$

Using the fact that $R_{k0ii} = 0$, we have

$$\begin{aligned}\log |\hat{g}| &= 2n \log(1 + v_0) - 2\varepsilon(1 + v_0) H(p) y^0 \\ &\quad + \frac{\varepsilon^2}{2} \left\{ \left[-\text{Ric}(N) + 4 \sum_{i \neq j} g(\nabla_{E_i} N, E_i) g(\nabla_{E_j} N, E_j) + \sum_i g(\nabla_{E_i} N, \nabla_{E_i} N) \right. \right. \\ &\quad \left. \left. - 4 \sum_{i \neq j} g(\nabla_{E_i} N, E_j) g(\nabla_{E_j} N, E_i) \right] (y^0)^2 + \frac{1}{3} \tilde{R}_{k\ell} y^k y^\ell \right\} + \mathcal{O}(\varepsilon^3)\end{aligned}$$

where Ric denotes the Ricci curvature of ∂M and

$$\tilde{R}_{k\ell} = \sum_{i=1}^n \tilde{R}_{ikil}.$$

A straightforward computation (still keeping in mind that $v_0 = \mathcal{O}(\varepsilon)$) shows that

$$\begin{aligned}(\Delta - \Delta_{\hat{g}})\phi_1 &= -\lambda_1 (1 - (1 + v_0)^{-2}) \phi_1 \\ &\quad + 2(1 + v_0)^{-1} \varepsilon \sum_{i,j} g(\nabla_{E_i} N, E_j) y^0 \left(\frac{y^i y^j}{r^2} \partial_r^2 \phi_1 + \frac{\delta_j^i}{r} \partial_r \phi_1 - \frac{y^i y^j}{r^3} \partial_r \phi_1 \right) \\ &\quad + \varepsilon (1 + v_0)^{-1} H(p) \frac{y^0}{r} \partial_r \phi_1 \\ &\quad + \varepsilon^2 \sum_{k,i,j,\ell} \left[[R_{0i0j} + g(\nabla_{E_i} N, \nabla_{E_j} N)] (y^0)^2 + 2R_{k0ij} y^k y^0 + \frac{1}{3} \tilde{R}_{ikj\ell} y^k y^\ell \right] \cdot \left(\frac{y^i y^j}{r^2} \partial_r^2 \phi_1 + \frac{\delta_j^i}{r} \partial_r \phi_1 - \frac{y^i y^j}{r^3} \partial_r \phi_1 \right) \\ &\quad + \varepsilon^2 \sum_{k,i,j} \left(2R_{i0ij} y^0 + \frac{1}{3} \tilde{R}_{ikj\ell} y^k + \frac{1}{6} \tilde{R}_{ik} y^k \right) \frac{y^j}{r} \partial_r \phi_1 \\ &\quad + \varepsilon^2 \left[-\text{Ric}(N) + 4 \sum_{i \neq j} g(\nabla_{E_i} N, E_i) g(\nabla_{E_j} N, E_j) + \sum_i g(\nabla_{E_i} N, \nabla_{E_i} N) \right. \\ &\quad \left. - 4 \sum_{i \neq j} g(\nabla_{E_i} N, E_j) g(\nabla_{E_j} N, E_i) \right] \cdot \frac{(y^0)^2}{r} \partial_r \phi_1\end{aligned}$$

where $i, j, k = 1, \dots, n$. Observe that we have used the fact that $R(X, X) \equiv 0$ and the symmetries of the curvature tensor for which $R_{ijkl} = R_{klji}$. Now, in the computation of

$$\int_{B_1^+} (\Delta - \Delta_{\hat{g}}) \phi_1 \langle \nabla \phi_1, b \rangle,$$

observe that the terms in the expansion of $(\Delta - \Delta_{\hat{g}}) \phi_1$ which contain an even number of coordinates different to y^0 , such as y^0 or $y^i y^j y^k y^\ell$ or $(y^0)^2 y^i y^j$ etc. do not contribute to the result since, once multiplied by $\langle \nabla \phi_1, b \rangle$ (keep

in mind that $b = (0, b')$, their average over S_+^n is 0. Therefore, we can write

$$\begin{aligned} \int_{B_1^+} (\Delta - \Delta_{\hat{g}}) \phi_1 \langle \nabla \phi_1, b \rangle &= \varepsilon^2 \sum_{\sigma \neq 0} \int_{B_1^+} \partial_r \phi_1 a_\sigma \frac{y^\sigma}{r} \cdot \\ &\quad \cdot \left(2 \sum_{k,i,j} R_{k0ij} \left(\frac{y^i y^j y^k y^0}{r^2} \partial_r^2 \phi_1 - \frac{y^i y^j y^k y^0}{r^3} \partial_r \phi_1 \right) + 2 \sum_{k,i,j} R_{i0ij} \frac{y^0 y^j}{r} \partial_r \phi_1 \right) \\ &\quad + \mathcal{O}(\varepsilon^3) \end{aligned}$$

We make use of the technical Lemmas 5.2 and 5.3 of the Appendix to conclude that

$$\int_{B_1^+} (\Delta - \Delta_{\hat{g}}) \phi_1 \langle \nabla \phi_1, b \rangle = \tilde{C} \varepsilon^2 \tilde{g}(\nabla^{\tilde{g}} H(p), \Theta(b')) + \mathcal{O}(\varepsilon^3). \quad (45)$$

where

$$\tilde{C} = -2 \left(\int_{S_+^n} y^0 (y^1)^2 \right) \int_{B_1^+} r |\partial_r \phi_1|^2.$$

Now we have to compute the terms

$$\int_{B_1^+} (\Delta - \Delta_{\hat{g}}) (\hat{\phi} - \phi_1) \langle \nabla \phi_1, b \rangle \quad \text{and} \quad (\lambda_1 - \hat{\lambda}) \int_{B_1^+} (\hat{\phi} - \phi_1) \langle \nabla \phi_1, a \rangle.$$

We observe that the coefficients of the metric, for $i, j = 1, \dots, n$, are given by

$$\hat{g}_{ij}(y) = \delta_{ij} + \hat{G}_{ij} \varepsilon y^0 + \mathcal{O}(\varepsilon^2)$$

for some constants \hat{G}_{ij} . Then the ε -first order term of $\hat{\phi} - \phi_1$ is radial in the coordinates y^1, \dots, y^n , i.e. there exists a function $h \in C^{2,\alpha}([0, 1]^2)$ such that

$$(\hat{\phi} - \phi_1)(y^0, y') = \varepsilon h(y^0, |y'|) + \mathcal{O}(\varepsilon^2).$$

Let $\rho := |y'|$. Using the same computation given above, we find

$$\begin{aligned} (\Delta - \Delta_{\hat{g}})(\hat{\phi} - \phi_1) &= (1 - (1 + v_0)^{-2}) \Delta (\hat{\phi} - \phi_1) \\ &\quad + \mathcal{O}(\varepsilon^2) \left(\frac{y^0 y^i y^j}{\rho^2} \partial_\rho^2 h + \frac{y^0}{\rho} \delta_j^i \partial_\rho h - \frac{y^0 y^i y^j}{\rho^3} \partial_\rho h + \partial_{y^0} h \right) + \mathcal{O}(\varepsilon^3) \\ &= \mathcal{O}(\varepsilon^2) \left(\tilde{h}(y^0, \rho) + \frac{y^0 y^i y^j}{\rho^2} \partial_\rho^2 h + \frac{y^0}{\rho} \delta_j^i \partial_\rho h - \frac{y^0 y^i y^j}{\rho^3} \partial_\rho h + \partial_{y^0} h \right) + \mathcal{O}(\varepsilon^3) \end{aligned}$$

for some function $\tilde{h} \in C^{0,\alpha}([0, 1]^2)$, and the terms $\mathcal{O}(\varepsilon^2)$ do not depend on the coordinates. As in the previous computation, terms which contain an even number of coordinates different to y^0 do not contribute to the result since, once multiplied by $\langle \nabla \phi_1, b \rangle$, their average over S_+^n is 0. Therefore

$$\int_{B_1^+} (\Delta - \Delta_{\hat{g}})(\hat{\phi} - \phi_1) \langle \nabla \phi_1, b \rangle = \mathcal{O}(\varepsilon^3).$$

For the last term we have to estimate, the previous computation immediately implies that

$$\int_{B_1^+} (\hat{\phi} - \phi_1) \langle \nabla \phi_1, b \rangle = \mathcal{O}(\varepsilon^2)$$

and then

$$(\lambda_1 - \hat{\lambda}) \int_{B_1^+} (\hat{\phi} - \phi_1) \langle \nabla \phi_1, b \rangle = \mathcal{O}(\varepsilon^3).$$

We conclude that

$$\int_{S_+^n} \frac{\partial \hat{\phi}}{\partial r} \langle \nabla \phi_1, b \rangle = \int_{B_1^+} (\Delta - \Delta_{\tilde{g}}) \phi_1 \langle \nabla \phi_1, b \rangle + |b| \mathcal{O}(\varepsilon^3) = \tilde{C} \varepsilon^2 \tilde{g}(\nabla^{\tilde{g}} H(p), \Theta(b')) + |b| \mathcal{O}(\varepsilon^3).$$

The Lemma follows at once from (44), keeping in mind that $v_0 = \mathcal{O}(\varepsilon)$. \square

Lemma 4.7. *Let $\bar{v} = \bar{v}(p, \varepsilon) \in C_{m, NC}^{2,\alpha}(S_+^n)$ such that in the coordinates $y = (y^0, y')$ we have*

$$\bar{v}(y^0, y') = \varepsilon \tilde{v}_p(y^0) + \mathcal{O}(\varepsilon^2)$$

for some function $\tilde{v}_p \in C^{2,\alpha}([0, 1])$. Then there exist two functions $\delta_p, \sigma_p \in C^{2,\alpha}([0, 1])$ such that

$$((L_\varepsilon - L_0) \bar{v})(y^0, y') = \varepsilon^2 \delta_p(y^0) + \mathcal{O}(\varepsilon^3)$$

and

$$F(p, \varepsilon, \bar{v}) - F(p, \varepsilon, 0) - L_\varepsilon \bar{v} = \varepsilon^2 \sigma_p(y^0) + \mathcal{O}(\varepsilon^3).$$

Proof. Clearly both L_ε and L_0 are first order differential operators, and the dependence on ε is smooth. Now, the difference between the coefficients of \bar{g} written in the coordinates y defined in (28) and the coefficient of the Euclidean metric can be estimated by

$$\bar{g}_{ij}(y^0, y') = \bar{G}_{ij} \varepsilon y^0 + \mathcal{O}(\varepsilon^2)$$

If the function \bar{v} is such that

$$\bar{v}(y^0, y') = \varepsilon \tilde{v}_p(y^0) + \mathcal{O}(\varepsilon^2)$$

for some function $\tilde{v}_p \in C^{2,\alpha}([0, 1])$, it is then clear that

$$((L_\varepsilon - L_0) \bar{v}) = \varepsilon ((L_\varepsilon - L_0) \tilde{v}_p) + \mathcal{O}(\varepsilon^3)$$

where now the function \tilde{v}_p is considered as a function on the coordinates (y^0, y') by the simple relation $\tilde{v}_p(y^0, y') = \tilde{v}_p(y^0)$. Moreover if we consider the operator F restricted to functions \bar{v} that depend only on the first variable y^0 , it is clear that the linearization of F at $(p, \varepsilon, 0)$ maps from the subset of functions in $C_{m, NC}^{2,\alpha}$ that depend only on the first variable y^0 into the subset of functions in $C_{m, NC}^{1,\alpha}$ that depend only on the first variable y^0 . Then there exists a function $\delta_p \in C^{1,\alpha}([0, 1])$ such that

$$((L_\varepsilon - L_0) \tilde{v}_p)(y^0, y') = \varepsilon \delta_p(y^0) + \mathcal{O}(\varepsilon^2)$$

and then

$$((L_\varepsilon - L_0) \bar{v})(y^0, y') = \varepsilon^2 \delta_p(y^0) + \mathcal{O}(\varepsilon^3).$$

Now let us estimate the second term. Taking in account that $\bar{v} = \mathcal{O}(\varepsilon)$ we have

$$F(p, \varepsilon, \bar{v}) = F(p, \varepsilon, 0) + L_\varepsilon \bar{v} + L_\varepsilon^2(\bar{v}, \bar{v}) + \mathcal{O}(\varepsilon^3)$$

and then

$$F(p, \varepsilon, \bar{v}) - F(p, \varepsilon, 0) - L_\varepsilon \bar{v} = L_\varepsilon^2(\bar{v}, \bar{v}) + \mathcal{O}(\varepsilon^3).$$

If the function \bar{v} is such that

$$\bar{v}(y^0, y') = \varepsilon \tilde{v}_p(y^0) + \mathcal{O}(\varepsilon^2)$$

then

$$F(p, \varepsilon, \bar{v}) - F(p, \varepsilon, 0) - L_\varepsilon \bar{v} = \varepsilon^2 L_\varepsilon^2(\tilde{v}_p, \tilde{v}_p) + \mathcal{O}(\varepsilon^3)$$

where again the function \tilde{v}_p is considered as a function on the coordinates (y^0, y') by $\tilde{v}(y^0, y') = \tilde{v}(y^0)$, and as for L_ε it is easy to see that L_ε^2 maps from the subset of functions in $C_{m, NC}^{2,\alpha}$ that depend only on the first variable

y^0 into the subset of functions in $C_{m,NC}^{1,\alpha}$ that depend only on the first variable y^0 . Then there exists a function $\sigma_p \in C^{1,\alpha}([0,1])$ such that

$$F(p, \varepsilon, \bar{v}) - F(p, \varepsilon, 0) - L_\varepsilon \bar{v} = \varepsilon^2 \sigma_p(y^0) + \mathcal{O}(\varepsilon^3).$$

This completes the proof of the Lemma. \square

We are now able to conclude the proof of Proposition 4.5. Using Lemma 4.7 we get

$$\int_{S_+^n} (F(p, \varepsilon, \bar{v}) - F(p, \varepsilon, 0) - L_\varepsilon \bar{v}) \langle b, \cdot \rangle + \int_{S_+^n} (L_\varepsilon - L_0) \bar{v} \langle b, \cdot \rangle = \mathcal{O}(\varepsilon^3).$$

Then, from (43) and using Lemma 4.6, we have that for all $p \in \partial M$ and all $b \in \partial \mathbb{R}_+^{n+1}$ with $|b| = 1$ the following ε -expansion holds:

$$\int_{S_+^n} F(p, \varepsilon, \bar{v}(p, \varepsilon)) \langle b, \cdot \rangle = C \varepsilon^2 \tilde{g}(\nabla^{\tilde{g}} H(p), \Theta(b')) + \mathcal{O}(\varepsilon^3).$$

This completes the proof of the Proposition. \square

4.5 Proof of Theorem 1.3

Let $b = (0, b') \in \partial \mathbb{R}_+^{n+1}$ with $|b| = 1$ and define

$$G_b(p, \varepsilon) := \varepsilon^{-2} \int_{S_+^n} F(p, \varepsilon, \bar{v}(p, \varepsilon)) \langle b, \cdot \rangle = C \tilde{g}(\nabla^{\tilde{g}} H(p), \Theta(b')) + \mathcal{O}(\varepsilon).$$

Clearly if $\varepsilon \neq 0$, we have that

$$\int_{S_+^n} F(p, \varepsilon, \bar{v}(p, \varepsilon)) \langle b, \cdot \rangle = 0 \iff G_b(p, \varepsilon) = 0.$$

G_b is a function defined on $\partial M \times [0, +\infty)$ into \mathbb{R} . By the assumption of our main Theorem 1.3, ∂M has a nondegenerate critical point p_0 of the mean curvature. Then the differential of G_b with respect to p computed at $(p_0, 0)$ is invertible and $G_b(p_0, 0) = 0$. By the implicit function theorem, for all ε small enough there exists $p_\varepsilon \in \partial M$ close to p_0 such that

$$G_b(p_\varepsilon, \varepsilon) = 0$$

for all $b \in \partial \mathbb{R}_+^{n+1}$ with $|b| = 1$. In addition we have

$$\text{dist}(p_0, p_\varepsilon) \leq c \varepsilon$$

We conclude then that

$$F(p_\varepsilon, \varepsilon, \bar{v}(p_\varepsilon, \varepsilon)) \in K^\perp$$

where K is the kernel of the operator L_0 . But by the construction of \bar{v} , we have also that

$$F(p_\varepsilon, \varepsilon, \bar{v}(p_\varepsilon, \varepsilon)) \in K$$

and then

$$F(p_\varepsilon, \varepsilon, \bar{v}(p_\varepsilon, \varepsilon)) = 0.$$

This means that the normal derivative of the first eigenfunction of the Laplace-Beltrami operator on $\Omega_\varepsilon = B_{g,\varepsilon}^+(p_\varepsilon)$ with mixed boundary condition is constant on $\partial \Omega_\varepsilon \cap \mathring{M}$ and then Ω_ε is extremal.

The only remaining point in the proof of Theorem 1.3, is the analyticity of $\partial \Omega_\varepsilon \cap \mathring{M}$ when M itself is analytic. This is a classical consequence of the extremality condition, see [23].

5 Appendix

5.1 Expansion of the metric

Take the local coordinates x^0, x^1, \dots, x^n in a neighborhood of a point $p \in \partial M$ that we introduced in (4). We denote the corresponding coordinate vector fields by

$$X_j := \Psi_*(\partial_{x^j})$$

for $j = 0, 1, \dots, n$. We want to write the expansion of the coefficients g_{ij} of the metric Ψ^*g in these coordinates. According with our notation, E_j are the coordinate vector field X_j evaluated at p .

Proposition 5.1. *At the point of coordinate $x = (x^0, x^1, \dots, x^n)$, the following expansion holds :*

$$\begin{aligned} g_{00} &= 1 \\ g_{0j} &= 0 \\ g_{ij} &= \delta_{ij} + 2g(\nabla_{E_i} N, E_j) x^0 + R_{0i0j} (x^0)^2 + g(\nabla_{E_i} N, \nabla_{E_j} N) (x^0)^2 \\ &\quad + 2 \sum_k R_{k0ij} x^k x^0 + \frac{1}{3} \sum_{k,\ell} \tilde{R}_{ijkl} x^k x^\ell + \mathcal{O}(|x|^3) \end{aligned}$$

for $i, j, k, l = 1, \dots, n$, where

$$\begin{aligned} R_{0i0j} &= g(R(N, E_i) N, E_j) \\ R_{k0ij} &= g(R(E_k, N) E_i, E_j) \\ \tilde{R}_{ijkl} &= \tilde{g}(\tilde{R}(E_i, E_k) E_j, E_\ell). \end{aligned}$$

Here R and \tilde{R} are respectively the curvature tensors of M and ∂M .

This result of this proposition is very well known. For example, the same kind of coordinates that we use in this paper are also used in [30], and Proposition 5.1 of [30] combined with the classical expansion of a metric in its geodesic normal coordinate (see for example [34]) immediately implies our Proposition 5.1. Nevertheless, in order to make the reading easier, we write the proof of the proposition.

Proof. We consider the mapping F . The curve $x^0 \mapsto F(x^0, x)$ being a geodesic we have $g(X_0, X_0) \equiv 1$. This also implies that $\nabla_{X_0} X_0 \equiv 0$ and hence we get

$$\partial_{x^0} g(X_0, X_j) = g(\nabla_{X_0} X_0, X_j) + g(\nabla_{X_0} X_j, X_0) = g(\nabla_{X_0} X_j, X_0).$$

The vector fields X_0 and X_j being coordinate vector fields we have $\nabla_{X_0} X_j = \nabla_{X_j} X_0$ and we conclude that

$$2 \partial_{x^0} g(X_0, X_j) = 2 g(\nabla_{X_j} X_0, X_0) = \partial_{x^j} g(X_0, X_0) = 0.$$

Therefore, $g(X_0, X_j)$ does not depend on x^0 and since on ∂M this quantity is 0 for $j = 1, \dots, n$, we conclude that the metric g can be written as

$$g = d(x^0)^2 + \bar{g}_{x^0},$$

where \bar{g}_{x^0} is a family of metrics on ∂M smoothly depending on x^0 (this is nothing but Gauss' Lemma). If \tilde{g} is the metric of ∂M induced by g , we certainly have

$$\bar{g}_{x^0} = \tilde{g} + \mathcal{O}(x^0).$$

We now derive the next term the expansion of \bar{g}_{x^0} in powers of x^0 . To this aim, we compute

$$\partial_{x^0} g(X_i, X_j) = g(\nabla_{X_i} X_0, X_j) + g(\nabla_{X_j} X_0, X_i),$$

for all $i, j = 1, \dots, n$. Since $X_0 = N$ on ∂M , we get

$$\partial_{x^0} \bar{g}_{x^0} |_{x^0=0} = 2 g(\nabla N, \cdot),$$

by definition of the second fundamental form. This already implies that

$$\bar{g}_{x^0} = \tilde{g} + 2g(\nabla \cdot N, \cdot) x^0 + \mathcal{O}((x^0)^2).$$

Using the fact that the X_0 and X_j are coordinate vector fields, we can compute

$$\partial_{x^0}^2 g(X_i, X_j) = g(\nabla_{X_0} \nabla_{X_i} X_0, X_j) + g(\nabla_{X_0} \nabla_{X_j} X_0, X_i) + 2g(\nabla_{X_i} X_0, \nabla_{X_j} X_0). \quad (46)$$

By definition of the curvature tensor, we can write

$$\nabla_{X_0} \nabla_{X_j} = R(X_0, X_j) + \nabla_{X_j} \nabla_{X_0} + \nabla_{[X_0, X_j]},$$

which, using the fact that X_0 and X_j are coordinate vector fields, simplifies into

$$\nabla_{X_0} \nabla_{X_j} = R(X_0, X_j) + \nabla_{X_j} \nabla_{X_0}.$$

Since $\nabla_{X_0} X_0 \equiv 0$, we get

$$\nabla_{X_0} \nabla_{X_j} X_0 = R(X_0, X_j) X_0.$$

Inserting this into (46) yields

$$\partial_{x^0}^2 g(X_i, X_j) = 2g(R(X_0, X_i) X_0, X_j) + 2g(\nabla_{X_i} X_0, \nabla_{X_j} X_0).$$

Evaluation at $x^0 = 0$ gives

$$\partial_{x^0}^2 \bar{g}_{x^0} |_{x^0=0} = 2g(R(N, \cdot) N, \cdot) + 2g(\nabla \cdot N, \nabla \cdot N).$$

This implies that

$$\bar{g}_{x^0} = \tilde{g} + 2g(\nabla \cdot N, \cdot) x^0 + [g(\nabla \cdot N, \nabla \cdot N) + g(R(N, \cdot) N, \cdot)] (x^0)^2 + \mathcal{O}((x^0)^3) \quad (47)$$

Now that we have the first terms of the expansion of \bar{g}_{x^0} in powers of x^0 we find the expansion of these term with respect to the geodesic coordinates (x^1, \dots, x^n) of ∂M in a neighborhood of p . Recall that for $i, j, k, l = 1, \dots, n$

$$\tilde{g}_{ij} = \delta_{ij} + \frac{1}{3} \sum_{k, \ell} \tilde{R}_{ikj\ell} x^k x^\ell + \mathcal{O}(|x|^3), \quad (48)$$

where

$$\tilde{R}_{ikj\ell} = \tilde{g}(\tilde{R}(E_i, E_k) E_j, E_\ell)$$

The proof of this fact can be found for example in [34]. Moreover for $k = 1, \dots, n$ we have

$$\begin{aligned} \partial_{x^k} g(\nabla_{X_i} N, X_j) &= g(\nabla_{X_k} \nabla_{X_i} N, X_j) + g(\nabla_{X_i} N, \nabla_{X_k} X_j) \\ &= g(\nabla_{X_k} \nabla_N X_i, X_j) + g(\nabla_{X_i} N, \nabla_{X_k} X_j) \\ &= g(R(X_k, N) X_i, X_j) + g(\nabla_N \nabla_{X_k} X_i, X_j) + g(\nabla_{X_i} N, \nabla_{X_k} X_j) \end{aligned}$$

and evaluated at p

$$\partial_{x^k} g(\nabla_{X_i} N, X_j)|_p = g(R(E_k, N) E_i, E_j) \quad (49)$$

From (47), using (48) and (49), we find the expansion of the metric in the coordinates x^0, x^1, \dots, x^n up to the term of order $|x|^2$. \square

5.2 Technical Lemmas

Lemma 5.2. *For all $\sigma = 1, \dots, n$, we have*

$$\sum_{i,j,k} \int_{S_+^n} R_{k0ij} x^0 x^i x^j x^k x^\sigma = 0.$$

Proof. To see that we consider all terms of the above sum, obtained fixing the 4-tuple (i, k, j, σ) . We observe that if in such a 4-tuple there is an element that appears an odd number of time then $\int_{S_+^n} x^0 x^i x^j x^k x^\sigma = 0$. Then

$$\sum_{i,j,k} \int_{S_+^n} R_{k0ij} x^0 x^i x^j x^k x^\sigma = \sum_i \int_{S_+^n} (R_{\sigma 0ii} + R_{i0i\sigma} + R_{i0\sigma i}) x^0 (x^i)^2 (x^\sigma)^2 = 0$$

by the symmetries of the curvature tensor. \square

Lemma 5.3. *For all $\sigma = 1, \dots, n$, we have*

$$\sum_{i,j} \int_{S_+^n} R_{i0ij} x^0 x^j x^\sigma = - \left(\int_{S_+^n} x^0 (x^1)^2 \right) H_\sigma$$

Proof. Again, we find that $\int_{S_+^n} x^0 x^j x^\sigma d\text{vol}_{\tilde{g}} = 0$ unless the indices j, σ are equal. Hence

$$\sum_{i,j} \int_{S_+^n} R_{i0ij} x^0 x^j x^\sigma = \left(\int_{S_+^n} x^0 (x^\sigma)^2 \right) \sum_i R_{i0i\sigma} = - \left(\int_{S_+^n} x^0 (x^1)^2 \right) H_\sigma$$

This completes the proof of the result. \square

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