# Geometric quantization on CR manifolds 

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Let $X$ be a compact connected orientable Cauchy-Riemann (CR) manifold with the action of a compact Lie group $G$. Under natural pseudoconvexity assumptions we show that the CR Guillemin-Sternberg map is an isomorphism at the level of Sobolev spaces of CR functions, modulo a finite-dimensional subspace. As application we study this map for holomorphic line bundles which are positive near the inverse image of 0 by the momentum map. We also show that "quantization commutes with reduction" for Sasakian manifolds.

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## 1. Introduction and Statement of the Main Results

The famous geometric quantization conjecture of Guillemin and Sternberg 19 states that for a compact pre-quantizable symplectic manifold admitting a Hamiltonian action of a compact Lie group, the principle of "quantization commutes with reduction" holds. This conjecture was first proved independently by Meinrenken [43] and Vergne [53] for the case where the Lie group is abelian, and by Meinrenken 44] in the general case, then Tian-Zhang [52] gave a purely analytic proof in general case with various generalizations, see [54] for a survey and complete references on this subject. In the case of a non-compact symplectic manifold $M$ with a compact Lie group action $G$, this question was solved by Ma-Zhang [39, 40] as a solution to a conjecture of Vergne in her ICM 2006 plenary lecture [55], see [36] for a survey. Paradan [47] gave a new proof, cf. also the recent work [24]. A natural choice for the quantum spaces of a compact symplectic manifold is the kernel of the Dirac operator.

In [38, Ma-Zhang established the asymptotic expansion of the $G$-invariant Bergman kernel for a positive line bundle $L$ over a compact symplectic manifold $M$ and by using the asymptotic expansion of $G$-invariant Bergman kernel, they could establish the "quantization commutes with reduction" theorem when the power of the line bundle $L$ is high enough.

On a compact Kähler manifold $M$ endowed with a prequantum line bundle $L$, a natural choice for the Hilbert space of quantum states is the space $H^{0}\left(M, L^{m}\right)$ of holomorphic sections of the tensor powers $L^{m}$. The family of quantum spaces $H^{0}\left(M, L^{m}\right)$ indexed by $m \in \mathbb{N}$ plays an essential role in geometric quantization and the semi-classical limit $m \rightarrow \infty$ allows to recover the classical mechanics of the phase space $M$.

One can wrap up the family of spaces $H^{0}\left(M, L^{m}\right), m \in \mathbb{N}$, as subspaces of a single Hilbert space by considering the $S^{1}$-bundle $X \subset L^{*}$, which is a strictly pseudoconvex CR manifold and identifying $H^{0}\left(M, L^{m}\right)$ with the $S^{1}$ isotypes of $m$-equivariant CR functions on $X$. The Hilbert space sum $\widehat{\oplus}_{m \in \mathbb{N}} H^{0}\left(M, L^{m}\right)$ can be identified to the space $H_{b}^{0}(X)$ of $L^{2} \mathrm{CR}$ functions on $X$ and the sum of the Bergman projections $B_{m}$ on $H^{0}\left(M, L^{m}\right)$ can be identified to the Szegő projector $S$ on $H_{b}^{0}(X)$. A fundamental fact is that the asymptotic behavior of $B_{m}$ is encoded in the singularities of the Szegő kernel $S(\cdot, \cdot)$. We can thus think of $X$ as the quantizing principal bundle of $M$ and of the space of $L^{2} \mathrm{CR}$ functions as the quantum space of $X$. In the presence of a $G$-action on $M$, which lifts to an action on $L$, we have an induced $G$-action on $X$ and on $H_{b}^{0}(X)$.

The quantization of strictly pseudoconvex or more generally contact manifolds via the Szegő projector or its generalizations was developed by Boutet de Monvel and Guillemin [7] and can be applied to the Kähler quantization by using the above construction. In this paper, we study the quantization of CR manifolds and the principle of "quantization commutes with reduction". For this purpose we develop a $G$-invariant Fourier integral operator calculus which will be used to obtain the
asymptotics of the $G$-invariant Szegő kernel. Our Theorem 1.2 is new even for compact strictly pseudoconvex CR manifolds and our results pertain also to Sasakian manifolds (see Theorem 1.3). Sasakian geometry is an important odd-dimensional counterpart of Kähler geometry. It is known that irregular Sasakian manifolds admit a compact CR torus action (see [23, Sec. 3] or Remark [2.2) and the study of $G$ equivariant CR functions on a Sasakian manifold is important in Sasaki geometry. We refer to [1, 9, 16, 17, 57] for the fundamentals of contact and Sasakian reduction and examples.

An important difference between the CR setting and the Kähler/symplectic setting is that the quantum spaces in the case of a compact Kähler/symplectic manifolds are finite-dimensional, whereas for the compact strictly pseudoconvex CR manifolds that we consider the quantum spaces consist of CR functions and are infinite-dimensional. In [20, Theorem 3.6], Guillemin-Sternberg proposed a version of quantization for symplectic cones associated with a homogeneous moment map. Again, in contrast to our Theorem 1.2 their invariant spaces are finite-dimensional.

We now formulate the main results. We refer to Sec. 2 for some notations and terminology used here. Let $\left(X, T^{1,0} X\right)$ be a compact orientable CR manifold of dimension $2 n+1, n \geq 1$, where $T^{1,0} X$ denotes the CR structure of $X$. Let $H X \subset$ $T X$ be the associated Levi distribution with complex structure $J \in \operatorname{End}(H X)$ and let $\omega_{0} \in \mathscr{C}^{\infty}\left(X, T^{*} X\right)$ be a non-vanishing real 1-form annihilating $H X$, called characteristic 1-form.

Let $G$ be a $d$-dimensional compact Lie group with Lie algebra $\mathfrak{g}$ acting on $X$ by preserving $J$ and $\omega_{0}$. Let $\mu: X \rightarrow \mathfrak{g}^{*}$ be the associated moment map $\mu: X \rightarrow \mathfrak{g}^{*}$ (cf. (2.27)). We will mainly work in the following setting.

Assumption 1.1. The $G$-action preserves the complex structure $J$ on $H X$ and the characteristic 1-form $\omega_{0}$, it is free on $\mu^{-1}(0)$, and one of the following conditions are fulfilled:
(i) $\operatorname{dim} X \geq 5$ and the Levi form of $X$ is positive definite near $\mu^{-1}(0)$.
(ii) $\operatorname{dim} X=3$, the Levi form of $X$ is positive definite everywhere and $\bar{\partial}_{b}$ has closed range in $L^{2}$ on $X$.

Due to Lemma 2.5. Assumption 1.1]implies that 0 is a regular value of $\mu$, hence $\mu^{-1}(0)$ is a $d$-codimensional submanifold of $X$. Let

$$
\begin{equation*}
Y:=\mu^{-1}(0), \quad X_{G}:=\mu^{-1}(0) / G \tag{1.1}
\end{equation*}
$$

The space $X_{G}$ is called the CR reduction. Under our hypotheses, if $\operatorname{dim} X_{G} \geq 3$, $X_{G}$ is a strictly pseudoconvex CR manifold with characteristic 1-form (in this case also contact form) $\omega_{0, G}$ induced canonically by $\omega_{0}$, see (2.38). If $\operatorname{dim} X_{G}=1$, then each of the finitely many components of $X_{G}$ is diffeomorphic to a circle.

Let $\bar{\partial}_{b}: \mathscr{C}^{\infty}(X) \rightarrow \Omega^{0,1}(X)$ and if $\operatorname{dim} X_{G} \geq 3$, let $\bar{\partial}_{b, X_{G}}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow \Omega^{0,1}\left(X_{G}\right)$ be the tangential Cauchy-Riemann operators on $X$ and $X_{G}$, respectively. We extend $\bar{\partial}_{b}$ and $\bar{\partial}_{b, X_{G}}$ to $L^{2}$ spaces by taking their weak maximal extension, see (2.13). We consider the spaces of $L^{2} \mathrm{CR}$ functions

$$
\begin{equation*}
H_{b}^{0}(X):=\left\{u \in L^{2}(X): \bar{\partial}_{b} u=0\right\}, \quad H_{b}^{0}\left(X_{G}\right):=\left\{u \in L^{2}\left(X_{G}\right): \bar{\partial}_{b, X_{G}} u=0\right\} . \tag{1.2}
\end{equation*}
$$

If $\operatorname{dim} X_{G}=1$, so $X_{G}$ is a finite union of circles, we set $H_{b}^{0}\left(X_{G}\right)$ to be the direct sum of the Hardy spaces of the components, that is, the $L^{2}$ subspaces of functions with vanishing Fourier coefficients of negative degree. The common feature of the spaces $H_{b}^{0}\left(X_{G}\right)$ for $\operatorname{dim} X_{G} \geq 3$ and $\operatorname{dim} X_{G}=1$ is the fact that they are boundary values of holomorphic functions in a filling of $X_{G}$ by a complex manifold (see Sec. 2.4).

Then $H_{b}^{0}(X)$ is a (possible infinite-dimensional) $G$-representation, its $G$ invariant part is the $G$-invariant $L^{2} \mathrm{CR}$ functions on $X$,

$$
\begin{equation*}
H_{b}^{0}(X)^{G}:=\left\{u \in H_{b}^{0}(X): h^{*} u=u, \text { for any } h \in G\right\} . \tag{1.3}
\end{equation*}
$$

For every $s \in \mathbb{R}$ let $H^{s}(X)$ and $H^{s}\left(X_{G}\right)$ denote the Sobolev spaces of $X$ and $X_{G}$ of order $s$ and let $(\cdot, \cdot)_{s}$ and $(\cdot, \cdot)_{X_{G}, s}$ be the inner products on $H^{s}(X)$ and $H^{s}\left(X_{G}\right)$, respectively (see (5.3)). For every $s \in \mathbb{R}$ put

$$
\begin{equation*}
H_{b}^{0}(X)_{s}:=\left\{u \in H^{s}(X): \bar{\partial}_{b} u=0 \text { in the sense of distributions }\right\} . \tag{1.4}
\end{equation*}
$$

We define $H_{b}^{0}\left(X_{G}\right)_{s}$ and $H_{b}^{0}(X)_{s}^{G}$ in the same way. If $\operatorname{dim} X_{G}=1$, we set $H_{b}^{0}\left(X_{G}\right)_{s}$ to be the direct sum of the Hardy-Sobolev spaces of the components, that is, the subspaces of $H^{s}\left(S^{1}\right)$ of distributions with vanishing Fourier coefficients of negative degree.

Let $\iota: Y \rightarrow X$ be the natural inclusion and let $\iota^{*}: \mathscr{C}^{\infty}(X) \rightarrow \mathscr{C}^{\infty}(Y)$ be the pull-back by $\iota$. Let $\iota_{G}: \mathscr{C}^{\infty}(Y)^{G} \rightarrow \mathscr{C}^{\infty}\left(X_{G}\right)$ be the natural identification. Let

$$
\begin{equation*}
\sigma_{G}: H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)^{G} \rightarrow H_{b}^{0}\left(X_{G}\right), \quad \sigma_{G}=\iota_{G} \circ \iota^{*} \tag{1.5}
\end{equation*}
$$

The map (1.5) is well defined, see the construction of the CR reduction in Sec. 2.3, The map $\sigma_{G}$ does not extend to a bounded operator on $L^{2}$, so it is necessary to consider its extension to Sobolev spaces. From Theorem 5.3, $\sigma_{G}$ extends by density to a bounded operator

$$
\begin{equation*}
\sigma_{G}=\sigma_{G, s}: H_{b}^{0}(X)_{s}^{G} \rightarrow H_{b}^{0}\left(X_{G}\right)_{s-\frac{d}{4}}, \quad \text { for every } s \in \mathbb{R} \tag{1.6}
\end{equation*}
$$

This operator can be thought as a Guillemin-Sternberg map in the CR setting. It maps the "first quantize and then reduce" space (the space of $G$-invariant Sobolev CR functions on $X$ ) to the "first reduce and then quantize" space (the space of Sobolev CR functions on $X_{G}$ ). Indeed, from the point of view of quantum mechanics, the Hilbert space structures play an essential role. It is natural, then, to investigate the extent to which the CR Guillemin-Sternberg map is Fredholm. The following main result of this work gives a CR analogue of the fact that
the canonical reduction map in the symplectic context is an isomorphism for sufficiently large tensor powers of the prequantum line bundle (cf. [38, Theorem 0.9] and Theorem (1.4).

Theorem 1.2. Let $X$ be a compact orientable $C R$ manifold and let $G$ be a compact Lie group acting on $X$ such that the $G$-action preserves $J$ and $\omega_{0}$ and Assumption 1.1 holds. Suppose that $\bar{\partial}_{b, X_{G}}$ has closed range in $L^{2}$. Then, for every $s \in \mathbb{R}$, the $C R$ Guillemin-Sternberg map (1.6) is Fredholm. Actually, $\operatorname{Ker} \sigma_{G, s}$ and $\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$ are finite-dimensional subspaces of $\mathscr{C}^{\infty}(X) \cap H_{b}^{0}(X)^{G}$ and $\mathscr{C}^{\infty}\left(X_{G}\right) \cap H_{b}^{0}\left(X_{G}\right)$, respectively, $\operatorname{Ker} \sigma_{G, s}$ and the index $\operatorname{dim} \operatorname{Ker} \sigma_{G, s}-\operatorname{dim}\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$ are independent of $s$.

We note that $\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$ is given by
$\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}:=\left\{u \in H_{b}^{0}\left(X_{G}\right)_{s-\frac{d}{4}}:\left(\sigma_{G, s} v, u\right)_{X_{G}, s-\frac{d}{4}}=0\right.$, for every $\left.v \in H_{b}^{0}(X)_{s}^{G}\right\}$.

Under Assumption 1.1(i) the hypothesis that $\operatorname{dim} X \geq 5$ is used in order to have local subelliptic Sobolev estimates on the set where the Levi form is positive definite (Theorem 3.10) and leads to the fact that the $G$-invariant Kohn Laplacian (3.26) has closed range in $L^{2}$. Note also that the Kohn Laplacian on strictly pseudoconvex CR manifolds of dimension greater than or equal to five has always closed range in $L^{2}$ but this is not true for all three-dimensional strictly pseudoconvex CR manifolds (a detailed discussion about the closed range of $\bar{\partial}_{b}$ in $L^{2}$ can be found in Sec. (2.4). In the case when $\operatorname{dim} X=3$ we will state in Theorem 5.7 a version of Theorem 1.2 under weaker hypotheses as Assumption 1.1(ii), namely that $X$ is pseudoconvex of finite type and $\bar{\partial}_{b, X}$ has closed range in $L^{2}$.

We turn now our attention to Sasakian manifolds. Let $\left(X, T^{1,0} X\right)$ be a compact connected Sasakian manifold (see Sec. 2.2), i.e. $\left(X, T^{1,0} X\right)$ is a compact connected strictly pseudoconvex CR manifold and we can fix a contact form $\omega_{0}$ and a Reeb vector field $R$ such that $i_{R} \omega_{0}=1, i_{R} d \omega_{0}=0$ and the flow associated with $R$ preserves $T^{1,0} X$ (cf. Remarks 2.2 and 2.3). Assume that the action of the compact Lie group $G$ on $X$ verifies Assumption 1.1. Moreover, we assume that

$$
\begin{equation*}
R \text { is } G \text {-invariant. } \tag{1.8}
\end{equation*}
$$

By Remark 2.2, the flow associated with $R$ preserves $H X, J$ and the natural metric $g_{\omega_{0}}$ on $T X$, in particular, $R$ is a Killing vector field. Since $X$ is compact, this implies the flow associated with $R$ generates a compact torus $\mathbb{T}$-action on $X$ and this $\mathbb{T}$ action commutes with the $G$-action. Thus it naturally induces a $\mathbb{T}$-action on $X_{G}$ and the generator $R$ induces the Reeb vector field $\widehat{R}$ on $X_{G}$. Since the flow associated with $R$ preserves the CR structure $T^{1,0} X$ and commutes with $G$, it follows that the flow associated with $\widehat{R}$ preserves the CR structure $T^{1,0} X_{G}$. Thus, $\left(X_{G}, T^{1,0} X_{G}\right)$ is a strictly pseudoconvex CR manifold with a CR Reeb vector field $\widehat{R}$. By (2.17) and Remark 2.3 $X_{G}$ is also a compact Sasakian manifold.

Now $H_{b}^{0}(X)^{G}$ and $H_{b}^{0}\left(X_{G}\right)$ are both $\mathbb{T}$-Hilbert spaces, thus we have the decomposition of Hilbert spaces via the weight $\alpha \in \mathbb{T}^{*}\left(\simeq \mathbb{Z}^{\operatorname{dim} \mathbb{T}}\right)$ of $\mathbb{T}$-action:

$$
\begin{equation*}
H_{b}^{0}(X)^{G}=\oplus_{\alpha \in \mathbb{T}^{*}} H_{b, \alpha}^{0}(X)^{G}, \quad H_{b}^{0}\left(X_{G}\right)=\oplus_{\alpha \in \mathbb{T}^{*}} H_{b, \alpha}^{0}\left(X_{G}\right) \tag{1.9}
\end{equation*}
$$

We refer to [23] for further applications of this weight decomposition. Both $H_{b, \alpha}^{0}(X)^{G}$ and $H_{b, \alpha}^{0}\left(X_{G}\right)$ are finite-dimensional subspaces of $\mathscr{C}^{\infty}(X)^{G}$ and $\mathscr{C}^{\infty}\left(X_{G}\right)$, respectively, as subspaces of the eigenspaces of the elliptic operators $\bar{\partial}_{b, X}^{*} \bar{\partial}_{b, X}-R^{2}, \bar{\partial}_{b, X}{ }_{G} \bar{\partial}_{b, X_{G}}-\widehat{R}^{2}$, respectively, of eigenvalues $|\alpha(R)|^{2}$. From the definition (1.5) of the map $\sigma_{G}$ we see that

$$
\begin{equation*}
\sigma_{G} R u=\widehat{R} \sigma_{G} u, \quad \text { for any } u \in H_{b}^{0}(X)^{G} \tag{1.10}
\end{equation*}
$$

and hence $\sigma_{G}$ maps $H_{b, \alpha}^{0}(X)^{G}$ to $H_{b, \alpha}^{0}\left(X_{G}\right)$. From this observation, Theorem 1.2 and the fact that $\bar{\partial}_{b}$ has closed range in $L^{2}$ on Sasakian manifolds (see 42, also Sec. (2.4), we deduce:

Theorem 1.3 (quantization commutes with reduction for Sasakian manifolds). Let $X$ be a compact connected Sasakian manifold with a CR Reeb vector field $R$. Suppose that $X$ admits a compact Lie group action $G$ which preserves the complex structure $J$ on $H X$ and the characteristic 1-form $\omega_{0}$, it is free on $\mu^{-1}(0)$, and the Reeb vector field $R$ is $G$-invariant. Then with the exception of finitely many $\alpha \in \mathbb{T}^{*}$ the map

$$
\begin{equation*}
\sigma_{G}: H_{b, \alpha}^{0}(X)^{G} \rightarrow H_{b, \alpha}^{0}\left(X_{G}\right) \tag{1.11}
\end{equation*}
$$

is an isomorphism.
We now apply Theorem 1.2 to the case of complex manifolds. Let ( $L, h^{L}$ ) be a Hermitian holomorphic line bundle over a connected compact complex manifold $(M, J)$ with $\operatorname{dim}_{\mathbb{C}} M=n$, where $J$ denotes the complex structure of $T M$ and $h^{L}$ is a Hermitian metric of $L$. We denote by $R^{L}$ the Chern curvature of $\left(L, h^{L}\right)$. We assume that $G$ acts holomorphically on $(M, J)$, and that the action lifts to a holomorphic action on $L$. We assume further that $h^{L}$ is preserved by the $G$-action. Then $R^{L}$ is a $G$-invariant form. Let $\tilde{\mu}: M \rightarrow \mathfrak{g}^{*}$ be the moment map defined by the Kostant formula (6.1).

Assume that $0 \in \mathfrak{g}^{*}$ is regular and the action of $G$ on $\tilde{\mu}^{-1}(0)$ is free. If $i R^{L}$ is positive near $\tilde{\mu}^{-1}(0)$, then the analogue of the Marsden-Weinstein reduction holds. More precisely, the complex structure $J$ on $M$ induces a complex structure $J_{G}$ on $M_{G}:=\tilde{\mu}^{-1}(0) / G$, for which the line bundle $L_{G}:=L / G$ is a holomorphic line bundle over $M_{G}$. For $m \in \mathbb{N}$, let $H^{0}\left(M, L^{m}\right)^{G}$ denote the space of $G$-invariant holomorphic sections with values in $L^{m}$.

Let $X:=\left\{v \in L^{*}:|v|_{h^{L^{*}}}^{2}=1\right\}$ be the circle bundle of $L^{*}$, where $h^{L^{*}}$ is the Hermitian metric on $L^{*}$ induced by $h^{L}$. Let $e^{i \theta}$ be the natural $S^{1}$-action on the
fibers of $X$. The action of $G$ on $M$ lifts to a CR action on $X$. For every $m \in \mathbb{N}$, put

$$
\begin{equation*}
H_{b, m}^{0}(X)^{G}:=\left\{u \in H_{b}^{0}(X)^{G}:\left(e^{i \theta}\right)^{*} u=e^{i m \theta} u \text { on } X, \text { for every } e^{i \theta} \in S^{1}\right\} \tag{1.12}
\end{equation*}
$$

It is easy to check that for every $m \in \mathbb{N}$ there are canonical isomorphisms

$$
\begin{equation*}
H_{b, m}^{0}(X)^{G} \cong H^{0}\left(M, L^{m}\right)^{G}, \quad H_{b, m}^{0}\left(X_{G}\right) \cong H^{0}\left(M_{G}, L_{G}^{m}\right) \tag{1.13}
\end{equation*}
$$

On account of (1.13), Theorem 1.2 (under Assumption 1.1 (i)) and the correspondence between the curvature $i R^{L}$ and the Levi form of $X$ (cf. Section 2.2) we deduce:

Theorem 1.4. Let $M$ be a compact connected complex manifold, $\operatorname{dim}_{\mathbb{C}} M \geq 2$, and $\left(L, h^{L}\right)$ be a Hermitian holomorphic line bundle over $M$. Let $G$ be a compact Lie group acting holomorphically on $M$ and whose action lifts to $\left(L, h^{L}\right)$. Suppose that $i R^{L}$ is positive near $\tilde{\mu}^{-1}(0)$ and $G$ acts freely on $\tilde{\mu}^{-1}(0)$. Then for $m$ large enough, the canonical map between $H^{0}\left(M, L^{m}\right)^{G}$ and $H^{0}\left(M_{G}, L_{G}^{m}\right)$ by restriction is an isomorphism, in particular

$$
\begin{equation*}
\operatorname{dim} H^{0}\left(M, L^{m}\right)^{G}=\operatorname{dim} H^{0}\left(M_{G}, L_{G}^{m}\right) \tag{1.14}
\end{equation*}
$$

If $L$ is positive on the whole $M$, this canonical isomorphism between $H^{0}\left(M, L^{m}\right)^{G}$ and $H^{0}\left(M_{G}, L_{G}^{m}\right)$ was constructed in [19, 58] for $m=1$ (see also for the metric aspect of this isomorphism [38, (0.27), Corollary 4.13] for $m$ large enough). This implies that the map $\sigma_{G}$ in Theorem 1.2 is actually an isomorphism in the case of the circle bundle $X$ of $L^{*}$. Thus we expect in many situations (in particular, if additionally $X$ is strictly pseudoconvex), that $\sigma_{G}$ in Theorem 1.2 is an isomorphism. Theorem 1.4 gives a version of the result of [19, 58] for large enough powers of $L$ by requiring the positivity of $i R^{L}$ only on $\tilde{\mu}^{-1}(0)$. In the case $\operatorname{dim}_{\mathbb{C}} M=1$, hence $\operatorname{dim} X=3$, Assumption 1.1(ii) corresponds to $L$ being positive everywhere on $M$, so an application of Theorem 1.2 does not bring anything new. We will give in Theorem 6.1] a version for almost complex manifolds of Theorem [1.4]

In the rest of Introduction we explain some technical aspects, in particular some results on the $G$-invariant Szegő projection, necessary to establish Theorem 1.2

We introduce a $G$-invariant Hermitian metric $g=g^{\mathbb{C} T X}$ on $X$ as in Lemma 2.7 which fixes the $L^{2}$ spaces on $X$. The $G$-invariant Szegő projection is the orthogonal projection

$$
\begin{equation*}
S_{G}: L^{2}(X) \rightarrow H_{b}^{0}(X)^{G} \tag{1.15}
\end{equation*}
$$

with respect to $(\cdot, \cdot)$. The $G$-invariant Szegö kernel $S_{G}(x, y) \in \mathscr{D}^{\prime}(X \times X)$ is the distribution kernel of $S_{G}$. In Theorem 3.25 we will prove that $S_{G}$ is a complex Fourier integral operator and in Theorem 3.27, we will show the regularity property

$$
\begin{equation*}
S_{G}: \mathscr{C}^{\infty}(X) \rightarrow H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X) \tag{1.16}
\end{equation*}
$$

From (1.16) we conclude that $H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)$ is dense in $H_{b}^{0}(X)^{G}$.

Let $S_{X_{G}}: L^{2}\left(X_{G}\right) \rightarrow H_{b}^{0}\left(X_{G}\right)$ be the orthogonal projection with respect to $(\cdot, \cdot)_{X_{G}}$ (cf. Convention 2.8). It follows from general results [8, 27] that if $\bar{\partial}_{b, X_{G}}$ has closed range in $L^{2}$, then $S_{X_{G}}$ is a Fourier integral operator with complex phase and also a pseudodifferential operator on $X_{G}$. In particular,

$$
\begin{equation*}
S_{X_{G}}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow H_{b}^{0}\left(X_{G}\right) \cap \mathscr{C}^{\infty}\left(X_{G}\right) \tag{1.17}
\end{equation*}
$$

which implies that $H_{b}^{0}\left(X_{G}\right) \cap \mathscr{C}^{\infty}\left(X_{G}\right)$ is dense in $H_{b}^{0}\left(X_{G}\right)$ if $\bar{\partial}_{b, X_{G}}$ has closed range in $L^{2}$.

We consider the linear map

$$
\begin{equation*}
\mathcal{R}_{x}: \underline{\mathfrak{g}}_{x} \rightarrow \underline{\mathfrak{g}}_{x}, \quad u \mapsto \mathcal{R}_{x} u, \quad\left\langle\mathcal{R}_{x} u, v\right\rangle=d \omega_{0}(u, J v)_{x} \tag{1.18}
\end{equation*}
$$

For $x \in Y$ we denote by $Y_{x}=\{h . x: h \in G\}$ the $G$-orbit of $Y$; then $Y_{x}$ is a $d$-dimensional submanifold of $X$. The $G$-invariant Hermitian metric $g$ induces a volume form $d v_{Y_{x}}$ on $Y_{x}$. Put

$$
\begin{equation*}
f_{G}(x)=\left|\operatorname{det} \mathcal{R}_{x}\right|^{-\frac{1}{4}} \sqrt{V_{\mathrm{eff}}(x)} \in \mathscr{C}^{\infty}(Y)^{G} \quad \text { with } V_{\mathrm{eff}}(x):=\int_{Y_{x}} d v_{Y_{x}} \tag{1.19}
\end{equation*}
$$

where $\mathscr{C}^{\infty}(Y)^{G}$ denotes the space of $G$-invariant smooth functions on $Y=\mu^{-1}(0)$.
Let $E: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow \mathscr{C}^{\infty}\left(X_{G}\right)$ be a classical elliptic pseudodifferential operator with principal symbol $p_{E}(x, \xi)=|\xi|^{-d / 4}$. Let

$$
\begin{equation*}
\sigma: H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)^{G} \rightarrow H_{b}^{0}\left(X_{G}\right), \quad \sigma=S_{X_{G}} \circ E \circ \iota_{G} \circ f_{G} \circ \iota^{*} \tag{1.20}
\end{equation*}
$$

It turns out that the operator $\sigma$ in (1.20) is bounded, see Corollary 4.16 and thus extends by density to a bounded operator

$$
\begin{equation*}
\sigma: H_{b}^{0}(X)^{G} \rightarrow H_{b}^{0}\left(X_{G}\right) \tag{1.21}
\end{equation*}
$$

We have encoded in the definition (1.20) some corrections in order to obtain good analytic properties of $\sigma$. One correction is the multiplication with the function $f_{G}$ from (1.19); this reflects the need to reconcile the volume forms on $\mu^{-1}(0)$ and on $X_{G}$. The multiplication by $f_{G}$ changes the CR character of the result, therefore the need to project back to the CR space by $S_{X_{G}}$. Here comes the role of $E$, which is more subtle (see also Remark (1.6). Ideally, the map $\sigma$ should be unitary. But we can content ourselves to require that $\sigma^{*} \sigma$ is "microlocally close" to $S_{G}$, where $\sigma^{*}: H_{b}^{0}\left(X_{G}\right) \rightarrow \mathscr{D}^{\prime}(X)$ is the adjoint of $\sigma$. In other words, we want $\sigma^{*} \sigma$ to be a complex Fourier integral operator with the same phase, the same order and the same leading symbol as $S_{G}$. To achieve this, we need to take $E$ to be a classical elliptic pseudodifferential operator with principal symbol $p_{E}(x, \xi)=|\xi|^{-d / 4}$.

The main technical result of this work is the following.
Theorem 1.5. Under the assumption of Theorem 1.2, the map $\sigma$ is Fredholm. Actually, $\operatorname{Ker} \sigma$ and $(\operatorname{Im} \sigma)^{\perp}$ are finite-dimensional subspaces of $\mathscr{C}^{\infty}(X) \cap H_{b}^{0}(X)^{G}$ and $\mathscr{C}^{\infty}\left(X_{G}\right) \cap H_{b}^{0}\left(X_{G}\right)$, respectively.

Remark 1.6. Note that the definition of $\sigma$ depends on the choice of the elliptic pseudodifferential operator $E$. We actually show that for any classical elliptic pseudodifferential operator $E$ with the same principal symbol $p_{E}(x, \xi)=|\xi|^{-d / 4}$, the map $\sigma: H_{b}^{0}(X)^{G} \rightarrow H_{b}^{0}\left(X_{G}\right)$ is Fredholm. Up to lower order terms of $E$, the map $\sigma$ is a canonical choice. The elliptic pseudodifferential operator $E$ corresponds to the power $m^{-d / 4}$ in the isomorphism map between $H^{0}\left(M, L^{m}\right)^{G}$ and $H^{0}\left(M_{G}, L_{G}^{m}\right)$ in complex case. Here we use the same notations as in the discussion after Theorem 1.3, More precisely, Ma-Zhang [38, Theorem 0.10] showed that the map

$$
\begin{equation*}
\sigma_{m}: H^{0}\left(M, L^{m}\right)^{G} \rightarrow H^{0}\left(M_{G}, L_{G}^{m}\right), \quad \sigma_{m}=m^{-d / 4} B_{M_{G}}^{m} \circ \iota_{G} \circ f_{G} \circ \iota^{*}, \tag{1.22}
\end{equation*}
$$

is an asymptotic isometry if $m$ is large enough, where $\iota_{G}, \iota^{*}$ and $f_{G} \in \mathscr{C}^{\infty}(M)^{G}$ are defined as in the discussion before (1.20) and $B_{M_{G}}^{m}: L^{2}\left(M_{G}, L_{G}^{m}\right) \rightarrow H^{0}\left(M_{G}, L_{G}^{m}\right)$ is the orthogonal projection. When we change $m^{-d / 4}$ in (1.22) to any $m$-depend function with order $m^{-d / 4}+O\left(m^{-\frac{d}{4}-1}\right)$, we still have an isomorphism between $H^{0}\left(M, L^{m}\right)^{G}$ and $H^{0}\left(M_{G}, L_{G}^{m}\right)$ for $m$ large. Moreover, in view of Theorem 1.2 to get $L^{2}$ isomorphism it makes sense to take an elliptic pseudodifferential operator $E$ of order $-\frac{d}{4}$.

Remark 1.7. (i) In this work, we do not assume that $\bar{\partial}_{b}$ has closed range in $L^{2}$ on $X$. We will show in Sec. 3.2 that under the assumption that the Levi form is positive on $Y=\mu^{-1}(0)$, the $G$-invariant Kohn Laplacian has closed range in $L^{2}$ and this is enough to obtain a full asymptotic expansion for the $G$ invariant Szegő kernel $S_{G}(x, y)$ (see Theorem 3.25). In order to show that the $G$-invariant Kohn Laplacian has closed range in $L^{2}$ we need the hypothesis that the dimension of $X$ is greater than or equal to five.
(ii) The asymptotic expansion for $S_{G}$ is also a new result. In [28], Hsiao and Huang obtained an asymptotic expansion for $S_{G}$ under the assumption that $\bar{\partial}_{b}$ has closed range in $L^{2}$ on $X$. In [28, Hsiao and Huang established "quantization commutes with reduction" results for CR manifolds with $S^{1}$-action. The spaces considered in 28 are finite-dimensional. To handle the infinite-dimensional case, we need to develop a new kind of calculus of complex Fourier integral operators.
(iii) In [19, Theorem A6], Guillemin and Sternberg showed that the underlying canonical relation of $S_{G}$ is given by $\left\{(h x, x): x \in \mu^{-1}(0), h \in G\right\}$ under the assumption that $X$ is strictly pseudoconvex. From our result about $S_{G}$ we can also deduce Guillemin and Sternberg's result and also generalize this result to some class of non-strictly pseudoconvex cases.
(iv) If we assume that the CR manifold $X$ is strictly pseudoconvex of dimension greater or equal to five, the proof is shorter. In this case, from the result of Boutet de Monvel-Sjöstrand [8, we see that the Szegő projection is a complex Fourier integral operator on $X$ and by using integration by parts, we can easily see that $S_{G}$ is smoothing away $\mu^{-1}(0)$. Hence, we do not need Secs. 3.1 3.3 For some class of CR manifolds, the Levi form is just positive near $\mu^{-1}(0)$ (see

Example (1) in Sec. (2.5), therefore it is natural to work with the assumption that the Levi form is just positive near $\mu^{-1}(0)$ since we cannot apply the result of Boutet de Monvel-Sjöstrand directly.

This paper is organized as follows. In Sec. 2, we review some facts on CR and Sasakian manifolds, CR reduction and Szegő kernel. In Sec. 3, we establish asymptotic expansion for the $G$-invariant Szegő kernel. In Sec. 4 we study the distribution kernel of the map $\sigma$ in (1.20) and establish Theorem 1.5 In Sec. 5 we establish Theorem 1.2. In Sec. 6, we establish a version for almost complex manifolds of Theorem 1.4

Notations. We denote by $\mathbb{N}=\{0,1,2, \ldots\}$ the set of natural numbers and set $\mathbb{N}^{*}=\mathbb{N} \backslash\{0\}, \mathbb{R}_{+}=[0, \infty)$. We use standard notations about distributions and Sobolev spaces on manifolds, as in 30, 37. In this paper we will systematically use the correspondence between operators $A$ and their kernels $A(\cdot, \cdot)=A(x, y)$ via the Schwartz kernel theorem [26, Theorems 5.2.1, 5.2.6], [37, Theorem B.2.7]. For two distributions $u, v$ we write $u \equiv v$ if $u-v$ is a smooth function. For two operators $A$, $B$, we write $A \equiv B$ if their Schwartz kernels satisfy $A(\cdot, \cdot) \equiv B(\cdot, \cdot)$, equivalently, if $A-B$ is a smoothing operator.

In the whole paper we will denote by $G$ a compact Lie group, by $\mathfrak{g}$ its Lie algebra, and by $d \mu$ the Haar measure on $G$ with $\int_{G} d \mu(h)=1$. If $E$ is a complex representation of $G$, we denote by $E^{G}$ the $G$-trivial component of $E$.

We denote by $\operatorname{Spec} A$ the spectrum of an operator $A$. For a real vector space/bundle $V$ we denote by $\mathbb{C} V=V \otimes_{\mathbb{R}} \mathbb{C}$ the associated complexified vector space/bundle.

## 2. Preliminaries

In this section, we explain some basic facts on CR and Sasakian manifolds, CR reduction and Szegő kernel.

### 2.1. CR manifolds and $C R$ functions

Let $\left(X, T^{1,0} X\right)$ be a compact, connected and orientable Cauchy-Riemann (CR) manifold of dimension $2 n+1, n \geq 1$, where $T^{1,0} X$ is a CR structure of $X$, that is, $T^{1,0} X$ is a complex vector sub-bundle of rank $n$ of the complexified tangent bundle $\mathbb{C} T X$, satisfying

$$
\begin{equation*}
T^{1,0} X \cap T^{0,1} X=\{0\}, \quad[\mathcal{V}, \mathcal{V}] \subset \mathcal{V}, \text { with } T^{0,1} X=\overline{T^{1,0} X}, \mathcal{V}=\mathscr{C}^{\infty}\left(X, T^{1,0} X\right) \tag{2.1}
\end{equation*}
$$

Denote by $T^{* 1,0} X$ and $T^{* 0,1} X$ the dual bundles of $T^{1,0} X$ and $T^{0,1} X$, respectively. Define the vector bundle of $(0, q)$-forms by

$$
\begin{equation*}
T^{* 0, q} X:=\Lambda^{q} T^{* 0,1} X \tag{2.2}
\end{equation*}
$$

The Levi distribution (or holomorphic tangent space) $H X$ of the CR manifold $X$ is the real part of $T^{1,0} X \oplus T^{0,1} X$, i.e. the unique sub-bundle $H X$ of $T X$ such that

$$
\begin{equation*}
\mathbb{C} H X=T^{1,0} X \oplus T^{0,1} X \tag{2.3}
\end{equation*}
$$

Let $J: H X \rightarrow H X$ be the complex structure given by $J(u+\bar{u})=i u-i \bar{u}$, for every $u \in T^{1,0} X$. If we extend $J$ complex linearly to $\mathbb{C} H X$ we have

$$
\begin{equation*}
T^{1,0} X=\{V \in \mathbb{C} H X: J V=i V\} \tag{2.4}
\end{equation*}
$$

Thus the CR structure $T^{1,0} X$ is determined by the Levi distribution and we shall also write $(X, H X, J)$ to denote the CR manifold $\left(X, T^{1,0} X\right)$.

The annihilator $(H X)^{0} \subset T^{*} X$ of $H X$ is called the characteristic conormal bundle of the CR manifold. Since $X$ is orientable, the characteristic bundle $(H X)^{0}$ is a trivial real line sub-bundle. We fix a global frame of $(H X)^{0}$, that is, a real nonvanishing 1-form $\omega_{0} \in \mathscr{C}^{\infty}\left(X, T^{*} X\right)$ such that $(H X)^{0}=\mathbb{R} \omega_{0}$, called characteristic 1-form. We have

$$
\begin{equation*}
\left\langle\omega_{0}(x), u\right\rangle=0, \quad \text { for any } u \in H_{x} X, \quad x \in X \tag{2.5}
\end{equation*}
$$

Then by (2.1), the restriction of $d \omega_{0}$ on $H X$ is a $(1,1)$-form. The Levi form $\mathscr{L}_{x}=$ $\mathscr{L}_{x}^{\omega_{0}}$ of $X$ at $x \in X$ associated to $\omega_{0}$ is the symmetric bilinear map

$$
\begin{equation*}
\mathscr{L}_{x}: H_{x} X \times H_{x} X \rightarrow \mathbb{R}, \quad \mathscr{L}_{x}(u, v)=\frac{1}{2} d \omega_{0}(u, J v), \quad \text { for } u, v \in H_{x} X \tag{2.6}
\end{equation*}
$$

It induces a Hermitian symmetric map

$$
\begin{equation*}
\mathscr{L}_{x}: T_{x}^{1,0} X \times T_{x}^{1,0} X \rightarrow \mathbb{C}, \quad \mathscr{L}_{x}(U, \bar{V})=\frac{1}{2 i} d \omega_{0}(U, \bar{V}), \quad \text { for } U, V \in T_{x}^{1,0} X \tag{2.7}
\end{equation*}
$$

A CR manifold $X$ is said to be strictly pseudoconvex if there exists a characteristic 1-form $\omega_{0}$ such that for every $x \in X$ the Levi form $\mathscr{L}_{x}^{\omega_{0}}$ is positive definite. In this case $\omega_{0}$ is a contact form and the Levi distribution $H X$ is a contact structure.

Back to the general case, given a characteristic 1-form $\omega_{0}$, let $T \in \mathscr{C}^{\infty}(X, T X)$ be a vector field, called characteristic vector field, such that

$$
\begin{equation*}
\mathbb{C} T X=T^{1,0} X \oplus T^{0,1} X \oplus \mathbb{C} T \tag{2.8}
\end{equation*}
$$

and

$$
\begin{equation*}
i_{T} \omega_{0}=1 \tag{2.9}
\end{equation*}
$$

Let $g^{\mathbb{C} T X}$ be a Hermitian metric on $\mathbb{C} T X$ such that the decomposition (2.8) is orthogonal. For $u, v \in \mathbb{C} T X$ we denote by $\langle u, v\rangle_{g}$ the inner product given by $g^{\mathbb{C} T X}$ and for $u \in \mathbb{C} T X$, we write $|u|_{g}^{2}:=\langle u, u\rangle_{g}$.

The determinant of the Levi form $\mathscr{L}_{x}$ at $x \in X$ with respect to $g^{\mathbb{C} T X}$ is defined by

$$
\begin{equation*}
\operatorname{det} \mathscr{L}_{x}=\lambda_{1}(x) \ldots \lambda_{n}(x) \tag{2.10}
\end{equation*}
$$

where $\lambda_{1}(x), \ldots, \lambda_{n}(x)$, are the eigenvalues of $\mathscr{L}_{x}$ as Hermitian form on $T_{x}^{1,0} X$ with respect to the inner product $\langle\cdot, \cdot\rangle_{g}$ on $T_{x}^{1,0} X$.

The Hermitian metric $g^{\mathbb{C} T X}$ on $\mathbb{C} T X$ induces, by duality, a Hermitian metric on $\mathbb{C} T^{*} X$ and also on the bundles of $(0, q)$ forms $T^{* 0, q} X, q=1,2, \ldots, n$. We shall also denote the inner product given by these metrics by $\langle\cdot, \cdot\rangle_{g}$. The metric $g^{\mathbb{C} T X}$ induces a Riemannian metric $g^{T X}$ on $T X$ and $g^{T X}$ induces in turn a Riemannian volume form $d v=d v(x)$ on $X$ and a distance function $d(\cdot, \cdot)$ on $X$.

The natural global $L^{2}$ inner product $(\cdot, \cdot)$ on $\Omega^{0, q}(X)$ induced by $d v(x)$ and $\langle\cdot, \cdot\rangle_{g}$ is given by

$$
\begin{equation*}
(u, v):=\int_{X}\langle u(x), v(x)\rangle_{g} d v(x), \quad u, v \in \Omega^{0, q}(X) \tag{2.11}
\end{equation*}
$$

We denote by $\left(L_{(0, q)}^{2}(X),(\cdot, \cdot)\right)$ the completion of $\Omega^{0, q}(X)$ with respect to $(\cdot, \cdot)$ and denote $\|\cdot\|$ the corresponding $L^{2}$ norm. We set $L^{2}(X):=L_{(0,0)}^{2}(X)$.

Let $\bar{\partial}_{b}: \Omega^{0, q}(X) \rightarrow \Omega^{0, q+1}(X)$ be the tangential CR operators on $X$ which is the composition of the exterior differential $d$ and the projection $\pi^{0, q+1}: \Lambda^{q+1}\left(\mathbb{C} T^{*} X\right) \rightarrow$ $T^{* 0, q+1} X$. We consider the weak maximal extension of $\bar{\partial}_{b}$ to $L^{2}$ spaces as follows:

$$
\begin{align*}
\operatorname{Dom} \bar{\partial}_{b} & =\left\{u \in L_{(0, q)}^{2}(X): \bar{\partial}_{b} u \in L_{(0, q+1)}^{2}(X)\right\},  \tag{2.12}\\
\bar{\partial}_{b}: \operatorname{Dom} \bar{\partial}_{b} \ni u & \mapsto \bar{\partial}_{b} u \in L_{(0, q+1)}^{2}(X),
\end{align*}
$$

where $\bar{\partial}_{b} u$ is defined in the sense of distributions. The space of $L^{2} \mathrm{CR}$ functions on $X$ is given by

$$
\begin{equation*}
H_{b}^{0}(X):=\left\{u \in L^{2}(X)=L_{(0,0)}^{2}(X): \bar{\partial}_{b} u=0\right\} \tag{2.13}
\end{equation*}
$$

Since differential operators are continuous on distributions, $H_{b}^{0}(X)$ is a closed subspace of $L^{2}(X)$. The Szegö projection is the orthogonal projection

$$
\begin{equation*}
S:\left(L^{2}(X),(\cdot, \cdot)\right) \rightarrow H_{b}^{0}(X) \tag{2.14}
\end{equation*}
$$

The Szegő kernel $S(x, y) \in \mathscr{D}^{\prime}(X \times X)$ is the distribution kernel of $S$.

### 2.2. Sasakian manifolds

We recall here some facts about Sasakian manifolds, cf. 9, 17. Recently, the subject of Sasakian geometry generated a great deal of interest due to the study of existence of Sasaki-Einstein metrics, and more generally, Sasakian metrics of constant scalar curvature, see for example [13].

Let $(X, H X, J)$ be an orientable strictly pseudoconvex CR manifold of dimension $2 n+1$ and let $\omega_{0}$ be a contact form whose Levi form (2.6) is positive definite, hence $g^{H X}=d \omega_{0}(\cdot, J \cdot)$ defines a $J$-invariant metric on $H X$. The Reeb vector field $R$ associated to the contact form $\omega_{0}$ is the vector field on $X$ defined by

$$
\begin{equation*}
i_{R} \omega_{0}=1, \quad i_{R} d \omega_{0}=0 \tag{2.15}
\end{equation*}
$$

From (2.15), $0=d \omega_{0}(R, u)=-\omega_{0}([R, u])$ for any $u \in \mathscr{C}^{\infty}(X, H X)$, thus we get

$$
\begin{equation*}
\left[R, \mathscr{C}^{\infty}(X, H X)\right] \subset \mathscr{C}^{\infty}(X, H X) \tag{2.16}
\end{equation*}
$$

We define a Riemannian metric $g_{\omega_{0}}$ on $X$ by $g_{\omega_{0}}(\cdot, \cdot)=d \omega_{0}(\cdot, J \cdot)+\omega_{0}(\cdot) \omega_{0}(\cdot)$. Associated to the data $\left(X, \omega_{0}, R, J, g_{\omega_{0}}\right)$, called contact metric manifold, there is a canonical connection $\nabla$ on $T X$, called the Tanaka-Webster connection (see Tanaka 51] and Webster [56]), that is the unique affine connection on $T X$ such that

- $\nabla g_{\omega_{0}}=0, \nabla J=0, \nabla d \omega_{0}=0$ and $\nabla$ preserves the decomposition $T X=H X \oplus$ $\mathbb{R} R$.
- For any $u, v$ in the Levi distribution $H X$, the torsion $T_{\nabla}$ of $\nabla$ satisfies $T_{\nabla}(u, v)=$ $d \omega_{0}(u, v) R$ and $T_{\nabla}(R, J u)=-J T_{\nabla}(R, u)$.

The torsion of the Tanaka-Webster connection in the direction of the Reeb vector field, $\tau: u \mapsto T_{\nabla}(R, u)$, is called pseudo-Hermitian torsion of $\nabla$. We see that $\nabla R=0$ and thus $\nabla_{R}=\tau+\mathcal{L}_{R}$, where $\mathcal{L}$ denotes the Lie derivative. By using $\nabla J=0$ we deduce that (cf. 51, Lemma 3.2(3)])

$$
\begin{equation*}
2 J \tau u=\left(\mathcal{L}_{R} J\right) u \quad \text { for any } u \in H X \tag{2.17}
\end{equation*}
$$

Definition 2.1. A contact metric manifold $\left(X, \omega_{0}, R, J, g_{\omega_{0}}\right)$ is called a Sasakian manifold if the pseudo-Hermitian torsion of its Tanaka-Webster connection vanishes: $\tau=T_{\nabla}(R, \cdot)=0$.

By (2.17), $\tau=0$ is equivalent to $\mathcal{L}_{R} J=0$ and by [9, Corollary 6.5.11] this is equivalent to the fact that the contact metric structure on $X$ is normal. Thus the definition above is equivalent to the definition of Sasaki manifolds given in 9, Definition 6.5.13] (which in turn is equivalent to the metric cone $\left(C(X)=\mathbb{R}_{+} \times\right.$ $X, d r^{2}+r^{2} g_{\omega_{0}}$ ) being a Kähler manifold [9, Definition 6.5.15]).

Sasakian manifolds can be classified in three categories based on the properties of the Reeb foliation $\mathcal{F}_{R}$ consisting of the orbits of the Reeb field (see [9, Definition 6.1.25]). If the orbits of the Reeb field are all closed, then the Reeb field $R$ generates a locally free, isometric $S^{1}$-action on $\left(X, g_{\omega_{0}}\right)$ and the Reeb foliation is called quasiregular. If this $S^{1}$-action is free, then the Reeb foliation is said to be regular. If $\mathcal{F}_{R}$ is not quasi-regular, it is said to be irregular. In this case, the flow associated with $R$ generates a transversal CR $\mathbb{R}$-action on $X$. We say that the $\mathbb{R}$-action $\eta$ on $X$ is CR transversal if $H X \oplus \mathbb{R} \eta_{X}=T X$, where $\eta_{X}$ denotes the infinitesimal generator field of the $\mathbb{R}$-action.

If $\mathcal{F}_{R}$ is quasi-regular, then by the structure theorem [9, Theorem 7.1.3] the quotient space $M:=X / \mathcal{F}_{R}=X / S^{1}$ is a Kähler orbifold and the quotient map $\pi: X \rightarrow M$ an orbifold Riemannian submersion. Moreover, $X$ is the total space of a principal $S^{1}$-bundle over $M$ with connection 1-form $\omega_{0}$; there exists an integral Kähler form $\omega$ on $M$ such that the curvature $d \omega_{0}$ of $\omega_{0}$ is the pullback by the quotient map of $\omega: d \omega_{0}=\pi^{*} \omega$. If $\mathcal{F}_{R}$ is regular, then $(M, \omega)$ is a Hodge manifold and the assertions above are the content of the Boothby-Wang theorem [5] (cf. also [22]).

Remark 2.2. For a Sasakian manifold $\left(X, \omega_{0}, R, J, g_{\omega_{0}}\right)$, from (2.16), the flow associated with $R$ preserves $H X$. Moreover, it follows from Definition 2.1 and (2.17)
that $\mathcal{L}_{R} J=0$. Combining with (2.15), we get $\mathcal{L}_{R} g_{\omega_{0}}=0$, i.e. the Reeb vector field is a Killing vector field on $\left(X, g_{\omega_{0}}\right)$.

Remark 2.3. Let $\left(X, T^{1,0} X\right)$ be a strictly pseudoconvex CR manifold. Then $X$ is Sasakian if and only if there is a global vector field $T \in \mathscr{C}^{\infty}(X, T X)$ such that

$$
\begin{equation*}
\left[T, \mathscr{C}^{\infty}\left(X, T^{1,0} X\right)\right] \subset \mathscr{C}^{\infty}\left(X, T^{1,0} X\right) \tag{2.18}
\end{equation*}
$$

and $T^{1,0} X \oplus T^{0,1} X$ and $T$ generate the complex tangent bundle of $X$ (cf. [9, Proposition 6.4.8], [46, Theorem 1.2]). In fact, if $X$ is Sasakian, we can take $T=R$ by (2.17). For the inverse direction, we define $\omega_{0} \in \Omega^{1}(X)$ by $i_{T} \omega_{0}=1,\left.\omega_{0}\right|_{H X}=0$. Now for any $V \in \mathscr{C}^{\infty}\left(X, T^{1,0} X\right)$ we get by (2.18) that $\left(i_{T} d \omega_{0}\right)(V)=-\omega_{0}([T, V])=0$. Thus $\omega_{0}$ verifies (2.15). Again by (2.18) we have $\mathcal{L}_{T} J=0$, thus $X$ is Sasakian by (2.17).

Related to Theorem 1.4 and motivated by the structure of quasi-regular Sasakian manifolds let us consider now the case of a circle bundle associated to a Hermitian holomorphic line bundle. Let $\left(L, h^{L}\right)$ be a Hermitian holomorphic line bundle over a connected compact complex manifold $(M, J)$. Let $h^{L^{*}}$ be the Hermitian metric on $L^{*}$ induced by $h^{L}$. Let

$$
\begin{equation*}
X:=\left\{v \in L^{*}:|v|_{h^{L^{*}}}^{2}=1\right\} \tag{2.19}
\end{equation*}
$$

be the circle bundle of $L^{*}$ (Grauert tube); it is isomorphic to the $S^{1}$ principal bundle associated to $L$. Since $X$ is a hypersurface in the complex manifold $L^{*}$, it has a CR structure inherited from the complex structure of $L^{*}$ by setting $T^{1,0} X=$ $T X \cap T^{1,0} L^{*}$.

In this situation, $S^{1}$ acts on $X$ by fiberwise multiplication, denoted $\left(x, e^{i \theta}\right) \mapsto$ $x e^{i \theta}$. A point $x \in X$ is a pair $x=(p, \lambda)$, where $\lambda$ is a linear functional on $L_{p}$, the $S^{1}$ action is $x e^{i \theta}=(p, \lambda) e^{i \theta}=\left(p, e^{i \theta} \lambda\right)$.

On $X$ we have a globally defined vector field $\partial_{\theta}$, the generator of the $S^{1}$ action. The span of $\partial_{\theta}$ defines a rank one subbundle $T^{V} X \cong T S^{1} \subset T X$, the vertical subbundle of the fibration $\pi: X \rightarrow M$. Moreover (2.8) holds for $T=\partial_{\theta}$.

For $m \in \mathbb{Z}$ the space $\mathscr{C}^{\infty}\left(X, L^{m}\right)$ of smooth sections of $L^{m}$ can be identified to the space $m$-equivariant smooth functions

$$
\mathscr{C}^{\infty}(X)_{m}=\left\{f \in \mathscr{C}^{\infty}(X, \mathbb{C}): f\left(x e^{i \theta}\right)=e^{i m \theta} f(x), \text { for } e^{i \theta} \in S^{1}, x \in X\right\}
$$

by

$$
\begin{equation*}
\mathscr{C}^{\infty}\left(M, L^{m}\right) \ni s \mapsto f \in \mathscr{C}^{\infty}(X)_{m}, \quad f(x)=f(p, \lambda)=\lambda^{\otimes m}(s(p)) \tag{2.20}
\end{equation*}
$$

where $\lambda^{m}=\lambda^{\otimes m}$ for $m \geq 0$ and $\lambda^{m}=\left(\lambda^{-1}\right)^{\otimes(-m)}$ for $m<0$. Through the identification (2.20) holomorphic sections correspond to CR functions

$$
\begin{equation*}
H^{0}\left(M, L^{m}\right) \cong H_{b, m}^{0}(X):=\left\{f \in \mathscr{C}^{\infty}(X)_{m}: \bar{\partial}_{b} f=0\right\} \tag{2.21}
\end{equation*}
$$

We construct now a Riemannian metric on $X$. Let $g^{T M}$ be a $J$-invariant metric on $M$. The Chern connection $\nabla^{L}$ on $L$ induces a connection on the $S^{1}$-principal bundle $\pi: X \rightarrow M$, and let $T^{H} X \subset T X$ be the corresponding horizontal bundle.

Let $g^{T X}=\pi^{*} g^{T M} \oplus \frac{d \theta^{2}}{4 \pi^{2}}$ be the metric on $T X=T^{H} X \oplus T S^{1}$, with $d \theta^{2}$ the standard metric on $S^{1}=\mathbb{R} / 2 \pi \mathbb{Z}$.

Pertaining to $g^{T X}$ we construct the $L^{2}$ inner product $(\cdot, \cdot)_{X}$ given by (2.11) on $X$. The metric $g^{T M}$ induces a Riemannian volume form $d v_{M}$ on $M$, which together with the fiber metric $h^{L^{m}}$ gives rise to an $L^{2}$ inner product $(\cdot, \cdot)_{m}$ on $\mathscr{C}^{\infty}\left(X, L^{m}\right)$. Then the isomorphism (2.20) becomes an isometry $\left(\mathscr{C}^{\infty}\left(M, L^{m}\right),(\cdot, \cdot)_{m}\right) \cong$ $\left(\mathscr{C}^{\infty}(X)_{m},(\cdot, \cdot)_{X}\right)$ and accordingly an isometry $L^{2}\left(M, L^{m}\right) \cong L^{2}(X)_{m}$, where the latter space is the completion of $\left(\mathscr{C}^{\infty}(X)_{m},(\cdot, \cdot)_{X}\right)$. Moreover, (2.20) induces an isometry

$$
\begin{equation*}
\left(H^{0}\left(M, L^{m}\right),(\cdot, \cdot)_{m}\right) \cong\left(H_{b, m}^{0}(X),(\cdot, \cdot)_{X}\right) \tag{2.22}
\end{equation*}
$$

The $S^{1}$-action gives rise to a Fourier decomposition $L^{2}(X) \cong \widehat{\bigoplus}_{m \in \mathbb{Z}} L^{2}(X)_{m}$ and this induces the following decomposition at the level of CR functions:

$$
\begin{equation*}
H_{b}^{0}(X) \cong \widehat{\bigoplus}_{m \in \mathbb{Z}} H_{b, m}^{0}(X) \cong \widehat{\bigoplus}_{m \in \mathbb{Z}} H^{0}\left(M, L^{m}\right) \tag{2.23}
\end{equation*}
$$

Let $\omega_{0}$ be the connection 1-form on $X$ associated to the Chern connection $\nabla^{L}$. Then $\omega_{0}\left(\partial_{\theta}\right)=1$, thus (2.8) and (2.9) are fulfilled and $T=\partial_{\theta}$ is a characteristic vector field on $X$ and $\omega_{0}$ is a characteristic 1-form for the CR structure on $X$. Moreover,

$$
\begin{equation*}
d \omega_{0}=\pi^{*}\left(i R^{L}\right) \tag{2.24}
\end{equation*}
$$

where $R^{L}$ is the curvature of $\nabla^{L}$. On account of (2.6) $X$ is strictly pseudoconvex at $x \in X$ if and only if $\left(L, h^{L}\right)$ is positive at $\pi(x) \in M$. In particular, if $\left(L, h^{L}\right)$ is positive on $M, X$ is a strictly pseudoconvex CR manifold, $\omega_{0}$ is a contact form and $\partial_{\theta}$ is the associated Reeb vector field.

Assume now $\left(L, h^{L}\right)$ is positive on $M$. We claim that $X$ is a Sasakian manifold. In fact, the $S^{1}$-action is fiberwise multiplication on $L^{*}$, thus $S^{1}$ acts holomorphically on $L^{*}$ and preserves $T^{1,0} L^{*}$. This means that the $S^{1}$-action preserves $T^{1,0} X=$ $T X \cap T^{1,0} L^{*}$, i.e. $\mathcal{L}_{\partial_{\theta}} J=0$, hence $X$ is Sasakian by (2.17). Since the $S^{1}$-action is free, $X$ is a regular Sasakian manifold. Conversely, any compact regular Sasakian manifold is CR-isomorphic to a $S^{1}$-fibration associated with a positive line bundle on a projective manifold by the Boothby-Wang theorem [5, [9, Theorem 6.1.26], [22.

Note also that if $\left(L, h^{L}\right)$ is positive on $M$ then $H^{0}\left(X, L^{m}\right)=0$ for $m<0$ by the Kodaira vanishing theorem, so the decomposition (2.23) becomes

$$
\begin{equation*}
H_{b}^{0}(X) \cong \widehat{\bigoplus}_{m \in \mathbb{N}} H_{b, m}^{0}(X) \cong \widehat{\bigoplus}_{m \in \mathbb{N}} H^{0}\left(M, L^{m}\right) \tag{2.25}
\end{equation*}
$$

Note further that $X$ is (weakly) pseudoconvex, that is, the Levi form is positive semidefinite on $X$ if and only if the curvature $i R^{L}$ is semi-positive on $M$. Moreover, $X$ has finite type if and only if $R^{L}$ vanishes to finite order at any point of $M$, cf. [31, 41, Proposition 11]. A CR manifold $X$ is said to be of finite type if at each point $p \in X$ if the space $T_{p} X$ is generated by vectors of the form $\left[V_{1},\left[V_{2},\left[\ldots\left[V_{k-1}, V_{k}\right] \ldots\right]\right]\right](p)$, where $V_{1}, \ldots, V_{k}, k \geq 2$, are sections of $H X$.

### 2.3. CR reduction

We refer to [1, 9, 17] for the fundamentals of contact geometry and examples. In this subsection, we extend the well-known symplectic reduction of Kähler manifolds to CR manifolds. We refer to [1, 9, 16, 17, 57, for constructions and examples of contact and Sasakian reduction. Our construction is a CR analogue of 48, §2.1].

Let $(X, H X, J)$ be a compact connected and orientable CR manifold of dimension $2 n+1, n \geq 1$, and let $\omega_{0}$ be a characteristic 1-form.

Let $G$ be a $d$-dimensional compact Lie group with Lie algebra $\mathfrak{g}$. We assume that $G$ acts smoothly on $X$ and that the $G$-action preserves $J$ and $\omega_{0}$.

For any $\xi \in \mathfrak{g}$, we denote $\xi_{X}(x)=\left.\frac{\partial}{\partial t} \exp (-t \xi) x\right|_{t=0}$ the vector field on $X$ induced by $\xi$. For $x \in X$, set

$$
\begin{equation*}
\underline{\mathfrak{g}}_{x}=\operatorname{Span}\left\{\xi_{X}(x): \xi \in \mathfrak{g}\right\} \tag{2.26}
\end{equation*}
$$

Definition 2.4. The moment map associated to the characteristic 1-form $\omega_{0}$ is the map $\mu: X \rightarrow \mathfrak{g}^{*}$ defined by

$$
\begin{equation*}
\langle\mu(x), \xi\rangle=\omega_{0}\left(\xi_{X}(x)\right), \quad x \in X, \quad \xi \in \mathfrak{g} . \tag{2.27}
\end{equation*}
$$

The moment map is $G$-equivariant, i.e. for $x \in X, h \in G$, we have

$$
\begin{equation*}
\mu(h . x)=\operatorname{Ad}_{h}^{*} \mu(x) . \tag{2.28}
\end{equation*}
$$

Relation (2.28) implies that $G$ acts on $\mu^{-1}(0)$. In fact, for any $\xi \in \mathfrak{g}$, we have

$$
\begin{align*}
\langle\mu(h \cdot x), \xi\rangle & \left.\left.=\omega_{0}\left(\xi_{X}(h \cdot x)\right)\right)=\omega_{0}\left(d h\left(\operatorname{Ad}_{h^{-1}} \xi\right)_{X}\right)\right)_{h \cdot x} \\
& =\left(h^{*} \omega_{0}\right)\left(\left(\operatorname{Ad}_{h^{-1}} \xi\right)_{X}\right)_{x}=\omega_{0}\left(\left(\operatorname{Ad}_{h^{-1}} \xi\right)_{X}\right)_{x} \\
& =\left\langle\mu(x), \operatorname{Ad}_{h^{-1}} \xi\right\rangle=\left\langle\operatorname{Ad}_{h}^{*} \mu(x), \xi\right\rangle . \tag{2.29}
\end{align*}
$$

Lemma 2.5. If $G$ acts freely on $\mu^{-1}(0)$ and the Levi form is positive on $\mu^{-1}(0)$, then 0 is a regular value of $\mu$.

Proof. Observe first that since $\omega_{0}$ is $G$-invariant, (2.5) yields that $H X$ is an $G$ equivariant sub-bundle of $T X$, thus

$$
\begin{equation*}
\left[\xi_{X}, U\right] \in \mathscr{C}^{\infty}(X, H X) \quad \text { for any } U \in \mathscr{C}^{\infty}(X, H X), \xi \in \mathfrak{g} \tag{2.30}
\end{equation*}
$$

Now (2.5) and (2.27) entail

$$
\begin{equation*}
\xi_{X}(x) \in H X \quad \text { if } x \in \mu^{-1}(0) \tag{2.31}
\end{equation*}
$$

Let $U_{x} \in H_{x} X$ that we extend to a section $U$ of $H X$ near $x$. Then (2.5) and (2.30) yield

$$
\begin{align*}
d \omega_{0}\left(U, \xi_{X}\right)_{x} & =U\left(\omega_{0}\left(\xi_{X}\right)\right)-\xi_{X}\left(\omega_{0}(U)\right)-\omega_{0}\left(\left[U, \xi_{X}\right]\right) \\
& =U_{x}\langle\mu, \xi\rangle=\left\langle U(\mu)_{x}, \xi\right\rangle \tag{2.32}
\end{align*}
$$

If $d \mu_{y}: H X \rightarrow \mathfrak{g}^{*}$ were not surjective for some $y \in \mu^{-1}(0)$, there would exist $\xi \in \mathfrak{g}$ such that $\left\langle Y(\mu)_{y}, \xi\right\rangle=0$ for any $Y \in H_{y} X$. This is a contradiction since $d \omega_{0}$ is nondegenerate on $H_{y} X$ and $0 \neq \xi_{X, y} \in H_{y} X$ by (2.31). Thus $d \mu_{y}: H X \rightarrow \mathfrak{g}^{*}$ is surjective for $y \in \mu^{-1}(0)$.

Set as in (1.1) $X_{G}=Y / G$ with $Y=\mu^{-1}(0)$. Let $\iota: Y \rightarrow X$ be the natural injection and let $\pi: Y \rightarrow X_{G}=Y / G$ be the natural projection.

Theorem 2.6. If $G$ acts freely on $Y=\mu^{-1}(0)$ and the Levi form is positive on $\mu^{-1}(0)$, then the reduced space $X_{G}=Y / G$ is a strictly pseudoconvex manifold with contact form $\omega_{0, G}$ satisfying $\iota^{*} \omega_{0}=\pi^{*} \omega_{0, G}$. Moreover, we can choose the characteristic vector field $T$ (cf. (2.8), (2.9) ) such that $\left.T\right|_{Y} \in \mathscr{C}^{\infty}(Y, T Y)$ and $T$ is $G$-invariant.

Proof. By Lemma $2.5 \mu^{-1}(0)$ is a smooth manifold. Since $G$ acts freely on $Y, X_{G}$ is a compact manifold. The positivity of the Levi form on $\mu^{-1}(0)$ means that

$$
\begin{equation*}
g^{H X}=d \omega_{0}(\cdot, J \cdot) \tag{2.33}
\end{equation*}
$$

is a $J$-invariant and $G$-equivariant metric on $H X$ on a neighborhood of $Y=\mu^{-1}(0)$.
Since $G$ acts freely on $Y$, the vector spaces $\underline{\mathfrak{g}}_{x}$ defined in (2.26) form a vector bundle $\underline{\mathfrak{g}}$ near $\mu^{-1}(0)$. We denote $\underline{\mathfrak{g}}_{Y}=\left.\underline{\mathfrak{g}}\right|_{Y}$. Then $\underline{\mathfrak{g}}_{Y} \subset T Y \cap H X$ by (2.31).

For $\bar{x} \in \mu^{-1}(0)$, by (2.31), (2.32), and the fact that $d \omega_{0}(\cdot, J \cdot)$ is a metric on $H_{x} X$ we have that $\left.d \mu\right|_{T Y}=0$ and $\left.d \mu\right|_{\underline{\mathfrak{g}}_{x}} \rightarrow \mathfrak{g}^{*}$ is surjective. Since $\operatorname{dim} Y+\operatorname{dim} \mathfrak{g}=$ $\operatorname{dim} T X$, we have

$$
\begin{equation*}
\left.J \underline{\mathfrak{g}}\right|_{Y} \oplus T Y=\left.T X\right|_{Y} \tag{2.34}
\end{equation*}
$$

From (2.34) and $\left.J \underline{\mathfrak{g}}\right|_{Y} \subset H X$, we know $\omega_{0}(T Y) \neq 0$. Thus $T Y \cap H X$ is a codimension 1 sub-bundle of $T Y$, and

$$
\begin{equation*}
\frac{T Y}{T Y \cap H X}=\left.\frac{T X}{H X}\right|_{Y} \tag{2.35}
\end{equation*}
$$

From (2.35), we can choose the vector field $T$ in (2.8) such that $\left.T\right|_{Y} \in \mathscr{C}^{\infty}(Y, T Y)$ and $T$ is $G$-invariant.

Let $T^{H} Y$ be the orthogonal complement of $\underline{\mathfrak{g}}_{Y}$ in $T Y \cap H X$ with respect to $g^{H X}$. By (2.32), (2.33), for any $U \in T Y \cap H X, \xi \in \mathfrak{g}$, as $J \underline{\mathfrak{g}}_{Y} \subset H X$, we have $g^{H X}\left(U, J \xi_{X}\right)=-d \omega_{0}\left(U, \xi_{X}\right)=0$. This means that $T Y \cap H X$ is orthogonal with $J \underline{\mathfrak{g}}_{Y}$ with respect to $g^{H X}$. Thus we have the $G$-equivariant orthogonal decomposition on $Y$,

$$
\begin{equation*}
T Y \cap H X=T^{H} Y \oplus \underline{\mathfrak{g}}_{Y},\left.\quad H X\right|_{Y}=\left.T^{H} Y \oplus \underline{\mathfrak{g}}_{Y} \oplus J \underline{\mathfrak{g}}\right|_{Y}, \tag{2.36}
\end{equation*}
$$

note that the second equation is from the dimension counting argument as $T Y \cap H X$ is a codimension 1 sub-bundle of $T Y$. Thus from (2.36) and the metric $g^{H X}$ on
$\left.H X\right|_{Y}$ is $J$-invariant, we get

$$
\begin{equation*}
J T^{H} Y=T^{H} Y=(T Y \cap H X) \cap J(T Y \cap H X) \tag{2.37}
\end{equation*}
$$

By (2.27) $\iota^{*} \omega_{0}$ is a $G$-invariant horizontal 1-form on $Y$, thus there exists a unique 1-form $\omega_{0, G} \in \Omega^{1}\left(X_{G}\right)$ such that

$$
\begin{equation*}
\iota^{*} \omega_{0}=\pi^{*} \omega_{0, G} \tag{2.38}
\end{equation*}
$$

We now define the Levi distribution on $X_{G}$ by

$$
\begin{equation*}
H X_{G}:=\operatorname{ker} \omega_{0, G} \tag{2.39}
\end{equation*}
$$

From (2.35), (2.36), $d \pi: \mathbb{R} T \oplus T^{H} Y \rightarrow T X_{G}$ is bijective, and we get the isomorphism $\mathbb{R} T \oplus T^{H} Y \simeq \pi^{*} T X_{G}$. Thus $d \pi$ maps $T^{H} Y$ onto $H X_{G}$ and this gives an isomorphism $T^{H} Y \simeq \pi^{*} H X_{G}$. Thus for $U \in H_{y} X_{G}$, we take $x \in \pi^{-1}(y)$ and $U^{H} \in T^{H} Y$ the lift of $U$, then by (2.37), we define $J_{G} \in \operatorname{End}\left(H X_{G}\right)$ by

$$
\begin{equation*}
\left(J_{G} U\right)^{H}=J U^{H} . \tag{2.40}
\end{equation*}
$$

From (2.38), we have

$$
\begin{equation*}
\iota^{*} d \omega_{0}=\pi^{*} d \omega_{0, G} \tag{2.41}
\end{equation*}
$$

Thus from (2.37), $\iota^{*} d \omega_{0}(\cdot, J \cdot)$ is positive and $G$-invariant on $T^{H} Y$ implies that $d \omega_{0, G}\left(\cdot, J_{G} \cdot\right)$ is positive and $J_{G}$-invariant on $H X_{G}$. We verify now that

$$
\begin{equation*}
T^{1,0} X_{G}=\left\{u-\sqrt{-1} J_{G} u: u \in H X_{G}\right\} . \tag{2.42}
\end{equation*}
$$

defines a CR structure on $X$. For $U, V \in \mathscr{C}^{\infty}\left(X_{G}, H X_{G}\right)$, from (2.40),

$$
\begin{equation*}
\left(U-\sqrt{-1} J_{G} U\right)^{H}=U^{H}-\sqrt{-1} J U^{H} \in \mathscr{C}^{\infty}\left(Y, T^{1,0} X \cap \mathbb{C} T Y\right), \tag{2.43}
\end{equation*}
$$

thus by (2.1), $\left[U^{H}-\sqrt{-1} J U^{H}, V^{H}-\sqrt{-1} J V^{H}\right] \in \mathscr{C}^{\infty}\left(Y, T^{1,0} X \cap \mathbb{C} T Y\right)$. By (2.36), $T^{1,0} X \cap \mathbb{C} T Y=\left\{v-\sqrt{-1} J v: v \in T^{H} Y\right\}$. Thus there exists $W \in \mathscr{C}{ }^{\infty}\left(X_{G}, H X_{G}\right)$ such that

$$
\begin{equation*}
\left[U^{H}-\sqrt{-1} J U^{H}, V^{H}-\sqrt{-1} J V^{H}\right]=W^{H}-\sqrt{-1} J W^{H} \tag{2.44}
\end{equation*}
$$

From (2.43), (2.44), we obtain

$$
\begin{align*}
{\left[U-\sqrt{-1} J_{G} U, V-\sqrt{-1} J_{G} V\right] } & =d \pi\left[U^{H}-\sqrt{-1} J U^{H}, V^{H}-\sqrt{-1} J V^{H}\right] \\
& =W-\sqrt{-1} J_{G} W \tag{2.45}
\end{align*}
$$

i.e. $\left[\mathscr{C}^{\infty}\left(X_{G}, T^{1,0} X_{G}\right), \mathscr{C}^{\infty}\left(X_{G}, T^{1,0} X_{G}\right)\right] \subset \mathscr{C}^{\infty}\left(X_{G}, T^{1,0} X_{G}\right)$. Let us finally note that (2.35) shows that we can choose the characteristic vector field $T$ (cf. (2.8), (2.9)) such that $\left.T\right|_{Y} \in \mathscr{C}^{\infty}(Y, T Y)$ and $T$ is $G$-invariant. The proof of Theorem 2.6 is completed.

In the rest of this paper we will work under Assumption 1.1 .

Lemma 2.7. Under Assumption 1.1 there is a $G$-invariant Hermitian metric $g=$ $g^{\mathbb{C} T X}$ on $\mathbb{C} T X$ so that
(i) $T^{1,0} \mathrm{X}$ is orthogonal to $T^{0,1} \mathrm{X}$,
(ii) $\mathfrak{g}$ is orthogonal to $H Y \cap J H Y$ at every point of $Y$,
(iii) $\overline{\langle } T, T\rangle_{g}=1$,
(iv) $T$ is orthogonal to $T^{1,0} X \oplus T^{0,1} X$,
where on $Y, H Y:=H X \cap T Y$.

Proof. This follows from the proof of Theorem [2.6. Let $U$ be a $G$-invariant neighborhood of $Y$ so that the Levi form is positive definite on $U$. Then the metric $g^{H X}=d \omega_{0}(\cdot, J \cdot)$ is a $J$-invariant and $G$-equivariant metric on $H X$ and we have the orthogonal decomposition (2.36) on $Y$. Now we extend the metric $g^{H X}$ from $U$ to $X$ as a $J$-invariant and $G$-equivariant metric on $H X$ by a partition of unity argument. Thus we can take $g^{T X}$ on $T X=\mathbb{R} T \oplus H X$ as the direct sum metric on $\left(H X, g^{H X}\right)$ and $(\mathbb{R} T,\langle T, T\rangle=1)$.

Convention 2.8. From now on we fix a $G$-invariant Hermitian metric $g=g^{\mathbb{C} T X}$ on $\mathbb{C} T X$ so that (i)-(iv) in Lemma 2.7 hold. This metric induces natural Hermitian metrics $\langle\cdot, \cdot\rangle_{X_{G}}$ on $\mathbb{C} T X_{G}$ and $\mathbb{C} T^{*} X_{G}$. As in (2.11) we define the $L^{2}$ inner products and spaces induced by $g$ on $X$ and $X_{G}$ by $\left(L_{(0, q)}^{2}(X),(\cdot, \cdot)\right)$ and $\left(L_{(0, q)}^{2}\left(X_{G}\right),(\cdot, \cdot)_{X_{G}}\right)$.

### 2.4. Closed range in $L^{2}$ for $\bar{\partial}_{b}$ and Szegö projections

The property of closed range in $L^{2}$ for $\bar{\partial}_{b}$ in (2.12) plays an important role in CR geometry. It follows from the works of Boutet de Monvel [6, Boutet de Monvel-Sjöstrand [8, Harvey-Lawson [21, Burns [10] and Kohn [33] that the conditions below are equivalent for a compact strictly pseudoconvex CR manifold $X$, $\operatorname{dim}_{\mathbb{R}} X \geq 3$ :
(a) $X$ is embeddable in the Euclidean space $\mathbb{C}^{N}$, for $N$ sufficiently large;
(b) $X$ bounds a strictly pseudoconvex complex manifold;
(c) The tangential CR operator $\bar{\partial}_{b}: \operatorname{Dom} \bar{\partial}_{b} \subset L^{2}(X) \rightarrow L_{(0,1)}^{2}(X)$ on functions has closed range.

If $X$ is a compact strictly pseudoconvex CR manifold of dimension greater than or equal to five, then $X$ satisfies condition (a), by the embedding theorem of Boutet de Monvel [6. However, there are examples of non-embeddable compact strictly pseudoconvex CR manifolds of dimension three given by Grauert et al. [2, 18, 49, see also Colţoiu-Tibăr [14. In fact this happens for arbitrarily small perturbations of the standard CR structure on the unit sphere in $\mathbb{C}^{2}$. For these examples the closed range in $L^{2}$ property fails.

Assume that condition (b) is satisfied and let $M$ be a strictly pseudoconvex complex manifold such that $\partial M=X$. If $u$ is continuous on $\bar{M}$ and holomorphic on $M$, then $\left.u\right|_{X}$ satisfies the tangential CR equations $\bar{\partial}_{b}\left(\left.u\right|_{X}\right)=0$. Conversely, by 34] any smooth function $u$ on $X$ satisfying $\bar{\partial}_{b} u=0$ admits a smooth extension $\widetilde{u}$ to $\bar{M}$ which is holomorphic in $M$. In this sense the space $H_{b}^{0}(X) \cap \mathscr{C}^{\infty}(X)$ is the space of boundary values of holomorphic functions $\mathcal{O}(M) \cap \mathscr{C}^{\infty}(\bar{M})$. Note also that the Hardy space $H_{b}^{0}\left(S^{1}\right)$ consists of boundary values of holomorphic functions on the unit disc in $\mathbb{C}$ cf. [50, Theorem 17.10]. This is the unifying feature of our definition of $H_{b}^{0}(X)$ for strictly pseudoconvex $X$ and $X=S^{1}$.

There are important classes of embeddable compact strictly pseudoconvex threedimensional CR manifolds (for which $\bar{\partial}_{b}$ has thus closed range in $L^{2}$ ) carrying interesting geometric structures such as

- transverse CR $S^{1}$-actions (4), 15, 35,
- conformal structures [3],
- Sasakian structures (transverse CR $\mathbb{R}$-actions) 42],

If $X$ is a compact strictly pseudoconvex CR manifold and $\bar{\partial}_{b}$ has closed range in $L^{2}$, Boutet de Monvel-Sjöstrand [8] showed that $S(x, y)$ is a Fourier integral operator with complex phase. In particular, $S(x, y)$ is smooth outside the diagonal of $X \times X$ and there is a precise description of the singularity on the diagonal $x=y$, where $S(x, x)$ has a certain asymptotic expansion. Hsiao [27, Theorem 1.2] generalized Boutet de Monvel-Sjöstrand's result to $(0, q)$ forms when the Levi form is non-degenerate and Kohn Laplacian for $(0, q)$ forms has closed range in $L^{2}$. If the Levi form is degenerate (for example $X$ is weakly pseudoconvex), Hsiao and Marinescu [30, Theorem 1.14] showed that the Szegő projector $S$ is a complex Fourier integral operator on the subset where the Levi form is positive definite if $\bar{\partial}_{b}$ has closed range in $L^{2}$.

Let

$$
\begin{equation*}
\bar{\partial}_{b}^{*}: \operatorname{Dom} \bar{\partial}_{b}^{*} \subset L_{(0,1)}^{2}(X) \rightarrow L^{2}(X) \tag{2.46}
\end{equation*}
$$

be the Hilbert space adjoint of $\bar{\partial}_{b}$ in the $L^{2}$ space with respect to $(\cdot, \cdot)$. Let $\square_{b}$ denote the (Gaffney extension) of the Kohn Laplacian on functions given by

$$
\begin{align*}
\operatorname{Dom} \square_{b} & =\left\{u \in L^{2}(X): u \in \operatorname{Dom} \bar{\partial}_{b}, \bar{\partial}_{b} u \in \operatorname{Dom} \bar{\partial}_{b}^{*}\right\} \\
\square_{b} u & =\bar{\partial}_{b}^{*} \bar{\partial}_{b} u \quad \text { for } u \in \operatorname{Dom} \square_{b}^{(0)} . \tag{2.47}
\end{align*}
$$

By a result of Gaffney, $\square_{b}$ is a positive self-adjoint operator (see [37, Proposition 3.1.2]). In particular, the spectrum $\operatorname{Spec} \square_{b}$ of $\square_{b}$ is contained in $[0, \infty)$. For a Borel set $B \subset \mathbb{R}$ we denote by $E(B)$ the spectral projection of $\square_{b}$ corresponding to the set $B$, where $E$ is the spectral measure of $\square_{b}$. For $\lambda \geq 0$, we set

$$
\begin{equation*}
H_{b, \leq \lambda}^{0}(X):=\operatorname{Im} E((-\infty, \lambda]) \subset L^{2}(X) \tag{2.48}
\end{equation*}
$$

and let

$$
\begin{equation*}
S_{\leq \lambda}: L^{2}(X) \rightarrow H_{b, \leq \lambda}^{0}(X) \tag{2.49}
\end{equation*}
$$

be the orthogonal projection with respect to the product $(\cdot, \cdot)$ and let

$$
\begin{equation*}
S_{\leq \lambda}(x, y) \in \mathscr{D}^{\prime}(X \times X) \tag{2.50}
\end{equation*}
$$

denote the distribution kernel of $S_{\leq \lambda}$. For $\lambda=0$, we write $S:=S_{\leq 0}, S(x, y):=$ $S_{\leq 0}(x, y)$.

Without the assumption that the range of $\bar{\partial}_{b}$ is closed, Ker $\bar{\partial}_{b}$ could be trivial and therefore it is natural to consider the spectral projection $S_{\leq \lambda}$ for $\lambda>0$.

Theorem 2.9 ([30, Theorem 1.5]). For any $\lambda>0$ the spectral projector $S_{\leq \lambda}$ is a complex Fourier integral operator on the subset where the Levi form is positive definite.

Theorems 3.22 and 3.24 are more detailed statements of this result. Since we do not assume that $\bar{\partial}_{b}$ has closed range in $L^{2}$ on $X$, Theorem 2.9 plays an important role in this work. We only assume that the Levi form is non-degenerate on $Y$ and we will show in Theorems 3.17 and 3.19 that the $G$-invariant tangential CR $\bar{\partial}_{b, G}$ has closed range in $L^{2}(X)^{G}$ and we have

$$
\begin{equation*}
S_{G}(x, y)=\int_{G} S_{\leq \lambda_{0}}(x, h \circ y) d \mu(h) \tag{2.51}
\end{equation*}
$$

for some $\lambda_{0}>0$, where $d \mu(h)$ is the Haar measure on $G$ with $\int_{G} d \mu(h)=1$. From (2.51), we can apply Theorem 2.9 to study $S_{G}$ without a closed range in $L^{2}$ assumption on $\bar{\partial}_{b}$ on $X$.

### 2.5. Examples

We give here some simple but non-trivial examples.
(1) Let $X:=\left\{\left(z_{1}, z_{2}, z_{3}\right) \in \mathbb{C}^{3}:\left|z_{1}\right|^{4}+\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}=1\right\}$. Then $X$ is a weakly pseudoconvex CR manifold of dimension five, and $X$ admits a $S^{1}$-action

$$
S^{1} \times X \rightarrow X, \quad e^{i \theta} \cdot\left(z_{1}, z_{2}, z_{3}\right)=\left(e^{-i \theta} z_{1}, e^{i \theta} z_{2}, e^{i \theta} z_{3}\right)
$$

Let $\xi_{X} \in \mathscr{C}^{\infty}(X, T X)$ be the vector field on $X$ induced by $\xi=-\frac{\partial}{\partial \theta}$ in the Lie algebra of $S^{1}$, then

$$
\xi_{X}=-i z_{1} \frac{\partial}{\partial z_{1}}+i \bar{z}_{1} \frac{\partial}{\partial \bar{z}_{1}}+\sum_{j=2}^{3}\left(i z_{j} \frac{\partial}{\partial z_{j}}-i \bar{z}_{j} \frac{\partial}{\partial \bar{z}_{j}}\right)
$$

Let $\omega_{0}:=J(d r)$, where $r:=\left|z_{1}\right|^{4}+\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}-1$ and $J$ is the complex structure map on $T \mathbb{C}^{3}$. It is straightforward to calculate that

$$
\omega_{0}=2 i z_{1} \bar{z}_{1}^{2} d z_{1}-2 i z_{1}^{2} \bar{z}_{1} d \bar{z}_{1}+\sum_{j=2}^{3}\left(i \bar{z}_{j} d z_{j}-i z_{j} d \bar{z}_{j}\right)
$$

and

$$
\left\langle\omega_{0}, \xi_{X}\right\rangle=4\left|z_{1}\right|^{4}-2\left|z_{2}\right|^{2}-2\left|z_{3}\right|^{2} .
$$

Hence,

$$
\mu^{-1}(0)=\left\{\left(z_{1}, z_{2}, z_{3}\right) \in \mathbb{C}^{3}:\left|z_{2}\right|^{2}+\left|z_{3}\right|^{2}=\frac{2}{3},\left|z_{1}\right|^{4}=\frac{1}{3}\right\} .
$$

Thus, $X$ is strictly pseudoconvex near $\mu^{-1}(0)$. Since $X$ is strictly pseudoconvex near $\mu^{-1}(0)$ and $S^{1}$ acts freely on $X$, from Lemma 2.5, zero is a regular value of the moment map $\mu$. Note that $\mu^{-1}(0) / S^{1} \cong S^{3}$. From our main results, we see that modulo some finite-dimensional subspaces of smooth functions, the space of $S^{1}$ invariant CR functions on $X$ is isomorphic to the space of CR functions on $S^{3}$. (2) Let

$$
\begin{aligned}
X:= & \left\{\left(z_{1}, z_{2}, z_{3}, z_{4}, z_{5}, z_{6}\right) \in \mathbb{C}^{6}:\left(\left|z_{5}\right|^{4}+\left|z_{6}\right|^{2}\right)\right. \\
& \left.\times\left(\sum_{j=1}^{4}\left|z_{j}\right|^{2}+z_{1} z_{3}+z_{2} z_{4}+\bar{z}_{1} \bar{z}_{3}+\bar{z}_{2} \bar{z}_{4}\right)=1\right\} .
\end{aligned}
$$

Then, $X$ admits a $G:=S^{1} \times S U(2)$ action:

$$
\begin{aligned}
\left(e^{i \theta}, g\right) \cdot z & =\left(w_{1}, w_{2}, \ldots, w_{6}\right) \\
\left(w_{1}, w_{2}\right)^{t} & :=g\left(z_{1}, z_{2}\right)^{t}, \quad\left(w_{3}, w_{4}\right)^{t}:=\bar{g}\left(z_{3}, z_{4}\right)^{t}, \quad\left(w_{5}, w_{6}\right)=\left(e^{-i \theta} z_{5}, e^{i \theta} z_{6}\right) \\
g & \in S U(2), \quad e^{i \theta} \in S^{1}, \quad z \in X
\end{aligned}
$$

where $z^{t}$ denotes the transpose of $z$. Then $X$ is a weakly pseudoconvex CR manifold and we set $\omega_{0}:=J(d r)$, where

$$
r:=\left(\left|z_{5}\right|^{4}+\left|z_{6}\right|^{2}\right)\left(\sum_{j=1}^{4}\left|z_{j}\right|^{2}+z_{1} z_{3}+z_{2} z_{4}+\bar{z}_{1} \bar{z}_{3}+\bar{z}_{2} \bar{z}_{4}\right)-1
$$

and $J$ is the complex structure map on $T \mathbb{C}^{6}$. As in Example (1) we can check that if $z=\left(z_{1}, z_{2}, z_{3}, z_{4}, z_{5}, z_{6}\right) \in \mu^{-1}(0)$, then $z_{5} \neq 0$. Thus, $X$ is strictly pseudoconvex near $\mu^{-1}(0)$. Since $X$ is strictly pseudoconvex near $\mu^{-1}(0)$ and $S^{1} \times S U(2)$ acts freely on $X$, from Lemma 2.5, zero is a regular value of the moment map $\mu$. Since $X$ is weakly pseudoconvex, it is difficult to understand the Szegő kernel. But from our main results, we see that the $G$-invariant Szegő kernel is a complex Fourier integral operator.

## 3. G-Invariant Szegő Kernel Asymptotics

In this section, we will establish an asymptotic expansion of the $G$-invariant Szegő kernel. From now on, we work under Assumption 1.1 and use the same notations as in Secs. 1 and 2, We do not assume that $\bar{\partial}_{b, X_{G}}$ has closed range in $L^{2}$.

### 3.1. Subelliptic estimates for $G$-invariant smooth functions away $Y$

In this section, we estimate the Szegő kernel outside $Y$. The manifold $X$ is supposed to have arbitrary dimension $\geq 3$. Let

$$
L_{1}, \ldots, L_{N} \in \mathscr{C}^{\infty}(X, H X), \quad N \in \mathbb{N}^{*}
$$

such that for any $x \in X,\left\{L_{1}(x), \ldots, L_{N}(x)\right\} \operatorname{span} H_{x} X$. Let $s \in \mathbb{N}^{*}$. For $u \in$ $\mathscr{C}^{\infty}(X)$, we define

$$
\begin{equation*}
\|u\|\left\|_{s}:=\sum_{\nu=1}^{s} \sum_{1 \leq j_{1}, \ldots, j_{\nu} \leq N}\right\| L_{j_{1}} L_{j_{2}} \ldots L_{j_{\nu}} u\|+\| u \| . \tag{3.1}
\end{equation*}
$$

Theorem 3.1. There exists $C>0$ such that for all $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\left\|\|u\|_{1}^{2} \leq C\left(\left(\square_{b} u, u\right)+|(T u, u)|+\|u\|^{2}\right) .\right. \tag{3.2}
\end{equation*}
$$

Proof. The proof uses the same method as the proof of [11, Theorem 8.3.5], so we only sketch the proof for the convenience of the reader. Let $D \Subset X$ be a small open set and let $\left\{Z_{1}, \ldots, Z_{n}\right\} \in \mathscr{C}^{\infty}\left(D, T^{1,0} X\right)$ be an orthonormal frame of $T^{1,0} X$ on $D$. Let $u \in \mathscr{C}^{\infty}(D)$. We have

$$
\begin{equation*}
\left(\square_{b} u, u\right)=\left\|\bar{\partial}_{b} u\right\|^{2}=\sum_{j=1}^{n}\left\|\bar{Z}_{j} u\right\|^{2} \tag{3.3}
\end{equation*}
$$

For every $j=1, \ldots, n$, by using integration by parts, we have

$$
\begin{align*}
\left\|Z_{j} u\right\|^{2} & =\left(Z_{j} u, Z_{j} u\right)=\left(Z_{j}^{*} Z_{j} u, u\right)=\left(-\bar{Z}_{j} Z_{j} u, u\right)+O\left(\left\|Z_{j} u\right\|\|u\|\right) \\
& =\left(-Z_{j} \bar{Z}_{j} u, u\right)+\left(\left[Z_{j}, \bar{Z}_{j}\right] u, u\right)+O\left(\left\|Z_{j} u\right\|\|u\|\right) \\
& =\left\|\bar{Z}_{j} u\right\|^{2}+\left(\left[Z_{j}, \bar{Z}_{j}\right] u, u\right)+O\left(\left\|Z_{j} u\right\|\|u\|\right)+O\left(\left\|\bar{Z}_{j} u\right\|\|u\|\right) \tag{3.4}
\end{align*}
$$

where $Z_{j}^{*}$ is the formal adjoint of $Z_{j}$. From (3.3) and (3.4), we get (3.2) for every $u \in \mathscr{C}_{c}^{\infty}(D)$. By using a partition of unity we get (3.2) for every $u \in \mathscr{C}^{\infty}(X)$.

Fix $x_{0} \notin Y$. By definition of $Y$, we can find a vector field $V \in \mathscr{C}^{\infty}(X, \underline{\mathfrak{g}})$ such that $\omega_{0}(V) \neq 0$ in an open neighborhood $D$ of $x_{0}$ with $\bar{D} \cap Y=\emptyset$. Then,

$$
\begin{equation*}
L:=V-\omega_{0}(V) T \in H X \tag{3.5}
\end{equation*}
$$

In the rest of this subsection we fix the neighborhood $D$ as above and let

$$
\begin{equation*}
\chi, \widetilde{\chi}, \chi_{1} \in \mathscr{C}_{0}^{\infty}(D), \quad \widetilde{\chi}=1 \text { near } \operatorname{supp} \chi \text { and } \chi_{1}=1 \text { near } \operatorname{supp} \widetilde{\chi} \tag{3.6}
\end{equation*}
$$

Let $u \in \mathscr{C}^{\infty}(X)^{G}$. From (3.5) and since that $V(u)=0$ and $\omega_{0}(V) \neq 0$ on $D$, we have

$$
\begin{align*}
T u & =\frac{-1}{\omega_{0}(V)} L(u) \quad \text { on } D, \\
T(\chi u) & =(T \chi) u+\chi T u=(T \chi) u+\chi \frac{-1}{\omega_{0}(V)} L(u) . \tag{3.7}
\end{align*}
$$

From (3.7), we deduce that there exists $C>0$ such that for all $u \in \mathscr{C}{ }^{\infty}(X)^{G}$,

$$
\begin{equation*}
\|T(\chi u)\| \leq C\left(\|\chi u\|_{1}+\left\|\chi_{1} u\right\|\right) \tag{3.8}
\end{equation*}
$$

For $k \in \mathbb{N}^{*}, U \subset X$ an open set, let $\mathscr{D}^{k}(U)$ be the set of differential operators which can be written as a linear combination of operators as $W_{1} \circ \cdots \circ W_{j}$ with $W_{1}, \ldots, W_{j} \in \mathscr{C}^{\infty}(U, \mathbb{C} H X), 1 \leq j \leq k$.

Lemma 3.2. With the notations above, fix $V_{1}, \ldots, V_{s} \in \mathscr{C}^{\infty}(X, \mathbb{C} H X), s \in \mathbb{N}^{*}$, then there exist $V_{1,1} \in \mathscr{C}^{\infty}(X, \mathbb{C} H X), Q_{1} \in \mathscr{D}^{s-1}(X), Q_{2} \in \mathscr{D}^{s}(X), a, b \in \mathscr{C}^{\infty}(X)$, such that

$$
\begin{align*}
T V_{1} & =V_{1} T+V_{1,1}+a(x) T, \quad \text { if } s=1,  \tag{3.9}\\
T V_{1} \ldots V_{s} & =V_{1} \ldots V_{s} T+Q_{1} T+Q_{2}+b(x) T, \quad \text { if } s \geq 2 \tag{3.10}
\end{align*}
$$

Proof. We first prove (3.9). Note that

$$
\begin{equation*}
T V_{1}=V_{1} T+\left[T, V_{1}\right] \tag{3.11}
\end{equation*}
$$

We have $\left[T, V_{1}\right]=\widetilde{V}_{1,1}+a(x) T$, where $\widetilde{V}_{1,1} \in \mathscr{C}^{\infty}(X, \mathbb{C} H X)$ and $a(x) \in \mathscr{C}^{\infty}(X)$. From this observation and (3.11), we get (3.9).

We now prove (3.10). Let $s=2$. By the argument after (3.11), we have

$$
\begin{align*}
T V_{1} V_{2} & =V_{1} T V_{2}+\left[T, V_{1}\right] V_{2} \\
& =V_{1} T V_{2}+\left(\widetilde{V}_{1,1}+a(x) T\right) V_{2} \tag{3.12}
\end{align*}
$$

From (3.9) and (3.12), we get (3.10) for $s=2$.
Assume that the claim (3.10) holds for $s=s_{0}$ for some $s_{0} \geq 2$. We are going to prove that the claim (3.10) holds for $s=s_{0}+1$. From the argument after (3.11), we have

$$
\begin{align*}
T V_{1} \ldots V_{s_{0}+1} & =V_{1} T V_{2} \ldots V_{s_{0}+1}+\left[T, V_{1}\right] V_{2} \ldots V_{s_{0}+1} \\
& =V_{1} T V_{2} \ldots V_{s_{0}+1}+\left(\widetilde{V}_{1,1}+a(x) T\right) V_{2} \ldots V_{s_{0}+1} \tag{3.13}
\end{align*}
$$

From (3.13) and the induction assumption we get the claim (3.10) for $s=s_{0}+1$. The lemma follows.

Theorem 3.3. Fix $s \in \mathbb{N}^{*}$. Let $V_{1}, \ldots, V_{s} \in \mathscr{C}^{\infty}(D, \mathbb{C} T X)$. Then there exists $Q \in \mathscr{D}^{s}(D)$ such that we have for every $u \in \mathscr{C}^{\infty}(X)^{G}$,

$$
\begin{equation*}
V_{1} \ldots V_{s} u=Q u \quad \text { on } D \tag{3.14}
\end{equation*}
$$

Proof. From (3.7), we see that (3.14) holds for $s=1$. Assume that (3.14) holds for $s=s_{0}$. We are going to prove that (3.14) holds for $s=s_{0}+1$. By induction assumption we only need to assume that $V_{1}=T$ and $V_{j} \in \mathscr{C}^{\infty}(D, \mathbb{C} H X)$,
$j=2,3, \ldots, s_{0}+1$. From (3.9) and (3.10), there exist $\widetilde{V}_{1,1} \in \mathscr{C}^{\infty}(D, \mathbb{C} H X)$, $Q_{1} \in \mathscr{D}^{s-1}(D), Q_{2} \in \mathscr{D}^{s}(D)$ such that

$$
\begin{align*}
T V_{2} & =V_{2} T+\widetilde{V}_{1,1}+a(x) T, \quad \text { if } s_{0}=1  \tag{3.15}\\
T V_{2} \ldots V_{s_{0}+1} & =V_{2} \ldots V_{s_{0}+1} T+Q_{1} T+Q_{2}+b(x) T, \quad \text { if } s_{0} \geq 2
\end{align*}
$$

where $a, b \in \mathscr{C}^{\infty}(D)$. From (3.7) and (3.15), we get (3.14).
From Theorem 3.3 and (3.6) we deduce:
Corollary 3.4. Let $s \in \mathbb{N}^{*}$ and $V_{1}, \ldots, V_{s} \in \mathscr{C}^{\infty}(D, \mathbb{C} T X)$. Then, there exists $C>0$ such that for all $u \in \mathscr{C}^{\infty}(X)^{G}$ we have

$$
\begin{align*}
\left\|\left(V_{1} \ldots V_{s}\right) \chi u\right\| & \leq C\left(\|\chi u\|_{s}+\left\|\chi_{1} u\right\|_{s-1}\right) \\
\left\|\chi\left(V_{1} \ldots V_{s}\right) u\right\| & \leq C\left(\|\chi u\|\left\|_{s}+\right\| \chi_{1} u \|_{s-1}\right) \tag{3.16}
\end{align*}
$$

For $s \in \mathbb{Z}$, let $\|\cdot\|_{s}$ denote the standard Sobolev norm of order $s$ on $X$. From Corollary 3.4 and (3.6), we deduce:

Corollary 3.5. For every $s \in \mathbb{N}$ there exists $C_{s}>0$ such that for any $u \in$ $\mathscr{C}^{\infty}(X)^{G}$,

$$
\|\chi u\|_{s} \leq C_{s}\| \| \chi_{1} u \|_{s}
$$

Theorem 3.6. For every $s \in \mathbb{N}$, there exists $C_{s}>0$ such that for any $u \in \mathscr{C}^{\infty}(X)^{G}$ we have

$$
\begin{equation*}
\left\|\|\chi u\|_{s+1}^{2} \leq C_{s}\left(\| \| \chi \square_{b} u\| \|_{s}^{2}+\| \| \chi_{1} u \|_{s}^{2}\right)\right. \tag{3.17}
\end{equation*}
$$

Proof. We prove (3.17) by induction over $s$. From (3.2), there exists $C>0$ such that for any $u \in \mathscr{C}^{\infty}(X)^{G}$ we have

$$
\begin{equation*}
\left\|\|\chi u\|_{1}^{2} \leq C\left(\left(\square_{b}(\chi u), \chi u\right)+|(T(\chi u), \chi u)|+\|\chi u\|^{2}\right) .\right. \tag{3.18}
\end{equation*}
$$

Now by (3.6),

$$
\begin{equation*}
\left(\square_{b}(\chi u), \chi u\right)=\left(\chi \square_{b} u, \chi u\right)+\left(\chi\left[\square_{b}, \chi\right] u, \chi_{1} u\right) \tag{3.19}
\end{equation*}
$$

From (3.7), (3.18), (3.19), and some elementary computation, we get (3.17) for $s=0$. We now assume that (3.17) holds for every $s<k$ and $k \geq 1$. We will prove that (3.17) holds for $s=k$. Let $Z_{1}, \ldots, Z_{k+1} \in \mathscr{C}^{\infty}(X, \mathbb{C} H X)$. By (3.2), there exist $C_{0}, C>0$ such that for any $u \in \mathscr{C}^{\infty}(X)^{G}$, we have

$$
\begin{align*}
\| Z_{1} \ldots & Z_{k+1}(\chi u) \|^{2} \\
\leq & C_{0}\| \| Z_{2} \ldots Z_{k+1}(\chi u) \|_{1}^{2} \\
\leq & C\left(\left(\square_{b} Z_{2} \ldots Z_{k+1}(\chi u), Z_{2} \ldots Z_{k+1}(\chi u)\right)+\left\|Z_{2} \ldots Z_{k+1}(\chi u)\right\|^{2}\right. \\
& \left.+\left|\left(T Z_{2} \ldots Z_{k+1}(\chi u), Z_{2} \ldots Z_{k+1}(\chi u)\right)\right|\right) . \tag{3.20}
\end{align*}
$$

We have

$$
\begin{align*}
&\left(\square_{b}\left(Z_{2} \ldots Z_{k+1}\right) \chi u,\left(Z_{2} \ldots Z_{k+1}\right) \chi u\right) \\
&=\left(Z_{2} \ldots Z_{k+1} \chi \square_{b} u+\left[\square_{b}, Z_{2} \ldots Z_{k+1}\right] \chi u\right. \\
&\left.+Z_{2} \ldots Z_{k+1}\left[\square_{b}, \chi\right] u, Z_{2} \ldots Z_{k+1} \chi u\right) \\
&=\left(Z_{2} \ldots Z_{k+1} \chi \square_{b} u, Z_{2} \ldots Z_{k+1} \chi u\right) \\
&+\left(\chi_{1} Z_{3} \ldots Z_{k+1}\left[\square_{b}, \chi\right] u, Z_{2}^{*} Z_{2} \ldots Z_{k+1} \chi u\right) \\
&+\left(\left[\square_{b}, Z_{2} \ldots Z_{k+1}\right] \chi u, Z_{2} \ldots Z_{k+1} \chi u\right) \tag{3.21}
\end{align*}
$$

where $Z_{2}^{*}$ denotes the adjoint of $Z_{2}$ and $Z_{2}^{*}=-Z_{2}+$ zero-order term. From (3.16) and (3.21) we see that there exists $C>0$ such that

$$
\begin{align*}
& \left|\left(\square_{b} Z_{2} \ldots Z_{k+1}(\chi u), Z_{2} \ldots Z_{k+1}(\chi u)\right)\right| \\
& \quad \leq C\left(\| \| \square _ { b } u \left\|_{k}^{2}+\frac{1}{\varepsilon}\left|\left\|\mid \chi_{1} u\right\|_{k}^{2}+\varepsilon\| \| \chi u \|_{k+1}^{2}\right)\right.\right. \tag{3.22}
\end{align*}
$$

for every $\varepsilon>0$. Similarly, from (3.16), there exists $\widehat{C}>0$ such that

$$
\begin{equation*}
\left|\left(T Z_{2} \ldots Z_{k+1}(\chi u), Z_{2} \ldots Z_{k+1}(\chi u)\right)\right| \leq \widehat{C}\left(\left.\frac{1}{\varepsilon} \right\rvert\,\left\|\chi_{1} u\right\|_{k}^{2}+\varepsilon\|\chi \chi u\|_{k+1}^{2}\right) \tag{3.23}
\end{equation*}
$$

for every $\varepsilon>0$. From (3.20), (3.22) and (3.23), we conclude that (3.17) holds for $s=k$ by applying (3.17) for $s<k$ for the term $\left\|Z_{2} \ldots Z_{k+1}(\chi u)\right\|^{2}$ in (3.20). The theorem follows.

## From Corollary 3.5 and Theorem 3.6 we get:

Theorem 3.7. For every $s \in \mathbb{N}$, there is $C_{s}>0$ such that for any $u \in \mathscr{C}{ }^{\infty}(X)^{G}$,

$$
\begin{equation*}
\|\chi u\|_{s+1}^{2} \leq C_{s}\left(\left\|\chi_{1} \square_{b} u\right\|_{s}^{2}+\left\|\chi_{1} u\right\|_{s}^{2}\right) \tag{3.24}
\end{equation*}
$$

where $\chi, \chi_{1} \in \mathscr{C}_{0}^{\infty}(D)$ are as in (3.6).

### 3.2. Closed range property for the G-invariant Kohn Laplacian

In this section, we will work in the setting of Assumption 1.1(i). Under these hypotheses we prove subelliptic estimates, regularity and the closed range in $L^{2}$ property for the $G$-invariant Kohn Laplacian.

Let $\bar{\partial}_{b}^{*}: \operatorname{Dom} \bar{\partial}_{b}^{*} \subset L_{(0, q+1)}^{2}(X) \rightarrow L_{(0, q)}^{2}(X)$ be the Hilbert space adjoint of $\bar{\partial}_{b}$ in (2.12) with respect to $(\cdot, \cdot)$ in (2.11). The operators $\bar{\partial}_{b}, \bar{\partial}_{b}^{*}$ commute with
$G$-action, thus we can define $\bar{\partial}_{b, G}$ and $\bar{\partial}_{b, G}^{*}$ with

$$
\begin{equation*}
\operatorname{Dom} \bar{\partial}_{b, G}:=\operatorname{Dom} \bar{\partial}_{b} \cap L_{(0, q)}^{2}(X)^{G}, \quad \operatorname{Dom} \bar{\partial}_{b, G}^{*}:=\operatorname{Dom} \bar{\partial}_{b}^{*} \cap L_{(0, q+1)}^{2}(X)^{G} \tag{3.25}
\end{equation*}
$$

and $\bar{\partial}_{b, G}^{*}: \operatorname{Dom} \bar{\partial}_{b, G}^{*} \rightarrow L_{(0, q)}^{2}(X)^{G}$ is the Hilbert space adjoint of $\bar{\partial}_{b, G}$. Let $\square_{b, G}^{(q)}$ denote the (Gaffney extension) of the $G$-invariant Kohn Laplacian given by

$$
\begin{align*}
\operatorname{Dom} \square_{b, G}^{(q)}= & \left\{u \in L_{(0, q)}^{2}(X)^{G}: u \in \operatorname{Dom} \bar{\partial}_{b, G} \cap \operatorname{Dom} \bar{\partial}_{b, G}^{*},\right. \\
& \left.\bar{\partial}_{b, G} u \in \operatorname{Dom} \bar{\partial}_{b, G}^{*}, \bar{\partial}_{b, G}^{*} u \in \operatorname{Dom} \bar{\partial}_{b, G}\right\}  \tag{3.26}\\
\square_{b, G}^{(q)} u= & \bar{\partial}_{b, G} \bar{\partial}_{b, G}^{*} u+\bar{\partial}_{b, G}^{*} \bar{\partial}_{b, G} u \quad \text { for } u \in \operatorname{Dom} \square_{b, G}^{(q)}
\end{align*}
$$

Lemma 3.8. Let $u \in \operatorname{Dom} \bar{\partial}_{b, G} \cap L_{(0, q)}^{2}(X)^{G}$. Then $\bar{\partial}_{b, G} u \in \operatorname{Dom} \bar{\partial}_{b, G}$ and $\bar{\partial}_{b, G}^{2}$ $u=0$.

Proof. By Friedrichs' lemma [11, Appendix D], there is a sequence $\left\{u_{j}\right\}_{j=1}^{\infty} \subset$ $\Omega^{0, q}(X)^{G}$ such that $u_{j} \rightarrow u$ in $L_{(0, q)}^{2}(X)^{G}$ as $j \rightarrow \infty$ and $\bar{\partial}_{b, G} u_{j} \rightarrow \bar{\partial}_{b, G} u$ in $L_{(0, q+1)}^{2}(X)^{G}$ as $j \rightarrow \infty$. Let $v \in \Omega^{0, q+2}(X)^{G}$. We have

$$
\begin{equation*}
\left(\bar{\partial}_{b, G} u, \bar{\partial}_{b, G}^{*} v\right)=\lim _{j \rightarrow \infty}\left(\bar{\partial}_{b, G} u_{j}, \bar{\partial}_{b, G}^{*} v\right)=\lim _{j \rightarrow \infty}\left(\bar{\partial}_{b, G}^{2} u_{j}, v\right)=0 \tag{3.27}
\end{equation*}
$$

Hence, $\bar{\partial}_{b, G} u \in \operatorname{Dom} \bar{\partial}_{b, G}$ and $\bar{\partial}_{b, G}^{2} u=0$.
Lemma 3.9. The operator $\square_{b, G}^{(q)}: \operatorname{Dom} \square_{b, G}^{(q)} \subset L_{(0, q)}^{2}(X)^{G} \rightarrow L_{(0, q)}^{2}(X)^{G}$ is closed.
Proof. Let $\left\{\left(f_{k}, \square_{b, G}^{(q)} f_{k}\right) \in L_{(0, q)}^{2}(X)^{G} \times L_{(0, q)}^{2}(X)^{G}: f_{k} \in \operatorname{Dom} \square_{b, G}^{(q)}\right\}_{k=1}^{\infty}$ with

$$
\begin{equation*}
\lim _{k \rightarrow \infty} f_{k}=f, \quad \lim _{k \rightarrow \infty} \square_{b, G}^{(q)} f_{k}=h \quad \text { in } L_{(0, q)}^{2}(X)^{G} \tag{3.28}
\end{equation*}
$$

By definition, to check that $\square_{b, G}^{(q)}$ is a closed operator, we need to show that $f \in$ $\operatorname{Dom} \square_{b, G}^{(q)}$ and $\square_{b, G}^{(q)} f=h$. Since $f_{k} \in \operatorname{Dom} \square_{b, G}^{(q)}$, for each $k$, we have

$$
\begin{align*}
\left\|\bar{\partial}_{b, G}^{*}\left(f_{j}-f_{k}\right)\right\|^{2} & \leq\left\|\bar{\partial}_{b, G}^{*}\left(f_{j}-f_{k}\right)\right\|^{2}+\left\|\bar{\partial}_{b, G}\left(f_{j}-f_{k}\right)\right\|^{2} \\
& =\left(\square_{b, G}^{(q)}\left(f_{j}-f_{k}\right), f_{j}-f_{k}\right) . \tag{3.29}
\end{align*}
$$

From (3.28) and (3.29), $\left\{\bar{\partial}_{b, G}^{*} f_{k}\right\}_{k=1}^{\infty}$ is a Cauchy sequence in $L^{2}$, hence $\lim _{k \rightarrow \infty} \bar{\partial}_{b, G}^{*} f_{k}=h_{1}$ in $L_{(0, q-1)}^{2}(X)^{G}$, for some $h_{1} \in L_{(0, q-1)}^{2}(X)^{G}$.

Let $v \in \operatorname{Dom} \bar{\partial}_{b, G} \cap L_{(0, q-1)}^{2}(X)^{G}$. We have

$$
\begin{equation*}
\left(f, \bar{\partial}_{b, G} v\right)=\lim _{k \rightarrow \infty}\left(f_{k}, \bar{\partial}_{b, G} v\right)=\lim _{k \rightarrow \infty}\left(\bar{\partial}_{b, G}^{*} f_{k}, v\right)=\left(h_{1}, v\right) \tag{3.30}
\end{equation*}
$$

Hence, $f \in \operatorname{Dom} \bar{\partial}_{b, G}^{*}$ and $\bar{\partial}_{b, G}^{*} f=\lim _{k \rightarrow \infty} \bar{\partial}_{b, G}^{*} f_{k}$. Similarly, we can repeat the procedure above and show that $f \in \operatorname{Dom} \bar{\partial}_{b, G}$ and $\bar{\partial}_{b, G} f=\lim _{k \rightarrow \infty} \bar{\partial}_{b, G} f_{k}$.

Now, we show that $\bar{\partial}_{b, G} f \in \operatorname{Dom} \bar{\partial}_{b, G}^{*}$. From Lemma 3.8, we know that

$$
\begin{equation*}
\left(\bar{\partial}_{b, G}^{*} \bar{\partial}_{b, G} u, \bar{\partial}_{b, G} \bar{\partial}_{b, G}^{*} v\right)=0, \text { for every } u, v \in \operatorname{Dom} \square_{b, G}^{(q)} \tag{3.31}
\end{equation*}
$$

From (3.31), we have

$$
\begin{equation*}
\left\|\bar{\partial}_{b, G}^{*} \bar{\partial}_{b, G}\left(f_{j}-f_{k}\right)\right\|^{2}+\left\|\bar{\partial}_{b, G} \bar{\partial}_{b, G}^{*}\left(f_{j}-f_{k}\right)\right\|^{2}=\left\|\square_{b, G}^{(q)}\left(f_{j}-f_{k}\right)\right\|^{2} \tag{3.32}
\end{equation*}
$$

From (3.28), (3.32), $\left\{\bar{\partial}_{b, G}^{*} \bar{\partial}_{b, G} f_{j}\right\}_{j=1}^{\infty}$ is a Cauchy sequence in $L^{2}$, hence $\lim _{k \rightarrow \infty} \bar{\partial}_{b, G}^{*} \bar{\partial}_{b, G} f_{k}=h_{2}$ in $L_{(0, q)}^{2}(X)^{G}$, for some $h_{2} \in L_{(0, q)}^{2}(X)^{G}$ and thus, $\bar{\partial}_{b, G} f \in \operatorname{Dom} \bar{\partial}_{b, G}^{*}$ and $\bar{\partial}_{b, G}^{*} \bar{\partial}_{b, G} f=\lim _{k \rightarrow \infty} \bar{\partial}_{b, G}^{*} \bar{\partial}_{b, G} f_{k}$. Similarly, we can repeat the process above and show that

$$
\begin{equation*}
\bar{\partial}_{b, G}^{*} f \in \operatorname{Dom} \bar{\partial}_{b, G}, \quad \bar{\partial}_{b, G} \bar{\partial}_{b, G}^{*} f=\lim _{k \rightarrow \infty} \bar{\partial}_{b, G} \bar{\partial}_{b, G}^{*} f_{k} . \tag{3.33}
\end{equation*}
$$

Hence, $f \in \operatorname{Dom} \square_{b, G}^{(q)}$ and $\square_{b, G}^{(q)} f=\lim _{k \rightarrow \infty} \square_{b, G}^{(q)} f_{k}=h$. The lemma follows.
Since the dimension of $X$ is greater or equal to five, we can repeat Kohn's method [32] (see also the proof of [11, Theorem 8.3.5]) and deduce the following subelliptic estimates:

Theorem 3.10. Under Assumption 1.1(i), let $\eta, \eta_{1} \in \mathscr{C}^{\infty}(X)$ such that $\eta=1$ near $Y, \eta_{1}=1$ near $\operatorname{supp} \eta$ and the Levi form is positive near supp $\eta_{1}$. Then for every $s \in \mathbb{N}$ there is $C_{s}>0$ such that for any $u \in \Omega^{0,1}(X)$,

$$
\begin{equation*}
\|\eta u\|_{s+1}^{2} \leq C_{s}\left(\left\|\eta_{1} \square_{b}^{(1)} u\right\|_{s}^{2}+\left\|\eta_{1} u\right\|_{s}^{2}\right) \tag{3.34}
\end{equation*}
$$

Repeating the proof of Theorem 3.7 with minor changes we get:
Theorem 3.11. Let $\gamma, \gamma_{1} \in \mathscr{C}^{\infty}(X)$ with $\gamma_{1}=1$ near supp $\gamma$ and $\operatorname{supp} \gamma_{1} \cap Y=\emptyset$. For every $s \in \mathbb{N}$, there exists $C_{s}>0$ such that for any $u \in \Omega^{0,1}(X)^{G}$,

$$
\begin{equation*}
\|\gamma u\|_{s+1}^{2} \leq C_{s}\left(\left\|\gamma_{1} \square_{b, G}^{(1)} u\right\|_{s}^{2}+\left\|\gamma_{1} u\right\|_{s}^{2}\right) . \tag{3.35}
\end{equation*}
$$

From Theorems 3.10, 3.11 and by using a partition of unity we obtain:
Theorem 3.12. Assume that Assumption $1.1(\mathrm{i})$ holds. Then for every $s \in \mathbb{N}$ and every $\gamma, \gamma_{1} \in \mathscr{C}^{\infty}(X)$ with $\gamma_{1}=1$ near $\operatorname{supp} \gamma$, there is $C_{s}>0$ such that for any $u \in \Omega^{0,1}(X)^{G}$,

$$
\begin{equation*}
\|\gamma u\|_{s+1}^{2} \leq C_{s}\left(\left\|\gamma_{1} \square_{b, G}^{(1)} u\right\|_{s}^{2}+\left\|\gamma_{1} u\right\|_{s}^{2}\right) . \tag{3.36}
\end{equation*}
$$

For every $s \in \mathbb{Z}$, let $H_{(0, q)}^{s}(X)^{G}$ be the completion of $\Omega^{0, q}(X)^{G}$ with respect to $\|\cdot\|_{s}$. From Theorem 3.12, we can repeat the technique of elliptic regularization (see the proof of [11, Theorem 8.4.2]) and conclude:

Theorem 3.13. Under Assumption 1.1(i) let $u \in \operatorname{Dom} \square_{b, G}^{(1)}$ and let $\mathcal{U}$ be an open set of $X$. Let $\square_{b, G}^{(1)} u=v \in L_{(0,1)}^{2}(X)^{G}$. If $\left.v\right|_{\mathcal{U}}$ is smooth, then $\left.u\right|_{\mathcal{U}}$ is smooth.

Let $\tau, \tau_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\tau_{1}=1$ near $\operatorname{supp} \tau$. If $\tau_{1} v \in H_{(0,1)}^{s}(X)^{G}$, for some $s \in \mathbb{N}$, then $\tau u \in H_{(0,1)}^{s+1}(X)^{G}$ and there is $C_{s}>0$ independent of $u$, $v$, such that

$$
\begin{equation*}
\|\tau u\|_{s+1} \leq C_{s}\left(\left\|\tau_{1} \square_{b, G}^{(1)} u\right\|_{s}+\left\|\tau_{1} u\right\|_{s}\right) \tag{3.37}
\end{equation*}
$$

Theorem 3.14. If Assumption 1.1(i) holds, then the operator $\square_{b, G}^{(1)}: \operatorname{Dom} \square_{b, G}^{(1)} \subset$ $L_{(0,1)}^{2}(X)^{G} \rightarrow L_{(0,1)}^{2}(X)^{G}$ has closed range.

Proof. We claim that there is $c>0$ such that

$$
\begin{equation*}
\left\|\square_{b, G}^{(1)} u\right\| \geq c\|u\|, \text { for all } u \in \operatorname{Dom} \square_{b, G}^{(1)}, u \perp \operatorname{Ker} \square_{b, G}^{(1)} \text {. } \tag{3.38}
\end{equation*}
$$

Suppose that the claim is not true. We can find $u_{j} \in \operatorname{Dom} \square_{b, G}^{(1)}, u_{j} \perp \operatorname{Ker} \square_{b, G}^{(1)}$, $\left\|u_{j}\right\|=1, j=1,2, \ldots$, such that

$$
\begin{equation*}
\left\|\square_{b, G}^{(1)} u_{j}\right\| \leq \frac{1}{j}\left\|u_{j}\right\|, \quad j=1,2, \ldots \tag{3.39}
\end{equation*}
$$

From Theorem 3.13, (3.39) and Rellich's lemma in functional analysis, we can find $\left\{j_{s}\right\}_{s=1}^{\infty} \subset \mathbb{N}, 1 \leq j_{1}<j_{2}<\ldots, \lim _{s \rightarrow \infty} j_{s}=\infty$, such that $\lim _{s \rightarrow \infty}\left\|u_{j_{s}}-u\right\|=0$, for some $u \in L_{(0,1)}^{2}(X)^{G}$. It is obvious that $\|u\|=1$ and $u \perp \operatorname{Ker} \square_{b, G}^{(1)}$. Let $g \in$ Dom $\bar{\partial}_{b, G} \bigcap L^{2}(X)^{G}$. We have

$$
\begin{aligned}
\left|\left(u, \bar{\partial}_{b, G} g\right)\right| & =\lim _{s \rightarrow \infty}\left|\left(u_{j_{s}}, \bar{\partial}_{b, G} g\right)\right|=\lim _{s \rightarrow \infty}\left|\left(\bar{\partial}_{b, G}^{*} u_{j_{s}}, g\right)\right| \\
& \leq \lim _{s \rightarrow \infty}\left\|\bar{\partial}_{b, G}^{*} u_{j_{s}}\right\|\|g\| \leq \lim _{s \rightarrow \infty}\left\|\square_{b, G}^{(1)} u_{j_{s}}\right\| \frac{1}{2}\|g\| \leq \lim _{s \rightarrow \infty} \frac{1}{\sqrt{j_{s}}}\|g\|=0
\end{aligned}
$$

Hence, $u \in \operatorname{Dom} \bar{\partial}_{b, G}^{*}$ and $\bar{\partial}_{b, G}^{*} u=0$. Let $f \in L_{(0,2)}^{2}(X)^{G}$. We have

$$
\begin{aligned}
\left|\left(u, \bar{\partial}_{b, G}^{*} f\right)\right| & =\lim _{s \rightarrow \infty}\left|\left(u_{j_{s}}, \bar{\partial}_{b, G}^{*} f\right)\right|=\lim _{s \rightarrow \infty}\left|\left(\bar{\partial}_{b, G} u_{j_{s}}, f\right)\right| \\
& \leq \lim _{s \rightarrow \infty}\left\|\bar{\partial}_{b, G} u_{j_{s}}\right\|\|f\| \leq \lim _{s \rightarrow \infty}\left\|\square_{b, G}^{(1)} u_{j_{s}}\right\| \frac{1}{2}\|f\| \leq \lim _{s \rightarrow \infty} \frac{1}{\sqrt{j_{s}}}\|f\|=0 .
\end{aligned}
$$

Hence, $u \in \operatorname{Dom} \bar{\partial}_{b, G}$ and $\bar{\partial}_{b, G} u=0$. We have proved that $u \in \operatorname{Ker} \square_{b, G}^{(1)}$. But $u \perp \operatorname{Ker} \square_{b, G}^{(1)}$, we get a contradiction. The claim (3.38) follows. From (3.38), we get the lemma.

Let $N_{G}^{(1)}: L_{(0,1)}^{2}(X)^{G} \rightarrow \operatorname{Dom} \square_{b, G}^{(1)}$ be the partial inverse of $\square_{b, G}^{(1)}$ and let $S_{G}^{(1)}: L_{(0,1)}^{2}(X)^{G} \rightarrow \operatorname{Ker} \square_{b, G}^{(1)}$ be the Szegő projection, i.e. the orthogonal projection onto $\operatorname{Ker} \square_{b, G}^{(1)}$ with respect to $(\cdot, \cdot)$. We have

$$
\begin{align*}
& \square_{b, G}^{(1)} N_{G}^{(1)}+S_{G}^{(1)}=I \quad \text { on } L_{(0,1)}^{2}(X)^{G},  \tag{3.40}\\
& N_{G}^{(1)} \square_{b, G}^{(1)}+S_{G}^{(1)}=I \quad \text { on } \operatorname{Dom} \square_{b, G}^{(1)} .
\end{align*}
$$

From Theorem 3.13 we deduce:
Theorem 3.15. Assume that Assumption 1.1(i) holds. Then for every $s \in \mathbb{Z}$ we can extend $N_{G}^{(1)}$ to a continuous operator $N_{G}^{(1)}: H_{(0,1)}^{s}(X)^{G} \rightarrow H_{(0,1)}^{s+1}(X)$. Moreover, $\operatorname{Ker} \square_{b, G}^{(1)}$ is a finite-dimensional subspace of $\Omega^{0,1}(X)^{G}$.

Let

$$
\begin{equation*}
\mathcal{P}_{G}^{(q)}: L_{(0, q)}^{2}(X) \rightarrow L_{(0, q)}^{2}(X)^{G}, \tag{3.41}
\end{equation*}
$$

be the orthogonal projection with respect to $(\cdot, \cdot)$. It is not difficult to see that for every $s \in \mathbb{Z}$ we can extend $\mathcal{P}_{G}^{(q)}$ to $H_{(0, q)}^{s}(X)$ and $\mathcal{P}_{G}^{(q)}: H_{(0, q)}^{s}(X) \rightarrow H_{(0, q)}^{s}(X)^{G}$ is continuous. We extend $N_{G}^{(1)}$ and $S_{G}^{(1)}$ to $H_{(0,1)}^{s}(X)$ by

$$
\begin{equation*}
N_{G}^{(1)} u:=N_{G}^{(1)} \mathcal{P}_{G}^{(1)} u, \quad S_{G}^{(1)} u:=S_{G}^{(1)} \mathcal{P}_{G}^{(1)} u, \quad \text { for } u \in H_{(0,1)}^{s}(X), \quad s \in \mathbb{Z} \tag{3.42}
\end{equation*}
$$

From Theorem 3.15 we see that $S_{G}^{(1)}$ is smoothing and

$$
\begin{equation*}
N_{G}^{(1)}: H_{(0,1)}^{s}(X) \rightarrow H_{(0,1)}^{s+1}(X)^{G} \text { is continuous for every } s \in \mathbb{Z} \tag{3.43}
\end{equation*}
$$

Theorem 3.16. Under Assumption 1.1(i), let $\tau, \tau_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\operatorname{supp} \tau \cap$ $\operatorname{supp} \tau_{1}=\emptyset$. Then $\tau N_{G}^{(1)} \tau_{1}$ is smoothing.

Proof. Let $\widetilde{\tau}, \hat{\tau} \in \mathscr{C}^{\infty}(X)^{G}$ with $\widetilde{\tau}=1$ near $\operatorname{supp} \hat{\tau}, \hat{\tau}=1$ near $\operatorname{supp} \tau$ and $\operatorname{supp} \widetilde{\tau} \cap$ $\operatorname{supp} \tau_{1}=\emptyset$. Let $v \in L_{(0,1)}^{2}(X)^{G}$ and put $\widetilde{\tau} N_{G}^{(1)} \tau_{1} v=u \in H_{(0,1)}^{1}(X)^{G}$. From (3.40), we have

$$
\begin{align*}
\square_{b, G}^{(1)} u & =\square_{b, G}^{(1)} \widetilde{\tau} N_{G}^{(1)} \tau_{1} v \\
& =\widetilde{\tau} \square_{b, G}^{(1)} N_{G}^{(1)} \tau_{1} v+\left[\square_{b, G}^{(1)}, \widetilde{\tau}\right] N_{G}^{(1)} \tau_{1} v \\
& =\widetilde{\tau}\left(I-S_{G}^{(1)}\right) \tau_{1} v+\left[\square_{b, G}^{(1)}, \widetilde{\tau}\right] N_{G}^{(1)} \tau_{1} v \\
& =-\widetilde{\tau} S_{G}^{(1)} \tau_{1} v+\left[\square_{b, G}^{(1)}, \widetilde{\tau}\right] N_{G}^{(1)} \tau_{1} v . \tag{3.44}
\end{align*}
$$

Since $S_{G}^{(1)}$ is smoothing, $-\widetilde{\tau} S_{G}^{(1)} \tau_{1} v \in \mathscr{C}^{\infty}(X)^{G}$. Since $\widetilde{\tau}=1$ near $\operatorname{supp} \hat{\tau}$, $\hat{\tau}\left[\square_{b, G}^{(1)}, \widetilde{\tau}\right] N_{G}^{(1)} \tau_{1} v=0$. From this observation, we deduce that

$$
\begin{equation*}
\hat{\tau} \square_{b, G}^{(1)} u \in \mathscr{C}^{\infty}(X)^{G} \tag{3.45}
\end{equation*}
$$

Fix $s \in \mathbb{N}$. From (3.37) and (3.44), there exist $C_{s}, \widetilde{C}_{s}>0$ such that for any $v \in$ $L_{(0,1)}^{2}(X)^{G}$, we have

$$
\begin{align*}
\left\|\tau N_{G}^{(1)} \tau_{1} v\right\|_{s+1} & =\|\tau u\|_{s+1} \leq C_{s}\left(\left\|\hat{\tau} \square_{b, G}^{(1)} u\right\|_{s}+\|\hat{\tau} u\|_{s}\right) \\
& \leq \widetilde{C}_{s}\left(\left\|-\hat{\tau} S_{G}^{(1)} \tau v\right\|_{s}+\left\|\hat{\tau} N_{G}^{(1)} \tau_{1} v\right\|_{s}\right) . \tag{3.46}
\end{align*}
$$

Take $s=1$ in (3.46), from Theorem 3.15 and note that $S_{G}^{(1)}$ is smoothing, we conclude that $\left\|\tau N_{G}^{(1)} \tau_{1} v\right\|_{2} \leq C\|v\|$. We have proved that for any $\gamma, \gamma_{1} \in \mathscr{C}^{\infty}(X)^{G}$
with supp $\gamma \cap \operatorname{supp} \gamma_{1}=\emptyset$, we have

$$
\begin{equation*}
\gamma N_{G}^{(1)} \gamma_{1}: L_{(0,1)}^{2}(X)^{G} \rightarrow H_{(0,1)}^{2}(X)^{G} \text { is continuous. } \tag{3.47}
\end{equation*}
$$

Take $s=2$ in (3.46), from (3.47) and note that $S_{G}^{(1)}$ is smoothing, we conclude that there exists $\hat{C}>0$ such that for any $v \in L_{(0,1)}^{2}(X)^{G},\left\|\tau N_{G}^{(1)} \tau_{1} v\right\|_{3} \leq \hat{C}\|v\|$. Continuing in this way, we conclude that for any $\gamma, \gamma_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\operatorname{supp} \gamma \cap$ $\operatorname{supp} \gamma_{1}=\emptyset$, we have

$$
\begin{equation*}
\gamma N_{G}^{(1)} \gamma_{1}: L_{(0,1)}^{2}(X)^{G} \rightarrow H_{(0,1)}^{s}(X)^{G} \text { is continuous, for every } s \in \mathbb{N}^{*} . \tag{3.48}
\end{equation*}
$$

By taking adjoint in (3.48), we deduce that for any $\gamma, \gamma_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\operatorname{supp} \gamma \cap$ $\operatorname{supp} \gamma_{1}=\emptyset$, we have

$$
\begin{equation*}
\gamma N_{G}^{(1)} \gamma_{1}: H_{(0,1)}^{-s}(X)^{G} \rightarrow L_{(0,1)}^{2}(X)^{G} \text { is continuous, for every } s \in \mathbb{N}^{*} \tag{3.49}
\end{equation*}
$$

Now, let $v \in H_{(0,1)}^{-s_{0}}(X)^{G}, s_{0} \in \mathbb{N}^{*}$, and put $\widetilde{\tau} N_{G}^{(1)} \tau_{1} v=u \in H_{(0,1)}^{-s_{0}+1}(X)^{G}$. Let $v_{j} \in \Omega^{0,1}(X)^{G}, j=1,2, \ldots$, with $v_{j} \rightarrow v$ in $H_{(0,1)}^{-s_{0}}(X)^{G}$ as $j \rightarrow \infty$. Taking $s=0$ in (3.46), we deduce from (3.49) and the fact that $S_{G}^{(1)}$ is smoothing that there exists $C>0$ such that for any $h \in L_{(0,1)}^{2}(X)^{G}$,

$$
\begin{equation*}
\left\|\tau N_{G}^{(1)} \tau_{1} h\right\|_{1} \leq C\|h\|_{-s_{0}} . \tag{3.50}
\end{equation*}
$$

From (3.50), we have

$$
\begin{equation*}
\lim _{j, k \rightarrow \infty}\left\|\tau N_{G}^{(1)} \tau_{1}\left(v_{j}-v_{k}\right)\right\|_{1} \leq C\left\|v_{j}-v_{k}\right\|_{-s_{0}}=0 \tag{3.51}
\end{equation*}
$$

From (3.51), we see that $\left\{\tau N_{G}^{(1)} \tau_{1} v_{j}\right\}_{j \geq 1}$ is a Cauchy sequence in $H_{(0,1)}^{1}(X)^{G}$. Since $\tau N_{G}^{(1)} \tau_{1} v_{j} \rightarrow \tau N_{G}^{(1)} \tau_{1} v$ in $H_{(0,1)}^{-s_{0}+1}(X)^{G}$ as $j \rightarrow \infty$ (see Theorem 3.15), we deduce that $\tau N_{G}^{(1)} \tau_{1} v \in H_{(0,1)}^{1}(X)^{G}$ and $\left\|\tau N_{G}^{(1)} \tau_{1} v\right\|_{1} \leq C\|v\|_{-s_{0}}$. We have proved that for any $\gamma, \gamma_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with

$$
\operatorname{supp} \gamma \cap \operatorname{supp} \gamma_{1}=\emptyset
$$

we have

$$
\begin{equation*}
\gamma N_{G}^{(1)} \gamma_{1}: H_{(0,1)}^{-s_{0}}(X)^{G} \rightarrow H_{(0,1)}^{1}(X)^{G} \text { is continuous. } \tag{3.52}
\end{equation*}
$$

Again, take $s=1$ in (3.46), from (3.52) and note that $S_{G}^{(1)}$ is smoothing, we conclude that there exists $C>0$ such that for any $h \in L_{(0,1)}^{2}(X)^{G}$,

$$
\begin{equation*}
\left\|\tau N_{G}^{(1)} \tau_{1} h\right\|_{2} \leq C\|h\|_{-s_{0}} . \tag{3.53}
\end{equation*}
$$

From (3.53), we can repeat the argument above and deduce that $\tau N_{G}^{(1)} \tau_{1} v \in$ $H_{(0,1)}^{2}(X)^{G}$ and $\left\|\tau N_{G}^{(1)} \tau_{1} v\right\|_{2} \leq C\|v\|_{-s_{0}}$. Continuing in this way, we conclude that for any $\gamma, \gamma_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with supp $\gamma \cap \operatorname{supp} \gamma_{1}=\emptyset$, we have

$$
\begin{equation*}
\gamma N_{G}^{(1)} \gamma_{1}: H_{(0,1)}^{-s_{0}}(X)^{G} \rightarrow H_{(0,1)}^{s}(X)^{G} \text { is continuous, for every } s \in \mathbb{N} . \tag{3.54}
\end{equation*}
$$

The theorem follows.

We return to the case of functions. As before, let $S_{G}: L^{2}(X) \rightarrow \operatorname{Ker} \square_{b, G}^{(0)}=$ $H_{b}^{0}(X)^{G}$ be the $G$-invariant Szegő projection. In view of Theorem 3.14 we can repeat the proof of [27, Proposition 6.15, p. 56] and deduce the following.

Theorem 3.17. Under Assumption 1.1(i) the operator $\square_{b, G}^{(0)}$ : $\operatorname{Dom} \square_{b, G}^{(0)} \subset$ $L^{2}(X)^{G} \rightarrow L^{2}(X)^{G}$ has closed range and

$$
\begin{equation*}
S_{G}=\mathcal{P}_{G}^{(0)}-\bar{\partial}_{b, G}^{*} N_{G}^{(1)} \bar{\partial}_{b, G} \mathcal{P}_{G}^{(0)} \text { on } L^{2}(X) \tag{3.55}
\end{equation*}
$$

Theorem 3.18. Let $\tau, \tau_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\operatorname{supp} \tau \cap \operatorname{supp} \tau_{1}=\emptyset$. If Assumption $1.1(\mathrm{i})$ holds, then $\tau S_{G} \tau_{1}$ is smoothing.

Proof. By (3.41) and (3.55) we have

$$
\begin{equation*}
\tau S_{G} \tau_{1}=-\tau \bar{\partial}_{b, G}^{*} N_{G}^{(1)} \bar{\partial}_{b, G} \tau_{1}=-\tau \bar{\partial}_{b, G}^{*} \widetilde{\tau} N_{G}^{(1)} \widetilde{\tau}_{1} \bar{\partial}_{b, G} \tau_{1}, \tag{3.56}
\end{equation*}
$$

where $\widetilde{\tau}, \widetilde{\tau}_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\widetilde{\tau}=1$ near $\operatorname{supp} \tau, \widetilde{\tau}_{1}=1$ near $\operatorname{supp} \tau_{1}, \operatorname{supp} \widetilde{\tau} \cap$ $\operatorname{supp} \widetilde{\tau}_{1}=\emptyset$. In view of Theorem 3.16, we see that $\widetilde{\tau} N_{G}^{(1)} \widetilde{\tau}_{1}$ is smoothing. From this observation and (3.56), the theorem follows.

Let $d \mu=d \mu(h)$ be the Haar measure on $G$ with $\int_{G} d \mu(h)=1$. We also need the following:

Theorem 3.19. Under Assumption 1.1(i) there exists $c_{0}>0$ such that for any $\lambda \in\left(0, c_{0}\right)$ such that

$$
S_{G}(x, y)=\int_{G} S_{\leq \lambda}(x, h . y) d \mu(h) \quad \text { on } X \times X
$$

where $S_{\leq \lambda}$ is the spectral projection given by (2.49).

Proof. Since $\square_{b, G}^{(0)}$ : $\operatorname{Dom} \square_{b, G}^{(0)} \subset L^{2}(X)^{G} \rightarrow L^{2}(X)^{G}$ has closed range, we can define the partial inverse

$$
N_{G}^{(0)}: L^{2}(X)^{G} \rightarrow \operatorname{Dom} \square_{b, G}^{(0)}
$$

of $\square_{b, G}^{(0)}$ as follows: Let $u \in L^{2}(X)^{G}$. Since $\operatorname{Im} \square_{b, G}^{(0)}$ is closed in $L^{2}(X)^{G}$, we have the orthogonal decomposition

$$
u=v_{0}+v_{1}, \quad v_{0} \in \operatorname{Im} \square_{b, G}^{(0)}, \quad v_{1} \perp \operatorname{Im} \square_{b, G}^{(0)} .
$$

There is a unique $\beta \in \operatorname{Dom} \square_{b, G}^{(0)}$, $\beta \perp \operatorname{Ker} \square_{b, G}^{(0)}$, such that $\square_{b, G}^{(0)} \beta=v_{0}$. Define $N_{G}^{(0)} u:=\beta$. Since $\square_{b, G}^{(0)}$ is self-adjoint, $v_{1} \in \operatorname{Ker} \square_{b, G}^{(0)}$. Thus, we have the Hodge
decomposition:

$$
\begin{aligned}
& \square_{b, G}^{(0)} N_{G}^{(0)}+S_{G}=I \text { on } L^{2}(X)^{G}, \\
& N_{G}^{(0)} \square_{b, G}^{(0)}+S_{G}=I \text { on } \operatorname{Dom} \square_{b, G}^{(0)} .
\end{aligned}
$$

Then, $N_{G}^{(0)}: L^{2}(X)^{G} \rightarrow \operatorname{Dom} \square_{b, G}^{(0)}$ is a linear operator. We claim that

$$
\begin{equation*}
N_{G}^{(0)}: L^{2}(X)^{G} \rightarrow \operatorname{Dom} \square_{b, G}^{(0)} \text { is a closed operator. } \tag{3.57}
\end{equation*}
$$

Let $\left\{\left(f_{k}, N_{G}^{(0)} f_{k}\right) \in L^{2}(X)^{G} \times L^{2}(X)^{G}: k \in \mathbb{N}\right\}$, with $\lim _{k \rightarrow \infty} f_{k}=f$ in $L^{2}(X)^{G}$ and $\lim _{k \rightarrow \infty} N_{G}^{(0)} f_{k}=h$ in $L^{2}(X)^{G}$. For every $k \in \mathbb{N}$, write

$$
\begin{equation*}
f_{k}=\square_{b, G}^{(0)} g_{k}+\xi_{k}, \quad g_{k} \in \operatorname{Dom} \square_{b, G}^{(0)}, \quad g_{k} \perp \operatorname{Ker} \square_{b, G}^{(0)}, \quad \xi_{k} \perp \operatorname{Im} \square_{b, G}^{(0)} \tag{3.58}
\end{equation*}
$$

From (3.58) and the definition of $N_{G}^{(0)}$, we see that

$$
\begin{equation*}
N_{G}^{(0)} f_{k}=g_{k}, \quad \lim _{k \rightarrow \infty} g_{k}=h \quad \text { in } L^{2}(X)^{G} \tag{3.59}
\end{equation*}
$$

Since $\left\{f_{k}\right\}_{k=1}^{\infty}$ is a Cauchy sequence, $\left\{\square_{b, G}^{(0)} g_{k}\right\}_{k=1}^{\infty}$ and $\left\{\xi_{k}\right\}_{k=1}^{\infty}$ are Cauchy sequences. Hence,

$$
\begin{equation*}
\lim _{k \rightarrow \infty} \square_{b, G}^{(0)} g_{k}=\eta, \quad \lim _{k \rightarrow \infty} \xi_{k}=\xi \quad \text { in } L^{2}(X)^{G} \tag{3.60}
\end{equation*}
$$

for some $\eta, \xi \in L^{2}(X)^{G}, \xi \perp \operatorname{Im} \square_{b, G}^{(0)}$, and we have the orthogonal decomposition

$$
\begin{equation*}
f=\eta+\xi \tag{3.61}
\end{equation*}
$$

From (3.59) and (3.60), we see that $\left(g_{k}, \square_{b, G}^{(0)} g_{k}\right) \rightarrow(h, \eta)$ in $L^{2}(X)^{G} \times L^{2}(X)^{G}$ as $k \rightarrow \infty$. Since $\square_{b, G}^{(0)}$ is a closed operator, we conclude that $h \in \operatorname{Dom} \square_{b, G}^{(0)}, h \perp$ $\operatorname{Ker} \square_{b, G}^{(0)}$, and $\square_{b, G}^{(0)} h=\eta$. From this observation and (3.61), we get the orthogonal decomposition

$$
f=\square_{b, G}^{(0)} h+\xi
$$

and hence $N_{G}^{(0)} f=h$. The claim (3.57) follows.
Since $N_{G}^{(0)}$ is a closed operator defined on the Banach space $L^{2}(X)^{G}$, by the closed graph theorem, $N_{G}^{(0)}$ is a continuous operator. Hence, there is $c_{0}>0$ such that

$$
\begin{equation*}
\left\|N_{G}^{(0)} \beta\right\| \leq \frac{1}{c_{0}}\|\beta\|, \quad \text { for every } \beta \in L^{2}(X)^{G} \tag{3.62}
\end{equation*}
$$

Let $\beta=\square_{b, G}^{(0)} u$ in (3.62), where $u \in \operatorname{Dom} \square_{b, G}^{(0)}, u \perp \operatorname{Ker} \square_{b, G}^{(0)}$, we get

$$
\begin{equation*}
\left\|\square_{b, G}^{(0)} u\right\| \geq c_{0}\|u\|, \quad \text { for every } u \in \operatorname{Dom} \square_{b, G}^{(0)}, u \perp \operatorname{Ker} \square_{b, G}^{(0)} . \tag{3.63}
\end{equation*}
$$

Fix $0<\lambda<c_{0}$. We claim that with the notation (3.41), we have

$$
\begin{equation*}
S_{G}=S_{\leq \lambda} \circ \mathcal{P}_{G}^{(0)} \quad \text { on } L^{2}(X) \tag{3.64}
\end{equation*}
$$

If the claim is not true, we can find a $u \in H_{b, \leq \lambda}^{0}(X) \cap L^{2}(X)^{G}$ with $u \perp \operatorname{Ker} \square_{b, G}^{(0)}$, $\|u\|=1$, where $H_{b, \leq \lambda}^{0}(X)$ is given by (2.48). Since $u \in H_{b, \leq \lambda}^{0}(X)$, we have
$\left\|\square_{b, G}^{(0)} u\right\| \leq \lambda\|u\|<c_{0}\|u\|$. From this observation and (3.63) we get a contradiction. The claim (3.64) follows and this yields the theorem.

### 3.3. G-invariant Szegö kernel asymptotics away $Y$

Let $S_{G}(x, y) \in \mathscr{D}^{\prime}(X \times X)$ be the distribution kernel of $S_{G}$. For any subset $A$ of $X$, put $G A:=\{g \cdot x: x \in A, g \in G\}$.

Theorem 3.20. Under Assumption 1.1, let $\tau \in \mathscr{C}{ }^{\infty}(X)$ with $\operatorname{supp} \tau \cap Y=\emptyset$. Then $\tau S_{G}$ and $S_{G} \tau$ are smoothing operators.

Proof. First we work under Assumption 1.1(i) as in the previous section. Let $v \in L^{2}(X)$. Take $v_{j} \in \mathscr{C}^{\infty}(X), j \in \mathbb{N}^{*}$, such that $\left\|v_{j}-v\right\| \rightarrow 0$ as $j \rightarrow \infty$. We have $\left\|S_{G} v_{j}-S_{G} v\right\| \rightarrow 0$ as $j \rightarrow \infty$. By (3.43) and (3.55), we see that

$$
\begin{equation*}
S_{G} v_{j} \in \mathscr{C}^{\infty}(X), \quad \text { for every } j \in \mathbb{N}^{*} \tag{3.65}
\end{equation*}
$$

For every $j$, put $u_{j}:=S_{G} v_{j}$. By Corollary 3.5 and (3.17), it is straightforward to see that for every $s \in \mathbb{N}$ there exists $C_{s}>0$ such that for $\chi$ in (3.6),

$$
\begin{equation*}
\left\|\chi u_{j}\right\|_{s} \leq C_{s}, \quad \text { for every } j \in \mathbb{N}^{*} \tag{3.66}
\end{equation*}
$$

From (3.66) we deduce that $\chi S_{G} v \in H^{s}(X)$ for every $s \in \mathbb{N}$ and

$$
\begin{equation*}
\chi S_{G}: L^{2}(X) \rightarrow H^{s}(X) \text { is continuous, for every } s \in \mathbb{N} . \tag{3.67}
\end{equation*}
$$

By using a partition of unity we conclude that for any $\tau, \tau_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\operatorname{supp} \tau \cap Y=\emptyset, \operatorname{supp} \tau_{1} \cap Y=\emptyset$, we have

$$
\begin{equation*}
\tau S_{G}: L^{2}(X) \rightarrow H^{s}(X) \text { is continuous, for every } s \in \mathbb{N} \tag{3.68}
\end{equation*}
$$

and hence

$$
\begin{equation*}
S_{G} \tau_{1}: H^{-s}(X) \rightarrow L^{2}(X) \text { is continuous, for every } s \in \mathbb{N} . \tag{3.69}
\end{equation*}
$$

From (3.68) and (3.69), we conclude that

$$
\begin{equation*}
\tau S_{G} \tau_{1}=\left(\tau S_{G}\right) \circ\left(S_{G} \tau_{1}\right): H^{-s}(X) \rightarrow H^{s}(X) \text { is continuous, for every } s \in \mathbb{N}, \tag{3.70}
\end{equation*}
$$

and hence $\tau S_{G} \tau_{1}$ is smoothing.
Since $G$ acts on $Y$, it is not difficult to see that there is a small neighborhood $W$ of $Y$ such that $\operatorname{supp} \tau \cap G W=\emptyset$. Let $\gamma_{0}, \gamma_{1} \in \mathscr{C}^{\infty}(X)^{G} \cap \mathscr{C}_{0}^{\infty}(G W)$ with $\gamma_{0}=1$ near supp $\gamma_{1}, \gamma_{1}=1$ near $Y$. For $i=0,1$, put $\tau_{i}:=1-\gamma_{i}$; then $\tau_{i} \in \mathscr{C}^{\infty}(X)^{G}$ with
$\operatorname{supp} \tau_{i} \cap Y=\emptyset$. Moreover, it is straightforward to check that

$$
\begin{equation*}
\tau_{0}=1 \text { on } \operatorname{supp} \tau, \quad \operatorname{supp} \tau_{0} \cap \operatorname{supp}\left(1-\tau_{1}\right)=\emptyset \tag{3.71}
\end{equation*}
$$

From (3.71) follows

$$
\begin{equation*}
\tau S_{G}=\tau \tau_{0} S_{G}=\tau \tau_{0} S_{G} \tau_{1}+\tau \tau_{0} S_{G}\left(1-\tau_{1}\right) \tag{3.72}
\end{equation*}
$$

By (3.70) $\tau_{0} S_{G} \tau_{1}$ is smoothing. In view of Theorem 3.18 and (3.71), we see that

$$
\begin{equation*}
\tau \tau_{0} S_{G}\left(1-\tau_{1}\right) \text { is smoothing. } \tag{3.73}
\end{equation*}
$$

From (3.72) and (3.73), we get that $\tau S_{G}$ is smoothing and hence $S_{G} \tau$ is smoothing.
Let us now work under the Assumption 1.1(ii). By [8, [30, Theorem 1.14] the Szegő projector $S$ is a Fourier integral operator with complex phase on $X$. We have $S_{G}=S \circ \mathcal{P}_{G}^{(0)}$. Since $S, \mathcal{P}_{G}^{(0)}$ map smooth functions to smooth functions, $S_{G}$ maps smooth functions to smooth functions. Thus, (3.65) still holds. Now,

$$
\begin{equation*}
\tau \tau_{0} S_{G}\left(1-\tau_{1}\right)=\tau \tau_{0} S \circ \mathcal{P}_{G}^{(0)}\left(1-\tau_{1}\right)=\tau \tau_{0} S\left(1-\tau_{1}\right) \circ \mathcal{P}_{G}^{(0)} \tag{3.74}
\end{equation*}
$$

Since $S$ is smoothing away the diagonal, $\tau \tau_{0} S \circ\left(1-\tau_{1}\right)$ is smoothing. From this observation and (3.74), we conclude that $\tau \tau_{0} S_{G}\left(1-\tau_{1}\right)$ is smoothing. Thus, (3.73) still holds. Since (3.65) and (3.73) hold so by repeating the proof above we conclude also in this case.

### 3.4. G-invariant Szegő kernel asymptotics near $Y$

In this section, $X$ has arbitrary dimension $\geq 3$. We first recall the definition of the Hörmander symbol spaces. Let $D \subset X$ be a local coordinate patch with local coordinates $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$.

Definition 3.21. For $m \in \mathbb{R}, S_{1,0}^{m}\left(D \times D \times \mathbb{R}_{+}\right)$is the space of all $a(x, y, t) \in$ $\mathscr{C}^{\infty}\left(D \times D \times \mathbb{R}_{+}\right)$such that, for all compact $K \Subset D \times D$ and all $\alpha, \beta \in \mathbb{N}^{2 n+1}$, $\gamma \in \mathbb{N}$, there exists $C_{\alpha, \beta, \gamma}>0$ such that

$$
\left|\partial_{x}^{\alpha} \partial_{y}^{\beta} \partial_{t}^{\gamma} a(x, y, t)\right| \leq C_{\alpha, \beta, \gamma}(1+|t|)^{m-\gamma}, \quad(x, y, t) \in K \times \mathbb{R}_{+}, \quad t \geq 1
$$

Put

$$
S^{-\infty}\left(D \times D \times \mathbb{R}_{+}\right):=\bigcap_{m \in \mathbb{R}} S_{1,0}^{m}\left(D \times D \times \mathbb{R}_{+}\right)
$$

Let $a_{j} \in S_{1,0}^{m_{j}}\left(D \times D \times \mathbb{R}_{+}\right), j=0,1,2, \ldots$ with $m_{j} \searrow-\infty$, as $j \rightarrow \infty$. Then there exists $a \in S_{1,0}^{m_{0}}\left(D \times D \times \mathbb{R}_{+}\right)$unique modulo $S^{-\infty}$, such that

$$
a-\sum_{j=0}^{k-1} a_{j} \in S_{1,0}^{m_{k}}\left(D \times D \times \mathbb{R}_{+}\right) \quad \text { for all } k \in \mathbb{N}^{*}
$$

If $a$ and $a_{j}$ have the properties above, we write $a \sim \sum_{j=0}^{\infty} a_{j}$ in $S_{1,0}^{m_{0}}\left(D \times D \times \mathbb{R}_{+}\right)$. We write

$$
\begin{equation*}
s(x, y, t) \in S_{\mathrm{cl}}^{m}\left(D \times D \times \mathbb{R}_{+}\right) \tag{3.75}
\end{equation*}
$$

if $s(x, y, t) \in S_{1,0}^{m}\left(D \times D \times \mathbb{R}_{+}\right)$and

$$
\begin{align*}
& s(x, y, t) \sim \sum_{j=0}^{\infty} s_{j}(x, y) t^{m-j} \quad \text { in } S_{1,0}^{m}\left(D \times D \times \mathbb{R}_{+}\right)  \tag{3.76}\\
& s_{j}(x, y) \in \mathscr{C}^{\infty}(D \times D), \quad \text { for any } j \in \mathbb{N}
\end{align*}
$$

Let $W_{1} \subset \mathbb{R}^{N_{1}}, W_{2} \subset \mathbb{R}^{N_{2}}$ be open sets. We can also define $S_{1,0}^{m}\left(W_{1} \times W_{2} \times \mathbb{R}_{+}\right)$, $S_{\mathrm{cl}}^{m}\left(W_{1} \times W_{2} \times \mathbb{R}_{+}\right)$and asymptotic sum in the similar way.

By Theorem [2.9 (cf. [30, Theorem 1.5]) the spectral projector $S_{\leq \lambda}$ is for every $\lambda>0$ a complex Fourier integral operator on the subset where the Levi form is positive definite. We give here a detailed description of the spectral kernel.

Theorem 3.22 ([30, Theorem 4.7]). Fix $\lambda>0$. Let $D \subset X$ be a local coordinate patch with local coordinates $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$. Assume that the Levi form of $X$ is positive definite at every point of $D$. Then,

$$
S_{\leq \lambda}(x, y) \equiv \int_{0}^{\infty} e^{i \varphi(x, y) t} s(x, y, t) d t \quad \text { on } D \times D
$$

with a symbol $s(x, y, t) \in S_{\mathrm{cl}}^{n}\left(D \times D \times \mathbb{R}_{+}\right)$such that the coefficient $s_{0}$ of the expansion (3.76) is given by

$$
\begin{equation*}
s_{0}(x, x)=\frac{1}{2} \pi^{-n-1}\left|\operatorname{det} \mathscr{L}_{x}\right|, \quad x \in D \tag{3.77}
\end{equation*}
$$

where $\operatorname{det} \mathscr{L}_{x}$ is the determinant of the Levi form, see (2.10), and the phase function $\varphi$ satisfies

$$
\begin{align*}
\varphi & \in \mathscr{C}^{\infty}(D \times D), \quad \operatorname{Im} \varphi(x, y) \geq 0 \\
\varphi(x, x) & =0, \quad \varphi(x, y) \neq 0 \quad \text { if } x \neq y \\
\left.d_{x} \varphi(x, y)\right|_{x=y} & =-\left.d_{y} \varphi(x, y)\right|_{x=y}=\lambda(x) \omega_{0}(x), \quad \lambda(x)>0  \tag{3.78}\\
\varphi(x, y) & =-\bar{\varphi}(y, x)
\end{align*}
$$

Moreover, let $X^{\prime}:=\{x \in X$ : the Levi form is non-degenerate at $x\}$. Then, $S_{\leq \lambda}$ is smoothing away the diagonal on $X^{\prime} \times X^{\prime}$.

Remark 3.23. With the same notations used in Theorem 3.22, let

$$
\Lambda_{\varphi}:=\left\{\left(x, t d_{x} \varphi, y, t d_{y} \varphi\right) \in T^{*} D \times T^{*} D: \varphi(x, y)=0, t>0\right\}
$$

From (3.78), we see that

$$
\Lambda_{\varphi}=\left\{\left(x, \lambda \omega_{0}(x), x,-\lambda \omega_{0}(x)\right) \in T^{*} D \times T^{*} D: x \in D, \lambda>0\right\}
$$

Hence, the canonical relation $C_{\varphi}$ of $S_{\leq \lambda}$ is given by

$$
C_{\varphi}=\left\{\left(x, \lambda \omega_{0}(x), x, \lambda \omega_{0}(x)\right) \in T^{*} D \times T^{*} D: x \in D, \lambda>0\right\}
$$

The following result describes the phase function in local coordinates (see [27, Chap. 8 of part I $]$ ). For a given point $p \in D$, let $\left\{W_{j}\right\}_{j=1}^{n}$ be an orthonormal frame
of $T^{1,0} X$ in a neighborhood of $p$ such that the Levi form is diagonal at $p$, i.e. $\mathscr{L}_{p}\left(W_{j}, \bar{W}_{s}\right)=\delta_{j, s} \mu_{j}, j, s=1, \ldots, n$, where $\delta_{j, s}=1$ if $j=s, \delta_{j, s}=0$ if $j \neq s$. We take local coordinates $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ defined on some neighborhood of $p$ such that $\omega_{0}(p)=-d x_{2 n+1}, x(p)=0$ and

$$
\begin{align*}
W_{j}= & \frac{\partial}{\partial z_{j}}-i \mu_{j} \bar{z}_{j} \frac{\partial}{\partial x_{2 n+1}}-c_{j} x_{2 n+1} \frac{\partial}{\partial x_{2 n+1}} \\
& +\sum_{k=1}^{2 n} a_{j, k}(x) \frac{\partial}{\partial x_{k}}+O\left(|x|^{2}\right), \quad j=1, \ldots, n \tag{3.79}
\end{align*}
$$

where $z_{j}=x_{2 j-1}+i x_{2 j}, c_{j} \in \mathbb{C}, a_{j, k}(x)$ is $\mathscr{C}^{\infty}, a_{j, k}(x)=O(|x|)$, for every $j=$ $1, \ldots, n, k=1, \ldots, 2 n$. Set $y=\left(y_{1}, \ldots, y_{2 n+1}\right), w_{j}=y_{2 j-1}+i y_{2 j}, j=1, \ldots, n$.

Theorem 3.24. With the same notations and assumptions used in Theorem 3.22 we have for the phase function $\varphi$ in some neighborhood of $(0,0)$,

$$
\begin{align*}
\operatorname{Im} \varphi(x, y) \geq & c \sum_{j=1}^{2 n}\left|x_{j}-y_{j}\right|^{2}, \quad c>0  \tag{3.80}\\
\varphi(x, y)= & -x_{2 n+1}+y_{2 n+1}+i \sum_{j=1}^{n}\left|\mu_{j}\right|\left|z_{j}-w_{j}\right|^{2} \\
& +\sum_{j=1}^{n}\left(i \mu_{j}\left(\bar{z}_{j} w_{j}-z_{j} \bar{w}_{j}\right)+c_{j}\left(-z_{j} x_{2 n+1}+w_{j} y_{2 n+1}\right)\right. \\
& \left.+\bar{c}_{j}\left(-\bar{z}_{j} x_{2 n+1}+\bar{w}_{j} y_{2 n+1}\right)\right)+\left(x_{2 n+1}-y_{2 n+1}\right) f(x, y) \\
& +O\left(|(x, y)|^{3}\right) \tag{3.81}
\end{align*}
$$

where $f$ is smooth and satisfies $f(0,0)=0, f(x, y)=\bar{f}(y, x)$. Moreover, we can take the phase $\varphi$ so that

$$
\begin{equation*}
\bar{\partial}_{b, x} \varphi(x, y) \text { vanishes to infinite order at } x=y \text {. } \tag{3.82}
\end{equation*}
$$

Furthermore, for any $\varphi_{1}(x, y) \in \mathscr{C}^{\infty}(D \times D)$, if $\varphi_{1}$ satisfies (3.78), (3.80), (3.81) and (3.82), then there is a function $h(x, y) \in \mathscr{C}^{\infty}(D \times D)$ with $h(x, x) \neq 0$, for every $x \in D$, such that $\varphi(x, y)-h(x, y) \varphi_{1}(x, y)$ vanishes to infinite order at $x=y$.

For the next result we recall that the map $\mathcal{R}_{x}$ and the function $V_{\text {eff }}$ were defined in (1.18) and (1.19). We denote by $\lambda_{1}^{\prime}(x), \ldots, \lambda_{d}^{\prime}(x)$ the eigenvalues of $\mathcal{R}_{x}$ with respect to the $G$-invariant Hermitian metric $g$ and we define the determinant of $\mathcal{R}_{x}$ by

$$
\begin{equation*}
\operatorname{det} \mathcal{R}_{x}=\lambda_{1}^{\prime}(x) \ldots \lambda_{d}^{\prime}(x) \tag{3.83}
\end{equation*}
$$

Theorem 3.25. Under the assumptions of Theorem 1.2, let $p \in Y$, let $U$ be an open neighborhood of $p$ and let $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ be local coordinates defined in $U$. Then,

$$
\begin{equation*}
S_{G}(x, y) \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} a(x, y, t) d t \quad \text { on } U \times U \tag{3.84}
\end{equation*}
$$

with a symbol $a(x, y, t) \in S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(U \times U \times \mathbb{R}_{+}\right)$such that the coefficient $a_{0}$ in its expansion (3.76) satisfies

$$
\begin{equation*}
a_{0}(x, x)=2^{d-1} \frac{1}{V_{\mathrm{eff}}(x)} \pi^{-n-1+\frac{d}{2}}\left|\operatorname{det} \mathcal{R}_{x}\right|^{-\frac{1}{2}}\left|\operatorname{det} \mathscr{L}_{x}\right|, \quad x \in U \cap Y \tag{3.85}
\end{equation*}
$$

and the phase function $\Phi$ satisfies

$$
\begin{align*}
\Phi(x, y) & \in \mathscr{C}^{\infty}(U \times U), \quad \operatorname{Im} \Phi(x, y) \geq 0, \\
d_{x} \Phi(x, x) & =-d_{y} \Phi(x, x)=\lambda(x) \omega_{0}(x), \quad \lambda(x)>0, \quad x \in U \cap Y . \tag{3.86}
\end{align*}
$$

Moreover, there exists $C \geq 1$ such that for all $(x, y) \in U \times U$,

$$
\begin{align*}
|\Phi(x, y)|+\operatorname{Im} \Phi(x, y) & \leq C\left(\inf \left\{d^{2}(g \cdot x, y): g \in G\right\}+d^{2}(x, Y)+d^{2}(y, Y)\right) \\
|\Phi(x, y)|+\operatorname{Im} \Phi(x, y) & \geq \frac{1}{C}\left(\inf \left\{d^{2}(g \cdot x, y): g \in G\right\}+d^{2}(x, Y)+d^{2}(y, Y)\right),  \tag{3.87}\\
C d^{2}(x, Y) & \geq \operatorname{Im} \Phi(x, x) \geq \frac{1}{C} d^{2}(x, Y), \quad x \in U
\end{align*}
$$

and $\Phi(x, y)$ satisfies (3.93), (3.94) and (3.95) below.
Proof. If Assumption 1.1(i) holds, then by Theorem 3.20 we can localize the study of the $G$-invariant Szegő kernel $S_{G}$ to $Y$ and from Theorems 3.19and 3.22, we repeat the proof of [28, Theorem 1.5] and conclude. If Assumption 1.1(ii) holds, we know by [8], 30, Theorem 1.14] that the Szegő projector $S$ is a Fourier integral operator with complex phase on $X$. Repeating the argument from [28, Theorem 1.5] we conclude.

Remark 3.26. With the same notations used in Theorem 3.25, let

$$
\Lambda_{\Phi}:=\left\{\left(x, t d_{x} \Phi, y, t d_{y} \Phi\right) \in T^{*} X \times T^{*} X: \Phi(x, y)=0, t>0\right\}
$$

From (3.87), we see that

$$
\Lambda_{\Phi}=\left\{\left(x, \lambda \omega_{0}(x), g \cdot x,-\lambda \omega_{0}(x)\right) \in T^{*} X \times T^{*} X: x \in Y, g \in G, \lambda>0\right\}
$$

Hence, the canonical relation $C_{\Phi}$ of $S_{G}$ is given by

$$
C_{\Phi}=\left\{\left(x, \lambda \omega_{0}(x), g \cdot x, \lambda \omega_{0}(x)\right) \in T^{*} X \times T^{*} X: x \in Y, g \in G, \lambda>0\right\}
$$

As applications of Theorems 3.18, 3.20 and [3.25, we establish the following regularity property for $S_{G}$.

Theorem 3.27. Under the assumptions of Theorem 1.2 we have

$$
S_{G}: \mathscr{C}^{\infty}(X) \rightarrow H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)
$$

In particular, $H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)$ is dense in $H_{b}^{0}(X)^{G}$.
Proof. Let $U$ be an open coordinate patch of $X$ and let $u \in \mathscr{C}_{0}^{\infty}(U)$. If $U \cap Y=\emptyset$, we see in view of Theorem 3.20 that $S_{G} u \in \mathscr{C}^{\infty}(X)$. Assume now $U \cap Y \neq \emptyset$. By

Theorem $3.25 S_{G}$ is a Fourier integral operator with complex phase on $U$ and hence $S_{G} u \in \mathscr{C}^{\infty}(U)$. We only need to show that $S_{G} u$ is smooth outside $U$. Let $x_{0} \notin U$. If $x_{0} \notin Y$, by Theorem 3.20 again we deduce that $S_{G} u$ is smooth near $x_{0}$. Now, we suppose that $x_{0} \in Y$.
Case I: $G x_{0} \cap U=\emptyset$. We can find $\tau, \tau_{1} \in \mathscr{C}^{\infty}(X)^{G}$ with $\operatorname{supp} \tau \cap \operatorname{supp} \tau_{1}=\emptyset, \tau \equiv 1$ near $x_{0}, \tau_{1} \equiv 1$ near $\operatorname{supp} u$. We have $\tau S_{G} u=\tau S_{G} \tau_{1} u$. In view of Theorem 3.18 under Assumption 1.1(i) or by the fact that $S_{G}$ is a Fourier integral operator under Assumption 1.1(ii), we see that $\tau S_{G} \tau_{1}$ is smoothing and we deduce that $\tau S_{G} \tau_{1} u \in$ $\mathscr{C}^{\infty}(X)$. In particular, $S_{G} u$ is smooth near $x_{0}$.
Case II: $G x_{0} \cap U \neq \emptyset$. There is a $\hat{g} \in G$ such that $\hat{g} \cdot x_{0} \in U$. Since $S_{G} u \in \mathscr{C}^{\infty}(U)$, $S_{G} u$ is smooth near $\hat{g} \cdot x_{0}$. Since $S_{G} u$ is $G$-invariant, $S_{G} u$ is smooth near $x_{0}$.

We have thus proved that $S_{G} u \in \mathscr{C}^{\infty}(X)$. By using a partition of unity we conclude.

Let $e_{0}$ be the identity element in $G$. Fix $p \in Y$. It was shown in [28, Theorem 3.6] that there exist local coordinates $v=\left(v_{1}, \ldots, v_{d}\right)$ on $G$ defined in a neighborhood $V$ of $e_{0}$ with $v\left(e_{0}\right)=(0, \ldots, 0)$ (until further notice, we will identify the element $h \in V$ with $v(h))$, local coordinates $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ of $X$ defined in a neighborhood $U=U_{1} \times U_{2}$ of $p$ with $0 \leftrightarrow p$, where $U_{1} \subset \mathbb{R}^{d}$ is a neighborhood of $0 \in \mathbb{R}^{d}$, $U_{2} \subset \mathbb{R}^{2 n+1-d}$ is an open neighborhood of $0 \in \mathbb{R}^{2 n+1-d}$ and a smooth function $\gamma=\left(\gamma_{1}, \ldots, \gamma_{d}\right) \in \mathscr{C}^{\infty}\left(U_{2}, U_{1}\right)$ with $\gamma(0)=0 \in \mathbb{R}^{d}$ such that for $\left(v_{1}, \ldots, v_{d}\right) \in V$, $\left(x_{d+1}, \ldots, x_{2 n+1}\right) \in U_{2}$,
$\left(v_{1}, \ldots, v_{d}\right) \cdot\left(\gamma\left(x_{d+1}, \ldots, x_{2 n+1}\right), x_{d+1}, \ldots, x_{2 n+1}\right)$
$=\left(v_{1}+\gamma_{1}\left(x_{d+1}, \ldots, x_{2 n+1}\right), \ldots, v_{d}+\gamma_{d}\left(x_{d+1}, \ldots, x_{2 n+1}\right), x_{d+1}, \ldots, x_{2 n+1}\right)$,
$Y \cap U=\left\{x_{d+1}=\ldots=x_{2 d}=0\right\}, \quad \underline{\mathfrak{g}}=\operatorname{span}\left\{\frac{\partial}{\partial x_{1}}, \ldots, \frac{\partial}{\partial x_{d}}\right\}$,
$J\left(\frac{\partial}{\partial x_{j}}\right)=\frac{\partial}{\partial x_{d+j}}+a_{j}(x) \frac{\partial}{\partial x_{2 n+1}}, \quad$ on $Y \cap U$, for $1 \leq j \leq d$,
where $a_{j}(x)$ are smooth functions on $Y \cap U$, independent of $x_{1}, \ldots, x_{2 d}, x_{2 n+1}$, $a_{j}(0)=0$, and $T_{p}^{1,0} X=\operatorname{span}\left\{Z_{1}, \ldots, Z_{n}\right\}$ with

$$
\begin{align*}
Z_{j} & =\frac{1}{2}\left(\frac{\partial}{\partial x_{j}}-i \frac{\partial}{\partial x_{d+j}}\right)(p), \quad \text { for } j=1, \ldots, d, \\
Z_{j} & =\frac{1}{2}\left(\frac{\partial}{\partial x_{2 j-1}}-i \frac{\partial}{\partial x_{2 j}}\right)(p), \quad \text { for } j=d+1, \ldots, n,  \tag{3.89}\\
\left\langle Z_{j}, Z_{k}\right\rangle_{g} & =\delta_{j, k}, \quad \mathscr{L}_{p}\left(Z_{j}, \bar{Z}_{k}\right)=\mu_{j} \delta_{j, k}, \quad \text { for } j, k=1,2, \ldots, n, \\
-\omega_{0}(x)= & (1+O(|x|)) d x_{2 n+1}+\sum_{j=1}^{d} 4 \mu_{j} x_{d+j} d x_{j}+\sum_{j=d+1}^{n} 2 \mu_{j} x_{2 j} d x_{2 j-1} \\
- & \sum_{j=d+1}^{n} 2 \mu_{j} x_{2 j-1} d x_{2 j}+\sum_{j=d+1}^{2 n} b_{j} x_{2 n+1} d x_{j}+O\left(|x|^{2}\right), \tag{3.90}
\end{align*}
$$

where $b_{d+1}, \ldots, b_{2 n} \in \mathbb{R}$. Put

$$
\begin{equation*}
x^{\prime \prime}=\left(x_{d+1}, \ldots, x_{2 n+1}\right), \quad \widehat{x}^{\prime \prime}=\left(x_{d+1}, \ldots, x_{2 d}\right), \quad \dot{x}^{\prime \prime}=\left(x_{d+1}, \ldots, x_{2 n}\right) . \tag{3.91}
\end{equation*}
$$

Theorem 3.28 ([28, Theorem 1.11]). The phase function $\Phi(x, y)$ appearing in the expression of the Szegő kernel (3.84), (3.86) is independent of $\left(x_{1}, \ldots, x_{d}\right)$ and $\left(y_{1}, \ldots, y_{d}\right)$. Hence,

$$
\begin{equation*}
\Phi(x, y)=\Phi\left(\left(0, x^{\prime \prime}\right),\left(0, y^{\prime \prime}\right)\right)=: \Phi\left(x^{\prime \prime}, y^{\prime \prime}\right) \tag{3.92}
\end{equation*}
$$

Moreover, there exists $c>0$ such that

$$
\begin{equation*}
\operatorname{Im} \Phi\left(x^{\prime \prime}, y^{\prime \prime}\right) \geq c\left(\left|\widehat{x}^{\prime \prime}\right|^{2}+\left|\widehat{y}^{\prime \prime}\right|^{2}+\left|\grave{x}^{\prime \prime}-\dot{y}^{\prime \prime}\right|^{2}\right), \quad \text { for }\left(0, x^{\prime \prime}\right),\left(0, y^{\prime \prime}\right) \in U \tag{3.93}
\end{equation*}
$$

and there exists a smooth function $g(x, y) \in \mathscr{C}^{\infty}\left(U \times U, T^{* 0,1} X\right)$ such that

$$
\begin{align*}
& \bar{\partial}_{b, x} \Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)-g(x, y) \Phi\left(x^{\prime \prime}, y^{\prime \prime}\right) \text { vanishes to } \\
& \quad \text { infinite order on } \operatorname{diag}((Y \cap U) \times(Y \cap U)), \tag{3.94}
\end{align*}
$$

and with the same $\mu_{j}, b_{d+1}, \ldots, b_{2 n} \in \mathbb{R}$ as in (3.90) we have

$$
\begin{align*}
\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)= & -x_{2 n+1}+y_{2 n+1}+2 i \sum_{j=1}^{d}\left|\mu_{j}\right| y_{d+j}^{2}+2 i \sum_{j=1}^{d}\left|\mu_{j}\right| x_{d+j}^{2} \\
& +i \sum_{j=d+1}^{n}\left|\mu_{j}\right|\left|z_{j}-w_{j}\right|^{2}+\sum_{j=d+1}^{n} i \mu_{j}\left(\bar{z}_{j} w_{j}-z_{j} \bar{w}_{j}\right) \\
& +\sum_{j=1}^{d}\left(-b_{d+j} x_{d+j} x_{2 n+1}+b_{d+j} y_{d+j} y_{2 n+1}\right) \\
& +\sum_{j=d+1}^{n} \frac{1}{2}\left(b_{2 j-1}-i b_{2 j}\right)\left(-z_{j} x_{2 n+1}+w_{j} y_{2 n+1}\right) \\
& +\sum_{j=d+1}^{n} \frac{1}{2}\left(b_{2 j-1}+i b_{2 j}\right)\left(-\bar{z}_{j} x_{2 n+1}+\bar{w}_{j} y_{2 n+1}\right) \\
& +\left(-x_{2 n+1}+y_{2 n+1}\right) f\left(x^{\prime \prime}, y^{\prime \prime}\right)+O\left(\left|\left(x^{\prime \prime}, y^{\prime \prime}\right)\right|^{3}\right), \tag{3.95}
\end{align*}
$$

where $z_{j}=x_{2 j-1}+i x_{2 j}, w_{j}=y_{2 j-1}+i y_{2 j}$, for $j=d+1, \ldots, n$, and $f\left(x^{\prime \prime}, y^{\prime \prime}\right) \in$ $\mathscr{C}^{\infty}(U \times U)$ with $f(0,0)=0$.

Remark 3.29. The phase function $\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)$ is not unique. For example, we can replace $\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)$ by $\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right) r\left(x^{\prime \prime}, y^{\prime \prime}\right)$, where $r\left(x^{\prime \prime}, y^{\prime \prime}\right) \in \mathscr{C}^{\infty}(U \times U), r(0,0)=1$. In [28. Theorem 5.2], the first author and Huang characterized the phase $\Phi$. Since $\partial_{y_{2 n+1}} \Phi(0,0) \neq 0$, the Malgrange preparation theorem [26, Theorem 7.57] implies that

$$
\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)=g\left(x^{\prime \prime}, y^{\prime \prime}\right)\left(y_{2 n+1}+\widehat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right)
$$

in some neighborhood of $(0,0)$, where $g\left(x^{\prime \prime}, y^{\prime \prime}\right), \widehat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right) \in \mathscr{C}^{\infty}(U \times U)$. We can replace $\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)$ by $y_{2 n+1}+\widehat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)$. From now on, we assume that

$$
\begin{equation*}
\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)=y_{2 n+1}+\widehat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right), \quad \widehat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right) \in \mathscr{C}^{\infty}(U \times U) . \tag{3.96}
\end{equation*}
$$

It is straightforward to check that the phase $\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)$ satisfies (3.86), (3.87), (3.93), and

$$
\begin{equation*}
\bar{\partial}_{b, x} \Phi\left(x^{\prime \prime}, y^{\prime \prime}\right) \text { vanishes to infinite order at } \operatorname{diag}((Y \cap U) \times(Y \cap U)) \tag{3.97}
\end{equation*}
$$ and with the same notations as in (3.95) we have

$$
\begin{align*}
\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right)= & -x_{2 n+1}+y_{2 n+1}+2 i \sum_{j=1}^{d}\left|\mu_{j}\right| y_{d+j}^{2}+2 i \sum_{j=1}^{d}\left|\mu_{j}\right| x_{d+j}^{2} \\
& +i \sum_{j=d+1}^{n}\left|\mu_{j}\right|\left|z_{j}-w_{j}\right|^{2}+\sum_{j=d+1}^{n} i \mu_{j}\left(\bar{z}_{j} w_{j}-z_{j} \bar{w}_{j}\right) \\
& +\sum_{j=1}^{d}\left(-b_{d+j} x_{d+j} x_{2 n+1}+b_{d+j} y_{d+j} x_{2 n+1}\right) \\
& +\sum_{j=d+1}^{n} \frac{1}{2}\left(b_{2 j-1}-i b_{2 j}\right)\left(-z_{j} x_{2 n+1}+w_{j} x_{2 n+1}\right) \\
& +\sum_{j=d+1}^{n} \frac{1}{2}\left(b_{2 j-1}+i b_{2 j}\right)\left(-\bar{z}_{j} x_{2 n+1}+\bar{w}_{j} x_{2 n+1}\right) \\
& +O\left(\left|\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right|^{3}\right) . \tag{3.98}
\end{align*}
$$

## 4. The Distribution Kernels of the Maps $\sigma$ and $\sigma^{*} \sigma$

In this section, we will study the map $\sigma$ defined in (1.20) and prove Theorem 1.5 We assume throughout that $\operatorname{dim} X_{G} \geq 3$ and $\bar{\partial}_{b, X_{G}}$ has closed range in $L^{2}$. The case when $\operatorname{dim} X_{G}=1$ will be treated separately.

Let $\iota^{*}: \mathscr{C}^{\infty}(X) \rightarrow \mathscr{C}^{\infty}(Y)$ be the pull-back of the inclusion $\iota: Y \rightarrow X$. Let $\iota_{G}: \mathscr{C}^{\infty}(Y)^{G} \rightarrow \mathscr{C}^{\infty}\left(X_{G}\right)$ be the natural identification. Let $S_{X_{G}}: L^{2}\left(X_{G}\right) \rightarrow$ $H_{b}^{0}\left(X_{G}\right)$ be the orthogonal projection with respect to $(\cdot, \cdot)_{X_{G}}$ (cf. Convention 2.8). By Theorem 3.27 and (1.17) (here we use that $\bar{\partial}_{b, X_{G}}$ has closed range in $L^{2}$ ), we can extend $\sigma$ to $\mathscr{C}^{\infty}(X)$ by

$$
\begin{align*}
\sigma: \mathscr{C}^{\infty}(X) & \rightarrow H_{b}^{0}\left(X_{G}\right) \cap \mathscr{C}^{\infty}\left(X_{G}\right) \subset \mathscr{C}^{\infty}\left(X_{G}\right),  \tag{4.1}\\
\sigma & =S_{X_{G}} \circ E \circ \iota_{G} \circ f_{G} \circ \iota^{*} \circ S_{G} .
\end{align*}
$$

Let $\sigma^{*}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow \mathscr{D}^{\prime}(X)$ be the formal adjoint of $\sigma$. We will show in Theorem 4.12 that $\sigma^{*}$ actually maps $\mathscr{C}^{\infty}\left(X_{G}\right)$ into $H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)^{G}$.

In this section, we will study the distribution kernels of $\sigma$ and $\sigma^{*} \sigma$. We explain briefly the role of the operator $E$ in (4.1). To prove our main result, we need to show that $\sigma^{*} \sigma$ is "microlocally close" to $S_{G}$. In other words, we want $\sigma^{*} \sigma$ to be a complex Fourier integral operator with the same phase, the same order and the same leading symbol as $S_{G}$. To achieve this we need to take $E$ to be a classical elliptic pseudodifferential operator with principal symbol $p_{E}(x, \xi)=|\xi|^{-d / 4}$, see also Remark 1.6

This section is organized as follows. In Sec. 4.1 we develop the calculus of Fourier integral operators of $G$-Szegő type. In Sec. 4.2, we study the distribution kernels of $\sigma$ and $\sigma^{*} \sigma$ and prove Theorem 1.5,

### 4.1. Calculus of complex Fourier integral operators

Let $p \in Y$ and $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ be the local coordinates as in the discussion before Theorem 3.28 defined in an open neighborhood $U$ of $p$. From now on, we change $x_{2 n+1}$ to $x_{2 n+1}-\sum_{j=1}^{d} a_{j}(x) x_{d+j}$, where $a_{j}(x)$ are as in (3.88). With this new local coordinates $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ on $Y \cap U$ we have

$$
\begin{equation*}
J\left(\frac{\partial}{\partial x_{j}}\right)=\frac{\partial}{\partial x_{d+j}}, \quad \text { for } j=1,2, \ldots, d \tag{4.2}
\end{equation*}
$$

Moreover, it is clear that $\Phi(x, y)+\sum_{j=1}^{d} a_{j}(x) x_{d+j}-\sum_{j=1}^{d} a_{j}(y) y_{d+j}$ satisfies (3.98). Note that $a_{j}(x)$ is a smooth function on $Y \cap U$, independent of $x_{1}, \ldots, x_{2 d}, x_{2 n+1}$ and $a_{j}(0)=0, j=1, \ldots, d$. We may assume that $U=\Omega_{1} \times \Omega_{2} \times \Omega_{3}$, where $\Omega_{1} \subset \mathbb{R}^{d}, \Omega_{2} \subset \mathbb{R}^{d}$ are open neighborhoods of $0 \in \mathbb{R}^{d}, \Omega_{3} \subset \mathbb{R}^{2 n+1-2 d}$ is an open neighborhood of $0 \in \mathbb{R}^{2 n+1-2 d}$. From now on, we identify $\Omega_{2}$ with

$$
\left\{\left(0, \ldots, 0, x_{d+1}, \ldots, x_{2 d}, 0, \ldots, 0\right) \in U:\left(x_{d+1}, \ldots, x_{2 d}\right) \in \Omega_{2}\right\}
$$

$\Omega_{3}$ with $\left\{\left(0, \ldots, 0, x_{2 d+1}, \ldots, x_{2 n+1}\right) \in U:\left(x_{2 d+1}, \ldots, x_{2 n+1}\right) \in \Omega_{3}\right\}, \Omega_{2} \times \Omega_{3}$ with

$$
\left\{\left(0, \ldots, 0, x_{d+1}, \ldots, x_{2 n+1}\right) \in U:\left(x_{d+1}, \ldots, x_{2 n+1}\right) \in \Omega_{2} \times \Omega_{3}\right\}
$$

For $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$, we set

$$
\begin{align*}
x^{\prime} & =\left(x_{1}, \ldots, x_{d}\right), \\
x^{\prime \prime} & =\left(x_{d+1}, \ldots, x_{2 n+1}\right), \quad \dot{x}^{\prime \prime}=\left(x_{d+1}, \ldots, x_{2 n}\right),  \tag{4.3}\\
\widehat{x}^{\prime \prime} & =\left(x_{d+1}, \ldots, x_{2 d}\right), \quad \underline{x}^{\prime \prime}=\left(x_{2 d+1}, \ldots, x_{2 n+1}\right) .
\end{align*}
$$

From now on, we identify

$$
\begin{align*}
& x^{\prime \prime} \text { with }\left(0, \ldots, 0, x_{d+1}, \ldots, x_{2 n+1}\right) \in U \\
& \widehat{x}^{\prime \prime} \text { with }\left(0, \ldots, 0, x_{d+1}, \ldots, x_{2 d}, 0, \ldots, 0\right) \in U \text {, }  \tag{4.4}\\
& \underline{x}^{\prime \prime} \text { with }\left(0, \ldots, 0, x_{2 d+1}, \ldots, x_{2 n+1}\right) \in U
\end{align*}
$$

Since $G$ acts freely on $Y$, we take $\Omega_{2}$ and $\Omega_{3}$ small enough so that if $x, x_{1} \in \Omega_{2} \times \Omega_{3}$ and $x \neq x_{1}$, then

$$
\begin{equation*}
g \cdot x \neq x_{1}, \quad \text { for all } g \in G \tag{4.5}
\end{equation*}
$$

Recall that we take $\Phi$ so that (3.96), (3.97), (3.98) hold. Put

$$
\begin{equation*}
\Phi^{*}(x, y):=-\bar{\Phi}(y, x) \tag{4.6}
\end{equation*}
$$

From (3.97) and notice that for $j=1, \ldots, d, x \in Y$, we have $\frac{\partial}{\partial x_{j}}+i \frac{\partial}{\partial x_{d+j}} \in T_{x}^{0,1} X$ and $\frac{\partial}{\partial x_{j}} \Phi(x, y)=\frac{\partial}{\partial y_{j}} \Phi^{*}(x, y)=0$, we conclude that for $j=1, \ldots, d$,

$$
\begin{equation*}
\left.\frac{\partial}{\partial x_{d+j}} \Phi(x, y)\right|_{x_{d+1}=\ldots=x_{2 d}=0} \text { and }\left.\frac{\partial}{\partial y_{d+j}} \Phi^{*}(x, y)\right|_{y_{d+1}=\ldots=y_{2 d}=0} \tag{4.7}
\end{equation*}
$$

vanish to infinite order at $\operatorname{diag}((Y \cap U) \times(Y \cap U))$.
Let

$$
\begin{aligned}
G_{j}(x, y) & :=\left.\frac{\partial}{\partial y_{d+j}} \Phi^{*}(x, y)\right|_{y_{d+1}=\ldots=y_{2 d}=0} \\
H_{j}(x, y) & :=\left.\frac{\partial}{\partial x_{d+j}} \Phi(x, y)\right|_{x_{d+1}=\ldots=x_{2 d}=0} .
\end{aligned}
$$

Put

$$
\begin{align*}
& \Phi_{1}(x, y):=\Phi^{*}(x, y)-\sum_{j=1}^{d} y_{d+j} G_{j}(x, y)  \tag{4.8}\\
& \Phi_{2}(x, y):=\Phi(x, y)-\sum_{j=1}^{d} x_{d+j} H_{j}(x, y) .
\end{align*}
$$

Then for $j=1,2, \ldots, d$,

$$
\begin{equation*}
\left.\frac{\partial}{\partial y_{d+j}} \Phi_{1}(x, y)\right|_{y_{d+1}=\ldots=y_{2 d}=0}=\left.\frac{\partial}{\partial x_{d+j}} \Phi_{2}(x, y)\right|_{x_{d+1}=\ldots=x_{2 d}=0}=0 \tag{4.9}
\end{equation*}
$$

and

$$
\begin{align*}
& \Phi^{*}(x, y)-\Phi_{1}(x, y) \text { vanishes to infinite order on } \operatorname{diag}((Y \cap U) \times(Y \cap U)) \\
& \Phi(x, y)-\Phi_{2}(x, y) \text { vanishes to infinite order on } \operatorname{diag}((Y \cap U) \times(Y \cap U)) \tag{4.10}
\end{align*}
$$

We also write $u=\left(u_{1}, \ldots, u_{2 n+1}\right)$ to denote the local coordinates of $U$. For any smooth function $h \in \mathscr{C}^{\infty}(U)$, let $\widetilde{h} \in \mathscr{C}^{\infty}\left(U^{\mathbb{C}}\right)$ be an almost analytic extension of
$h$ (see [45], Sec. 1]). Let $\vartheta$ be a local coordinate of $\mathbb{R}$. Let

$$
\begin{equation*}
F\left(\widetilde{x}, \widetilde{y}, \widetilde{u^{\prime \prime}}, \widetilde{\vartheta}\right):=\widetilde{\Phi}_{1}\left(\widetilde{x}, \widetilde{u^{\prime \prime}}\right)+\widetilde{\vartheta} \widetilde{\Phi}_{2}\left(\widetilde{u^{\prime \prime}}, \widetilde{y}\right) \tag{4.11}
\end{equation*}
$$

We consider the following two systems for $j=1,2, \ldots, 2 n-2 d+1$ and $j=1$, $2, \ldots, 2 n-d+1$, respectively,

$$
\begin{align*}
\frac{\partial F}{\partial \widetilde{\vartheta}}\left(\widetilde{x}, \widetilde{y}, \widetilde{u^{\prime \prime}}, \widetilde{\vartheta}\right) & =\widetilde{\Phi}_{2}\left(\widetilde{\underline{u}^{\prime \prime}}, \widetilde{y}\right)=0, \\
\frac{\partial F}{\partial \widetilde{u}_{2 d+j}}\left(\widetilde{x}, \widetilde{y}, \widetilde{u^{\prime \prime}}, \widetilde{\vartheta}\right) & =\frac{\partial \widetilde{\Phi}_{1}}{\partial \widetilde{y}_{2 d+j}}\left(\widetilde{x}, \widetilde{u^{\prime \prime}}\right)+\widetilde{\vartheta} \frac{\partial \widetilde{\Phi}_{2}}{\partial \widetilde{x}_{2 d+j}}\left(\widetilde{u^{\prime \prime}}, \widetilde{y}\right)=0, \tag{4.12}
\end{align*}
$$

and

$$
\begin{align*}
\frac{\partial F}{\partial \widetilde{\vartheta}}\left(\widetilde{x}, \widetilde{y}, \widetilde{u^{\prime \prime}}, \widetilde{\vartheta}\right) & =\widetilde{\Phi}_{2}\left(\widetilde{u^{\prime \prime}}, \widetilde{y}\right)=0 \\
\frac{\partial F}{\partial \widetilde{u}_{d+j}}\left(\widetilde{x}, \widetilde{y}, \widetilde{u^{\prime \prime}}, \widetilde{\vartheta}\right) & =\frac{\partial \widetilde{\Phi}_{1}}{\partial \widetilde{y}_{d+j}}\left(\widetilde{x}, \widetilde{u^{\prime \prime}}\right)+\widetilde{\vartheta} \frac{\partial \widetilde{\Phi}_{2}}{\partial \widetilde{x}_{d+j}}\left(\widetilde{u^{\prime \prime}}, \widetilde{y}\right)=0 \tag{4.13}
\end{align*}
$$

where $\widetilde{u^{\prime \prime}}=\left(0, \ldots, 0, \widetilde{u}_{2 d+1}, \ldots, \widetilde{u}_{2 n+1}\right), \widetilde{u^{\prime \prime}}=\left(0, \ldots, 0, \widetilde{u}_{d+1}, \ldots, \widetilde{u}_{2 n+1}\right)$. Here we always use $\frac{\partial}{\partial \widetilde{x}_{k}}$ for first variable, $\frac{\partial}{\partial \tilde{y}_{k}}$ for second variable. From (3.96) and (4.9), we can take $\widetilde{\Phi}_{1}$ and $\widetilde{\Phi}_{2}$ so that for every $j=1,2, \ldots, d$,

$$
\begin{equation*}
\frac{\partial \widetilde{\Phi}_{1}}{\partial \widetilde{y}_{d+j}}\left(\widetilde{x}, \widetilde{u^{\prime \prime}}\right)=\frac{\partial \widetilde{\Phi}_{2}}{\partial \widetilde{x}_{d+j}}\left(\widetilde{u^{\prime \prime}}, \widetilde{y}\right)=0, \quad \text { if } \widetilde{u}_{d+1}=\ldots=\widetilde{u}_{2 d}=0 \tag{4.14}
\end{equation*}
$$

and $\widetilde{\widehat{\Phi}}_{1}, \widetilde{\widehat{\Phi}}_{2} \in \mathscr{C}^{\infty}\left(U^{\mathbb{C}} \times U^{\mathbb{C}}\right)$ such that

$$
\begin{equation*}
\widetilde{\Phi}_{1}(\widetilde{x}, \widetilde{y})=-\widetilde{x}_{2 n+1}+\widetilde{\widehat{\Phi}}_{1}\left(\widetilde{\stackrel{\rightharpoonup}{x}}^{\prime \prime}, \widetilde{y^{\prime \prime}}\right), \quad \widetilde{\Phi}_{2}(\widetilde{x}, \widetilde{y})=\widetilde{y}_{2 n+1}+\widetilde{\hat{\Phi}}_{2}\left(\widetilde{x^{\prime \prime}}, \widetilde{y}^{\prime \prime}\right) \tag{4.15}
\end{equation*}
$$

where $\widetilde{\mathscr{x}}^{\prime \prime}=\left(0, \ldots, 0, \widetilde{x}_{d+1}, \ldots, \widetilde{x}_{2 n}, 0\right), \widetilde{y}^{\prime \prime}=\left(0, \ldots, 0, \widetilde{y}_{d+1}, \ldots, \widetilde{y}_{2 n}, 0\right)$.
By (3.86), (4.6), (4.7) and $\underline{x}^{\prime \prime} \in Y$, we know that $d_{x} \Phi(x, x), d_{y} \Phi(x, x)$ are real for $x \in Y$, and for $j=1,2, \ldots, 2 n-d$,

$$
\begin{aligned}
\frac{\partial \widetilde{\Phi}_{1}}{\partial \widetilde{y}_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) & =-\frac{\partial \bar{\Phi}}{\partial \widetilde{x}_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=-\frac{\partial \Phi}{\partial x_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) \\
& =\frac{\partial \Phi}{\partial y_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=\frac{\partial \Phi_{2}}{\partial y_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right),
\end{aligned}
$$

$$
\frac{\partial \widetilde{\Phi}_{2}}{\partial \widetilde{y}_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)+\frac{\partial \widetilde{\Phi}_{2}}{\partial \widetilde{x}_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=0
$$

Note that by (3.87), $\Phi(x, x)=0$ for $x \in Y$. Combining it with the above equation, we get

$$
\begin{align*}
\widetilde{\Phi}_{2}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) & =0 \\
\frac{\partial \widetilde{\Phi}_{1}}{\partial \widetilde{y}_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)+\left.\widetilde{\vartheta} \frac{\partial \widetilde{\Phi}_{2}}{\partial \widetilde{x}_{d+j}}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)\right|_{\tilde{\vartheta}=1} & =0, \quad j=1,2, \ldots, 2 n-d, \tag{4.16}
\end{align*}
$$

and the Hessians

$$
F_{\widetilde{\vartheta}, \widetilde{u^{\prime \prime}}}(0,0,0,1)=\left.\left(\begin{array}{cc}
\frac{\partial^{2} F}{\partial^{2} \widetilde{\vartheta}} & \frac{\partial^{2} F}{\partial \widetilde{\vartheta} \partial \widetilde{u^{\prime \prime}}}  \tag{4.17}\\
\frac{\partial^{2} F}{\partial \widetilde{u^{\prime \prime}} \partial \widetilde{\vartheta}} & \frac{\partial^{2} F}{\partial^{2} \widetilde{u^{\prime \prime}}}
\end{array}\right)\right|_{(0,0,0,1)}
$$

and

$$
F_{\widetilde{\vartheta}, \widetilde{u^{\prime \prime}}}(0,0,0,1)=\left.\left(\begin{array}{cc}
\frac{\partial^{2} F}{\partial^{2} \widetilde{\vartheta}} & \frac{\partial^{2} F}{\partial \widetilde{\vartheta} \partial \widetilde{u^{\prime \prime}}}  \tag{4.18}\\
\frac{\partial^{2} F}{\partial \widetilde{u^{\prime \prime}} \partial \widetilde{\vartheta}} & \frac{\partial^{2} F}{\partial^{2} \widetilde{u^{\prime \prime}}}
\end{array}\right)\right|_{(0,0,0,1)}
$$

are non-singular. Moreover, from (3.98), we calculate

$$
\begin{align*}
& \operatorname{det}\left(\frac{1}{2 \pi i} F_{\widetilde{\vartheta}, \widetilde{\underline{u}^{\prime \prime}}}\right)(0,0,0,1)=\frac{2^{2 n-2 d-2}}{\pi^{2 n-2 d+2}}\left(\left|\mu_{d+1}\right| \ldots\left|\mu_{n}\right|\right)^{2} \\
& \operatorname{det}\left(\frac{1}{2 \pi i} F_{\widetilde{\vartheta}, \widetilde{u^{\prime \prime}}}\right)(0,0,0,1)=\frac{2^{2 n-2}}{\pi^{2 n-d+2}}\left(\left|\mu_{1}\right| \ldots\left|\mu_{d}\right|\right)\left(\left|\mu_{d+1}\right| \ldots\left|\mu_{n}\right|\right)^{2} . \tag{4.19}
\end{align*}
$$

Hence, near $(p, p)$ and $\widetilde{\vartheta}=1$, we can solve (4.12) and (4.13) and the solutions are unique. Let

$$
\begin{align*}
\widetilde{u^{\prime \prime}} & =\alpha(x, y)  \tag{4.20}\\
\widetilde{\vartheta} & \left.=\gamma(x, y) \in \alpha_{2 d+1}(x, y), \ldots, \alpha_{2 n+1}(x, y)\right) \in \mathscr{C}^{\infty}(U \times U, \mathbb{C})
\end{align*}
$$

and

$$
\begin{align*}
\widetilde{u^{\prime \prime}} & =\beta(x, y)=\left(\beta_{d+1}(x, y), \ldots, \beta_{2 n+1}(x, y)\right) \in \mathscr{C}^{\infty}\left(U \times U, \mathbb{C}^{2 n-d+1}\right) \\
\widetilde{\vartheta} & =\delta(x, y) \in \mathscr{C}^{\infty}(U \times U, \mathbb{C}) \tag{4.21}
\end{align*}
$$

be the solutions of (4.12) and (4.13), respectively. From (4.14), it is easy to see that

$$
\begin{align*}
\beta(x, y) & =\left(\beta_{d+1}(x, y), \ldots, \beta_{2 n+1}(x, y)\right) \\
& =\left(0, \ldots, 0, \alpha_{2 d+1}(x, y), \ldots, \alpha_{2 n+1}(x, y)\right)  \tag{4.22}\\
\gamma(x, y) & =\delta(x, y)
\end{align*}
$$

From (4.22), we see that the value of $\widetilde{\Phi}_{1}\left(x, \widetilde{u^{\prime \prime}}\right)+\widetilde{\vartheta}_{2}\left(\widetilde{u^{\prime \prime}}, y\right)$ at critical points $\widetilde{u^{\prime \prime}}=$ $\alpha(x, y), \widetilde{\vartheta}=\gamma(x, y)$ is equal to the value of $\widetilde{\Phi}_{1}\left(x, \widetilde{u^{\prime \prime}}\right)+\widetilde{\vartheta} \widetilde{\Phi}_{2}\left(\widetilde{u^{\prime \prime}}, y\right)$ at critical points

$$
\begin{align*}
& \widetilde{u^{\prime \prime}}=\beta(x, y), \widetilde{\vartheta}=\delta(x, y) . \text { Put } \\
& \qquad \begin{aligned}
\Phi_{3}(x, y) & :=\widetilde{\Phi}_{1}(x, \alpha(x, y))+\gamma(x, y) \widetilde{\Phi}_{2}(\alpha(x, y), y) \\
& =\widetilde{\Phi}_{1}(x, \beta(x, y))+\delta(x, y) \widetilde{\Phi}_{2}(\beta(x, y), y)
\end{aligned}
\end{align*}
$$

Then $\Phi_{3}(x, y)$ is a complex phase function, $\operatorname{Im} \Phi_{3}(x, y) \geq 0$ and $\Phi_{3}(x, y)=$ $\Phi_{3}\left(x^{\prime \prime}, y^{\prime \prime}\right)$. It is easy to check that

$$
\begin{align*}
d_{x} \Phi_{3}(x, x) & =-d_{y} \Phi_{3}(x, x)=d_{x} \Phi(x, x) \\
& =-d_{y} \Phi(x, x)=-\lambda(x) \omega_{0}(x), \quad \lambda(x)>0, \quad \text { for } x \in U \cap Y . \tag{4.24}
\end{align*}
$$

From now on we take $U$ small enough so that the Levi form is positive on $U$ and

$$
\begin{equation*}
d_{x} \Phi_{3}(x, y) \neq 0, \quad d_{y} \Phi_{3}(x, y) \neq 0, \quad \text { for every }(x, y) \in U \times U \tag{4.25}
\end{equation*}
$$

and $a_{0}(x, y) \neq 0$, for every $(x, y) \in U \times U$, where $a_{0}(x, y) \in \mathscr{C}^{\infty}(U \times U)$ is as in (3.85).

Fix an open neighborhood $\mathcal{U} \Subset U$ of $p$ with $\widehat{\Omega}_{2} \times \widehat{\Omega}_{3} \subset \mathcal{U}$, where $\widehat{\Omega}_{2} \Subset \Omega_{2} \subset \mathbb{R}^{d}$ is an open neighborhood of $0 \in \mathbb{R}^{d}$ and $\widehat{\Omega}_{3} \Subset \Omega_{3} \subset \mathbb{R}^{2 n+1-2 d}$ is an open neighborhood of $0 \in \mathbb{R}^{2 n+1-2 d}$.

Theorem 4.1. The phase functions $\Phi$ and $\Phi_{3}$ are equivalent on $\mathcal{U}$, that is, for any $b_{1} \in S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$there exist $\hat{b}_{1}, b_{2} \in S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$such that

$$
\begin{align*}
& \int_{0}^{\infty} e^{i \Phi(x, y) t} b_{1}(x, y, t) d t \equiv \int_{0}^{\infty} e^{i \Phi_{3}(x, y) t} b_{2}(x, y, t) d t \quad \text { on } \mathcal{U} \times \mathcal{U}  \tag{4.26}\\
& \int_{0}^{\infty} e^{i \Phi(x, y) t} \hat{b}_{1}(x, y, t) d t \equiv \int_{0}^{\infty} e^{i \Phi_{3}(x, y) t} b_{1}(x, y, t) d t \quad \text { on } \mathcal{U} \times \mathcal{U}
\end{align*}
$$

Proof. We consider the kernel $\left(S_{G} \circ S_{G}\right)(\cdot, \cdot)$ on $U$. Let $\mathcal{U} \Subset U_{1} \Subset U$ be open neighborhoods of $p$. Let $\chi\left(x^{\prime \prime}\right) \in \mathscr{C}_{0}^{\infty}\left(\Omega_{2} \times \Omega_{3}\right)$. By (4.5) we can extend $\chi\left(x^{\prime \prime}\right)$ to

$$
W:=\left\{g \cdot x: g \in G, \quad x \in \Omega_{2} \times \Omega_{3}\right\}
$$

by $\chi\left(g \cdot x^{\prime \prime}\right):=\chi\left(x^{\prime \prime}\right)$, for every $g \in G$. Assume that $\chi=1$ on some neighborhood of $\bar{U}_{1}$. Let $\chi_{1} \in \mathscr{C}_{0}^{\infty}(U)$ with $\chi_{1}=1$ on some neighborhood of $\bar{U}_{1}$ and $\operatorname{supp} \chi_{1} \subset$ $\{x \in X: \chi(x)=1\}$. We have

$$
\begin{equation*}
\chi_{1} S_{G} \circ S_{G}=\chi_{1} S_{G} \chi \circ S_{G}+\chi_{1} S_{G}(1-\chi) \circ S_{G} \tag{4.27}
\end{equation*}
$$

Let us first consider $\chi_{1} S_{G}(1-\chi) \circ S_{G}$. We have

$$
\begin{align*}
\left(\chi_{1} S_{G}(1-\chi)\right)(x, u) & =\chi_{1}(x) \int_{G} S_{\leq \lambda_{0}}(x, g \cdot u)(1-\chi(u)) d \mu(g) \\
& =\chi_{1}(x) \int_{G} S_{\leq \lambda_{0}}(x, u)\left(1-\chi\left(g^{-1} \cdot u\right)\right) d \mu(g) \tag{4.28}
\end{align*}
$$

where $\lambda_{0}>0$ is a small constant as in Theorem3.19 If $g^{-1} \cdot u \notin\{x \in X: \chi(x)=1\}$, since supp $\chi_{1} \subset\{x \in X: \chi(x)=1\}$ and $\chi(x)=\chi(g \cdot x)$, for every $g \in G, x \in X$, we
conclude that $u \notin \operatorname{supp} \chi_{1}$. From this observation and since that $S_{\leq \lambda_{0}}$ is smoothing away the diagonal on $G U$ (see Theorem (3.22), we deduce that $\chi_{1} S_{G}(1-\chi)$ is smoothing and hence

$$
\begin{equation*}
\chi_{1} S_{G}(1-\chi) \circ S_{G} \equiv 0 \quad \text { on } X \times X \tag{4.29}
\end{equation*}
$$

From (4.27) and (4.29), we get

$$
\begin{equation*}
\chi_{1} S_{G} \circ S_{G} \equiv \chi_{1} S_{G} \chi \circ S_{G} \quad \text { on } X \times X \tag{4.30}
\end{equation*}
$$

From Theorem 3.25 and using that $S_{G}^{*}=S_{G}$, where $S_{G}^{*}$ is the adjoint of $S_{G}$ with respect to $(\cdot, \cdot)$, we obtain that on $U$,

$$
\begin{align*}
& \left(\chi_{1} S_{G} \chi \circ S_{G}\right)(x, y) \\
& \equiv \\
& \equiv \int_{\Omega_{2} \times \Omega_{3}} \int_{0}^{\infty} \int_{0}^{\infty} e^{i \Phi^{*}\left(x, u^{\prime \prime}\right) t+i \Phi\left(u^{\prime \prime}, y\right) s} \chi_{1}(x) a^{*}\left(x^{\prime \prime}, u^{\prime \prime}, t\right) V_{\mathrm{eff}}\left(u^{\prime \prime}\right) \\
& \quad \times \chi\left(u^{\prime \prime}\right) a\left(u^{\prime \prime}, y^{\prime \prime}, s\right) d s d t d v\left(u^{\prime \prime}\right) \\
& \equiv \\
& \equiv \int_{\Omega_{2} \times \Omega_{3}} \int_{0}^{\infty} \int_{0}^{\infty} e^{i \Phi_{1}\left(x, u^{\prime \prime}\right) t+i \Phi_{2}\left(u^{\prime \prime}, y\right) s} \chi_{1}(x) b\left(x, u^{\prime \prime}, t\right) \\
& \quad \times \chi\left(u^{\prime \prime}\right) c\left(u^{\prime \prime}, y, s\right) d s d t d v\left(u^{\prime \prime}\right) \quad(\text { here we used (4.10) })  \tag{4.31}\\
& \equiv \\
& \quad \int_{\Omega_{2} \times \Omega_{3}} \int_{0}^{\infty} \int_{0}^{\infty} e^{i\left(\Phi_{1}\left(x, u^{\prime \prime}\right)+\Phi_{2}\left(u^{\prime \prime}, y\right) \vartheta\right) t} \\
& \quad \times \chi_{1}(x) b\left(x, u^{\prime \prime}, t\right) \chi\left(u^{\prime \prime}\right) c\left(u^{\prime \prime}, y, t \vartheta\right) t d \vartheta d t d v\left(u^{\prime \prime}\right)
\end{align*}
$$

where $s=\vartheta t, d \mu(g) d v\left(u^{\prime \prime}\right)=d v(x)$ on $U, a^{*}\left(x^{\prime \prime}, u^{\prime \prime}, t\right)=\bar{a}\left(u^{\prime \prime}, x^{\prime \prime}, t\right)$ and

$$
\begin{align*}
& b(x, y, t), c(x, y, t) \in S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(U \times U \times \mathbb{R}_{+}\right)  \tag{4.32}\\
& b_{0}(x, x) \neq 0, c_{0}(x, x) \neq 0, \text { for any } x \in U \cap Y(\text { cf. Notation (3.76) }) .
\end{align*}
$$

We apply the complex stationary phase formula of Melin-Sjöstrand 45 to carry out the $d v\left(u^{\prime \prime}\right) d \vartheta$ integration in 4.31). This yields the existence of a symbol $d(x, y, t) \in S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(U \times U \times \mathbb{R}_{+}\right)$with the expansion $d(x, y, t) \sim \sum_{j=0}^{\infty} t^{n-\frac{d}{2}-j} d_{j}(x, y)$ in $S_{1,0}^{n-\frac{d}{2}}\left(U \times U \times \mathbb{R}_{+}\right)$(cf. (3.76)) with $d_{0}(x, x) \neq 0$ for $x \in U \cap Y$ and

$$
\begin{equation*}
\left(\chi_{1} S_{G} \chi \circ S_{G}\right)(x, y) \equiv \int_{0}^{\infty} e^{i \Phi_{3}(x, y) t} d(x, y, t) d t \quad \text { on } U \times U \tag{4.33}
\end{equation*}
$$

From (4.30), (4.33), we deduce that

$$
\begin{equation*}
\int_{0}^{\infty} e^{i \Phi_{3}(x, y) t} d(x, y, t) d t \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} \chi_{1}(x) a(x, y, t) d t \quad \text { on } U \times U \tag{4.34}
\end{equation*}
$$

Note that $\Phi_{3}$ and $\Phi$ are independent of $x^{\prime}$ and $y^{\prime}$. By the Malgrange preparation theorem we may assume that $\Phi_{3}$ and $\Phi$ have the form

$$
\begin{aligned}
\Phi_{3}\left(x^{\prime \prime}, y^{\prime \prime}\right) & =y_{2 n+1}+h\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right) \\
\Phi\left(x^{\prime \prime}, y^{\prime \prime}\right) & =y_{2 n+1}+h_{1}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)
\end{aligned}
$$

where $\hat{\Phi}_{3}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right), \hat{\Phi}\left(x^{\prime \prime}, \hat{y}^{\prime \prime}\right) \in \mathscr{C}^{\infty}(U \times U)$. From (4.34), we can repeat the proof of [28, Theorem 5.2] and deduce that $\Phi_{3}(x, y)-\Phi(x, y)$ vanishes to infinite order at $\operatorname{diag}(Y \times Y)$. Since $\Phi_{3}(x, y)$ and $\Phi(x, y)$ are independent of $x^{\prime}, y^{\prime}$, we conclude that $\Phi_{3}(x, y)-\Phi(x, y)$ vanishes to infinite order on the underlying canonical relation

$$
\{(x, g \cdot x): x \in Y, g \in G\} \cap(U \times U)
$$

Hence, $\Phi$ and $\Phi_{3}$ are equivalent on $\mathcal{U}$. For the convenience of the reader, we sketch briefly the idea of the proof. We will use some semi-classical notations as in the proof of [28, Theorem 5.2]. Suppose that $U=U^{\prime} \times(-\varepsilon, \varepsilon), \varepsilon>0, U^{\prime}$ is an open set of $\mathbb{R}^{2 n}$. Let $\tau \in \mathscr{C}_{c}^{\infty}((-\varepsilon, \varepsilon)), \tau \geq 0, \tau \equiv 1$ on $\left[-\frac{\varepsilon}{2}, \frac{\varepsilon}{2}\right]$. For each $k>0$, we consider the distributions (see the proof of [28, Theorem 5.2])

$$
\begin{aligned}
& A_{k}: u \mapsto \iint_{0}^{\infty} e^{i\left(y_{2 n+1}+h\left(x^{\prime \prime}, \mathfrak{y}^{\prime \prime}\right)\right) t-i k y_{2 n+1}} d(x, y, t) \tau\left(y_{2 n+1}\right) u(\grave{y}) d y d t \\
& B_{k}: u \mapsto \iint_{0}^{\infty} e^{i\left(y_{2 n+1}+h_{1}\left(x^{\prime \prime}, \hat{y}^{\prime \prime}\right)\right) t-i k y_{2 n+1}} \tau(x) \chi_{1}(x) a(x, y, t) \tau_{1} \tau\left(y_{2 n+1}\right) u(\grave{y}) d y d t
\end{aligned}
$$

for $u \in \mathscr{C}_{0}^{\infty}\left(U^{\prime}\right)$, where $\dot{y}=\left(y_{1}, \ldots, y_{2 n}\right)$. By using the stationary phase formula of Melin-Sjöstrand [45], we can show that (cf. the proof of [29, Theorem 3.12]) $A_{k}$ and $B_{k}$ are smoothing operators and

$$
\begin{aligned}
A_{k}(x, \dot{y}) & \equiv e^{i k h\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)} g(x, \dot{y}, k)+O\left(k^{-\infty}\right), \\
B_{k}(x, \dot{y}) & \equiv e^{i k h_{1}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)} p(x, \check{y}, k)+O\left(k^{-\infty}\right), \\
g(x, \check{y}, k), p(x, \check{y}, k) & \in S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(U \times U^{\prime}\right), \\
g(x, \dot{y}, k) & \sim \sum_{j=0}^{\infty} g_{j}(x, \stackrel{y}{y}) k^{n-\frac{d}{2}-j} \quad \text { in } S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(U \times U^{\prime}\right), \\
p(x, \dot{y}, k) & \sim \sum_{j=0}^{\infty} p_{j}\left(x, y^{\prime}\right) k^{n-\frac{d}{2}-j} \quad \text { in } S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(U \times U^{\prime}\right), \\
g_{j}(x, \grave{y}), p_{j}(x, \dot{y}) & \in \mathscr{C}^{\infty}\left(U \times U^{\prime}\right), \quad j=0,1, \ldots \\
g_{0}\left(x_{0}, \mathscr{x}_{0}\right) & \neq 0
\end{aligned}
$$

Since

$$
\int_{0}^{\infty} e^{i\left(y_{2 n+1}+h\left(x^{\prime \prime}, \tilde{y}^{\prime \prime}\right)\right) t} d(x, y, t) d t-\int_{0}^{\infty} e^{i\left(y_{2 n+1}+h_{1}\left(x^{\prime \prime}, \dot{y}^{\prime \prime \prime}\right)\right) t} \chi_{1}(x) a(x, y, t) d t
$$

is smoothing, by using integration by parts with respect to $y_{2 n+1}$, it is easy to see that $A_{k}-B_{k}=O\left(k^{-\infty}\right)$ (see [29, Sec. 3]). From this observation, we can repeat the argument as in the discussion after [28, (5.2)] and deduce that $h-h_{1}$ vanishes to infinite order at diag $(Y \times Y)$. Since $\Phi_{3}(x, y)$ and $\Phi(x, y)$ are independent of $x^{\prime}$, $y^{\prime}$, we conclude that $\Phi_{3}(x, y)-\Phi(x, y)$ vanishes to infinite order on the underlying canonical relation $\{(x, g \cdot x): x \in Y, g \in G\} \cap(U \times U)$.

The following two theorems follow from the proof of Theorem 4.1 (4.10), (4.23), the proof of (4.34), the complex stationary phase formula of Melin-Sjöstrand 45.

Theorem 4.2. Consider the Fourier integral operators

$$
A(x, y)=\int_{0}^{\infty} e^{i \Phi(x, y) t} a(x, y, t) d t, \quad B(x, y)=\int_{0}^{\infty} e^{i \Phi(x, y) t} b(x, y, t) d t
$$

with symbols $a(x, y, t) \in S_{\mathrm{cl}}^{k}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$and $b(x, y, t) \in S_{\mathrm{cl}}^{\ell}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$. Consider $\chi\left(x^{\prime \prime}\right) \in \mathscr{C}_{0}^{\infty}\left(\widehat{\Omega}_{2} \times \widehat{\Omega}_{3}\right)$. Then, we have

$$
\int A\left(x, u^{\prime \prime}\right) \chi\left(u^{\prime \prime}\right) B\left(u^{\prime \prime}, y\right) d v\left(u^{\prime \prime}\right) \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} c(x, y, t) d t \quad \text { on } \mathcal{U} \times \mathcal{U}
$$

with $c(x, y, t) \in S_{\mathrm{cl}}^{k+\ell-\left(n-\frac{d}{2}\right)}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$. For $x \in Y \cap U$ we have

$$
\begin{equation*}
c_{0}(x, x)=2^{-n-\frac{d}{2}+1} \pi^{n-\frac{d}{2}+1}\left|\operatorname{det} \mathscr{L}_{x}\right|^{-1}\left|\operatorname{det} \mathcal{R}_{x}\right|^{\frac{1}{2}} a_{0}(x, x) b_{0}(x, x) \chi\left(x^{\prime \prime}\right) \tag{4.35}
\end{equation*}
$$

where $\operatorname{det} \mathcal{R}_{x}$ is the determinant $\mathcal{R}_{x}$ cf. (3.83). Moreover, if there are $N_{1}, N_{2} \in \mathbb{N}^{*}$, $C>0$, such that for all $x_{0} \in Y \cap \mathcal{U}$,

$$
\begin{equation*}
\left|a_{0}(x, y)\right| \leq C\left|(x, y)-\left(x_{0}, x_{0}\right)\right|^{N_{1}}, \quad\left|b_{0}(x, y)\right| \leq C\left|(x, y)-\left(x_{0}, x_{0}\right)\right|^{N_{2}} \tag{4.36}
\end{equation*}
$$

then there exists $\widehat{C}>0$ such that for all $x_{0} \in Y \cap \mathcal{U}$

$$
\begin{equation*}
\left|c_{0}(x, y)\right| \leq \widehat{C}\left|(x, y)-\left(x_{0}, x_{0}\right)\right|^{N_{1}+N_{2}} \tag{4.37}
\end{equation*}
$$

Proof. By (4.10) and since $\Phi, \Phi_{1}, \Phi_{2}$ and $\Phi^{*}$ are independent of $\left(x^{\prime}, y^{\prime}\right), \Phi-\Phi_{2}$, $\Phi^{*}-\Phi_{1}$ vanish to infinite order on the underlying canonical relation

$$
\{(g \cdot x, x): x \in Y, g \in G\}
$$

Hence, we can change the phase of $A$ to $\Phi_{1}$ and the phase of $B$ to $\Phi_{2}$. Thus,

$$
\begin{aligned}
& \int A\left(x, u^{\prime \prime}\right) \chi\left(u^{\prime \prime}\right) B\left(u^{\prime \prime}, y\right) d v\left(u^{\prime \prime}\right) \\
& \quad \equiv \int d t \int_{0}^{\infty} d s \int_{0}^{\infty} e^{i \Phi_{1}\left(x, u^{\prime \prime}\right) t+i \Phi_{2}\left(u^{\prime \prime}, y\right) s} e\left(x, u^{\prime \prime}, t\right) f\left(u^{\prime \prime}, y, s\right) \chi\left(u^{\prime \prime}\right) d u^{\prime \prime}
\end{aligned}
$$

for some $e \in S_{\mathrm{cl}}^{k}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right), f \in S_{\mathrm{cl}}^{\ell}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$such that the leading terms $e_{0}, f_{0}$ of $e, f$, respectively, satisfy (4.36). From (4.23), the proof of (4.34) and the complex stationary phase formula of Melin-Sjöstrand [45], we get

$$
\begin{aligned}
& \int A\left(x, u^{\prime \prime}\right) \chi\left(u^{\prime \prime}\right) B\left(u^{\prime \prime}, y\right) d v\left(u^{\prime \prime}\right) \\
& \quad \equiv \int_{0}^{\infty} d t \iint_{0}^{\infty} e^{i t\left(\Phi_{1}\left(x, u^{\prime \prime}\right)+\gamma \Phi_{2}\left(u^{\prime \prime}, y\right)\right)} t e\left(x, u^{\prime \prime}, t\right) f\left(u^{\prime \prime}, y, \gamma t\right) \chi\left(u^{\prime \prime}\right) d u^{\prime \prime} d \gamma \\
& \quad \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} c(x, y, t) d t \quad \text { on } \mathcal{U} \times \mathcal{U}
\end{aligned}
$$

with $c(x, y, t) \in S_{\mathrm{cl}}^{k+\ell-\left(n-\frac{d}{2}\right)}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$(as we integrate over $u^{\prime \prime}$ and $\gamma$, whose total dimension is $2 n-d+2$, we need to lower the order by $n+1-d / 2$; however, due to
the factor $t$ in the integral we gain one order, thus we get the order $k+l-(n-d / 2))$ and

$$
\begin{equation*}
c_{0}(x, y)=g(x, y) \widetilde{e}_{0}(x, \beta(x, y)) \widetilde{f}_{0}(\beta(x, y), y) \tag{4.38}
\end{equation*}
$$

for some smooth function $g$, where $\beta(x, y)$ is as in (4.21), $\widetilde{e}_{0}, \widetilde{f}_{0}$ denote almost analytic extensions of $e_{0}, f_{0}$, respectively. Since $\beta(x, x)=x$ for every $x \in Y$, we conclude that $\widetilde{e}_{0}(x, \beta(x, y))$ and $\widetilde{f}_{0}(\beta(x, y), y)$ have vanishing order at least $N_{1}$ and $N_{2}$ on $\operatorname{diag}(Y \times Y)$, respectively. This observation and 4.38) yield (4.37).

Theorem 4.3. Consider the Fourier integral operators

$$
\mathcal{A}\left(x, \underline{y}^{\prime \prime}\right)=\int_{0}^{\infty} e^{i \Phi\left(x, \underline{y}^{\prime \prime}\right) t} \alpha\left(x, \underline{y}^{\prime \prime}, t\right) d t, \quad \mathcal{B}\left(\underline{x}^{\prime \prime}, y\right)=\int_{0}^{\infty} e^{i \Phi\left(\underline{x}^{\prime \prime}, y\right) t} \beta\left(\underline{x}^{\prime \prime}, y, t\right) d t
$$

with symbols $\alpha\left(x, \underline{y}^{\prime \prime}, t\right) \in S_{\mathrm{cl}}^{k}\left(\mathcal{U} \times \Omega_{3} \times \mathbb{R}_{+}\right)$and $\beta\left(\underline{x}^{\prime \prime}, y, t\right) \in S_{\mathrm{cl}}^{\ell}\left(\Omega_{3} \times \mathcal{U} \times \mathbb{R}_{+}\right)$. Let $\chi_{1}\left(\underline{x}^{\prime \prime}\right) \in \mathscr{C}_{0}^{\infty}\left(\bar{\Omega}_{3}\right)$. Then, we have

$$
\int \mathcal{A}\left(x, \underline{u}^{\prime \prime}\right) \chi_{1}\left(\underline{u}^{\prime \prime}\right) \mathcal{B}\left(\underline{u}^{\prime \prime}, y\right) d v\left(\underline{u}^{\prime \prime}\right) \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} \gamma(x, y, t) d t \quad \text { on } \mathcal{U} \times \mathcal{U}
$$

with $\gamma(x, y, t) \in S_{\mathrm{cl}}^{k+\ell-(n-d)}\left(\mathcal{U} \times \mathcal{U} \times \mathbb{R}_{+}\right)$where
$\gamma_{0}(x, x)=2^{-n+1} \pi^{n-d+1}\left|\operatorname{det} \mathscr{L}_{x}\right|^{-1}\left|\operatorname{det} \mathcal{R}_{x}\right| \alpha_{0}\left(x, \underline{x}^{\prime \prime}\right) \beta_{0}\left(\underline{x}^{\prime \prime}, x\right) \chi_{1}\left(\underline{x}^{\prime \prime}\right), \quad x \in Y \cap U$.

Moreover, if there are $N_{1}, N_{2} \in \mathbb{N}^{*}, C>0$, such that for all $x_{0} \in Y \cap \mathcal{U}$ we have

$$
\left|\alpha_{0}\left(x, \underline{y}^{\prime \prime}\right)\right| \leq C\left|\left(x, \underline{y}^{\prime \prime}\right)-\left(x_{0}, x_{0}\right)\right|^{N_{1}}, \quad\left|\beta_{0}\left(x, \underline{y}^{\prime \prime}\right)\right| \leq C\left|\left(x, \underline{y}^{\prime \prime}\right)-\left(x_{0}, x_{0}\right)\right|^{N_{2}}
$$

then there exists $\widehat{C}>0$ such that for all $x_{0} \in Y \cap \mathcal{U}$,

$$
\begin{equation*}
\left|\gamma_{0}(x, y)\right| \leq \widehat{C}\left|(x, y)-\left(x_{0}, x_{0}\right)\right|^{N_{1}+N_{2}} \tag{4.40}
\end{equation*}
$$

The proof of Theorem 4.3 is similar to the proof of Theorem 4.2. We omit the details.

Remark 4.4. (i) The reason why in Theorem 4.3 we integrate only with respect to $\underline{y}^{\prime \prime}$ is that in the proof of our main results, we need to integrate over the reduced space and this corresponds to the integration only with respect to $\underline{y}^{\prime \prime}$ variables.
(ii) Theorems 4.2 and 4.3 are about compositions of Fourier integral operators. In general, the composition of Fourier integral operators correspond to composition of canonical relation but we do not use this point of view since in the proof of our main results we need to know the precise form of the phase function of the composition of our Fourier integral operators.

We introduce next the following notion.

Definition 4.5. Let $H: \mathscr{C}^{\infty}(X) \rightarrow \mathscr{C}^{\infty}(X)^{G}$ be a continuous operator with distribution kernel $H(x, y) \in \mathscr{D}^{\prime}(X \times X)$ and $k \in \mathbb{R}, \ell \in \mathbb{N}$.
(i) We say that $H$ is a complex Fourier integral operator of $G$-Szegő type of leading $\operatorname{order}(k, \ell)$, if for every open set $D$ of $X$ with $D \cap Y=\emptyset$,

$$
\begin{equation*}
\chi H \text { and } H \chi \text { are smoothing operators on } X, \text { for every } \chi \in \mathscr{C}_{0}^{\infty}(D) \tag{4.41}
\end{equation*}
$$

and for every $p \in Y$ and any open coordinate neighborhood ( $U, x=$ $\left.\left(x_{1}, \ldots, x_{2 n+1}\right)\right)$ of $p$, we have

$$
\begin{equation*}
H(x, y) \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} a(x, y, t) d t \quad \text { on } U \times U \tag{4.42}
\end{equation*}
$$

with $\Phi(x, y) \in \mathscr{C}^{\infty}(U \times U)$ as in Theorem 3.25, and

$$
\begin{equation*}
a(x, y, t) \in S_{\mathrm{cl}}^{k}\left(U \times U \times \mathbb{R}_{+}\right) \tag{4.43}
\end{equation*}
$$

and under the notation (3.76),

$$
\begin{equation*}
\left.\frac{\partial^{|\alpha|+|\beta|} a_{0}(x, y)}{\partial x^{\alpha} \partial y^{\beta}}\right|_{x=y \in Y}=0, \quad \text { for } \alpha, \beta \in \mathbb{N}^{2 n+1}, \quad|\alpha|+|\beta| \leq \ell-1 \tag{4.44}
\end{equation*}
$$

(ii) We say that $H$ is a complex Fourier integral operator of $G$-Szegő type of order $(k, \ell), k \in \mathbb{R}, \ell \in \mathbb{N}$, if (4.41), (4.42), 4.43) hold and there is a $r(x, y, t) \in$ $S_{\mathrm{cl}}^{-\infty}\left(U \times U \times \mathbb{R}_{+}\right)$such that

$$
\begin{gather*}
\left.\frac{\partial^{|\alpha|+|\beta|}(a(x, y, t)-r(x, y, t))}{\partial x^{\alpha} \partial y^{\beta}}\right|_{x=y \in Y}=0, \quad \text { for } \alpha, \beta \in \mathbb{N}^{2 n+1} \\
|\alpha|+|\beta| \leq \ell-1 \tag{4.45}
\end{gather*}
$$

Let $G_{k, \ell}(X)$ denote the space of all complex Fourier integral operators of $G$ Szegő type of leading order $(k, \ell)$ and let $\widehat{G}_{k, \ell}(X)$ denote the space of all complex Fourier integral operators of $G$-Szegő type of order $(k, \ell)$.

Note that 4.45) means that each coefficient $a_{j}$ has vanishing order at least $\ell$ at diag $(Y \times Y)$. In Definition 4.5 we use the terminology $G$-Szegő type in order to stress the dependence on the set $Y$ and thus on the group $G$.

Let us explain briefly the role of the spaces $G_{k, \ell}(X)$ and $\widehat{G}_{k, \ell}(X)$. Our goal is to study distribution kernel of $\sigma^{*} \sigma$. In Theorem 4.13, we will show that $C_{0} \sigma^{*} \sigma$ is of the same type as $S_{G}$ and with the same leading term, where $C_{0}$ is a constant. In the terminology introduced in Definition 4.5, $C_{0} \sigma^{*} \sigma-S_{G} \in G_{n-(d / 2), 1}(X)$. To prove our main result, we need to show that $C_{0} \sigma^{*} \sigma-S_{G}$ is "microlocally small", and it suffices to prove that elements $H \in G_{n-(d / 2), 1}(X)$ have good regularity properties (see (4.58)).

Let $H=H_{0}+H_{1}$, where $H_{0} \in G_{n-(d / 2), 1}(X)$ is the leading term of $H$ and $H_{1} \in G_{n-(d / 2)-1,0}(X)$ is the lower order term of $H$. By using calculus of complex Fourier integral operators, we can show that when we compose $H_{1}$ with itself, the order of the composition will decrease. More precisely, $H_{1}^{N} \in G_{n-(d / 2)-N, 0}(X)$, for
every $N \in \mathbb{N}^{*}$. Hence, for large $N, H_{1}^{N}$ has good regularity properties and hence $H_{1}$ itself has good regularity properties.

In order to handle $H_{0}$ we observe that when we compose $H_{0}$ with itself the order of the composition will not decrease, that is, $H_{0}^{N}$ is still in $G_{n-(d / 2), 1}(X)$ for every $N \in \mathbb{N}^{*}$. To get good regularity properties, we need the space $\widehat{G}_{k, \ell}(X)$. Note that the space $\widehat{G}_{k, \ell}(X)$ is the subspace of $G_{k, \ell}(X)$ whose elements have full symbols have vanishing order at least $\ell$ on $\operatorname{diag}(Y \times Y)$. The key observation is that the leading symbol of $H_{0}^{N}$ has vanishing order at least $N$ on $\operatorname{diag}(Y \times Y)$. We write $H_{0}^{N}=R_{0, N}+R_{1, N}$, where $R_{1, N} \in G_{n-(d / 2)-1,0}(X)$ is the lower order term of $H_{0}^{N}$ and $R_{0, N} \in \widehat{G}_{n-(d / 2), N}(X)$ is the the leading term of $H_{0}^{N}$. Since the full symbol of $R_{0, N}$ has vanishing order at least $N$ at $\operatorname{diag}(Y \times Y)$, even the order of $R_{0, N}$ is high, $R_{0, N}$ still has good regularity properties if $N$ is large. Note that for an element $A \in G_{k, \ell}(X)$, only the leading symbol of $A$ has vanishing order at least $\ell$ on $\operatorname{diag}(Y \times Y)$. Hence even for large $\ell$ we still do not have good regularity property for $A$ in general. That is why we need the space $\widehat{G}_{k, \ell}(X)$.

Remark 4.6. In the following, we will establish $L^{2}$ continuity and regularity properties for $H \in G_{k, \ell}(X)$ under certain assumptions on $k, \ell$ (see Theorem 4.10). Hörmander 25] already established the $L^{2}$ continuity for complex Fourier integral operators but we cannot apply Hörmander's results to our situation. More precisely, from the main results in [25], we have that for $H \in G_{n-(d / 2)-m, 0}(X)$ with $m \in \mathbb{N}$, the map $H: H^{s}(X) \rightarrow H^{s+m}(X)$ is continuous for every $s \in \mathbb{R}$. In this work, we need the regularity property for $H \in G_{n-(d / 2), \ell}(X)$ for $\ell$ large (see Lemma 4.9 and (4.58)) but Hörmander's result can only be applied to $\ell=0$.

From Theorem 4.2 we deduce the following.
Theorem 4.7. Let $H_{1} \in G_{k_{1}, \ell_{1}}(X), H_{2} \in G_{k_{2}, \ell_{2}}(X)$, where $k_{1}, \ell_{1}, k_{2}, \ell_{2} \in \mathbb{R}$. Then,

$$
H_{1} \circ H_{2} \in G_{k_{1}+k_{2}-\left(n-\frac{d}{2}\right), \ell_{1}+\ell_{2}}(X)
$$

Recall that $\|\cdot\|_{s}$ denotes the standard Sobolev norm on $X$ of order $s$.
Theorem 4.8. Let $H \in G_{k, 0}(X)$ with $k \leq n-\frac{d}{2}-1$. Then, there exists $C>0$ such that for any $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\|H u\| \leq C\|u\| . \tag{4.46}
\end{equation*}
$$

Moreover, for every $s \in \mathbb{N}^{*}$, there exist $N_{s} \in \mathbb{N}^{*}$ and $C_{s}>0$ such that for any $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\left\|H^{N_{s}} u\right\|_{s} \leq C_{s}\|u\| \tag{4.47}
\end{equation*}
$$

Proof. Fix $s \in \mathbb{N}^{*}$. By Theorem 4.7, for any $N_{s} \in \mathbb{N}^{*}$ we have $H^{N_{s}} \in$ $G_{N_{s} k-\left(N_{s}-1\right)\left(n-\frac{d}{2}\right), 0}(X)$. Taking now $N_{s} \in \mathbb{N}^{*}$ such that $N_{s} k-\left(N_{s}-1\right)\left(n-\frac{d}{2}\right)<$ $-s-2$. We claim that $H^{N_{s}}(x, y) \in \mathscr{C}^{s}(X \times X)$. Since $H^{N_{s}}$ is smoothing outside
$Y$, we only need to check this property near $Y$. Let $U$ be an open set of $p \in Y$. Fix $\alpha, \beta \in \mathbb{N}^{2 n+1}$ with $|\alpha|+|\beta| \leq s$. Since $H^{N_{s}} \in G_{N_{s} k-\left(N_{s}-1\right)\left(n-\frac{d}{2}\right), 0}(X)$, we have

$$
\partial_{x}^{\alpha} \partial_{y}^{\beta} H^{N_{s}}(x, y) \equiv \int_{0}^{\infty} e^{i t \Phi(x, y)} a(x, y, t)
$$

$a(x, y, t) \in S_{\mathrm{cl}}^{N_{s} k-\left(N_{s}-1\right)\left(n-\frac{d}{2}\right)+|\alpha|+|\beta|}\left(U \times U \times \mathbb{R}_{+}\right)$. Since there exist $C, \hat{C}>0$ such that

$$
\begin{aligned}
\int_{1}^{\infty} e^{-t \operatorname{Im} \Phi(x, y)}|a(x, y, t)| d t & \leq C \int_{1}^{\infty} t^{N_{s} k-\left(N_{s}-1\right)\left(n-\frac{d}{2}\right)+|\alpha|+|\beta|} d t \\
& \leq \hat{C} \int_{1}^{\infty} t^{-2} d t<\infty
\end{aligned}
$$

we conclude that $H^{N_{s}}(x, y) \in \mathscr{C}^{s}(X \times X)$ and (4.47) follows.
We now prove (4.46). We claim that for every $\ell \in \mathbb{N}^{*}$, we have for any $u \in$ $\mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\|H u\|^{2} \leq\left\|\left(H^{*} H\right)^{2^{\ell}} u\right\|^{2^{-\ell}}\|u\|^{2-2^{-\ell}} \tag{4.48}
\end{equation*}
$$

where $H^{*}$ is the adjoint of $H$. We have for any $u \in \mathscr{C}{ }^{\infty}(X)$,

$$
\begin{equation*}
\|H u\|^{2}=(H u, H u)=\left(H^{*} H u, u\right) \leq\left\|H^{*} H u\right\|\|u\| \tag{4.49}
\end{equation*}
$$

and

$$
\begin{equation*}
\left\|H^{*} H u\right\|^{2}=\left(H^{*} H u, H^{*} H u\right)=\left(\left(H^{*} H\right)^{2} u, u\right) \leq\left\|\left(H^{*} H\right)^{2} u\right\|\|u\| . \tag{4.50}
\end{equation*}
$$

We prove (4.48) by induction on $\ell$. From (4.49) and (4.50), we get (4.48) for $\ell=1$. Suppose that (4.48) holds for $\ell \in \mathbb{N}^{*}$. We have for every $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{align*}
\left\|\left(H^{*} H\right)^{2^{\ell}} u\right\|^{2} & =\left(\left(H^{*} H\right)^{2^{\ell}} u,\left(H^{*} H\right)^{2^{\ell}} u\right) \\
& =\left(\left(H^{*} H\right)^{2^{\ell+1}} u, u\right) \leq\left\|\left(H^{*} H\right)^{2^{\ell+1}} u\right\|\|u\|, \tag{4.51}
\end{align*}
$$

From the induction hypothesis and (4.51), we get (4.48) for $\ell+1$.
It is obvious that $H^{*} \in G_{k, 0}$ and hence $H^{*} H \in G_{2 k-\left(n-\frac{d}{2}\right), 0}$. From this observation and (4.47), we deduce that for $\ell$ large, there exists $C>0$ such that for any $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\left\|\left(H^{*} H\right)^{2^{\ell}} u\right\| \leq C\|u\| \tag{4.52}
\end{equation*}
$$

From (4.48) and (4.52), we get (4.46).
Lemma 4.9. Let $H \in \widehat{G}_{k, 2 \ell}(X)$. If $k-\ell \leq-s-2$, for some $s \in \mathbb{N}$, then $H(x, y) \in$ $\mathscr{C}^{s}(X \times X)$.

Proof. Since $H$ is smoothing away $Y$, we only need to prove that $H(x, y)$ is in $\mathscr{C}^{s}$ near $Y$. Let $p \in Y$ and $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ be local coordinates as in the discussion
before Theorem 3.28 defined in an open set $U$ of $p$. We will use the same notations as in the beginning of Sec.4.1. On $U$, write

$$
H(x, y)=\int_{0}^{\infty} e^{i \Phi(x, y) t} a(x, y, t) d t+F(x, y)
$$

$F(x, y) \in \mathscr{C}^{\infty}(U \times U), a \in S_{\mathrm{cl}}^{k}\left(U \times U \times \mathbb{R}_{+}\right)$. Let $u \in \mathscr{C}_{c}^{\infty}(U)$. Since $H u$ is $G$ invariant, $H u$ is independent of $x^{\prime}$ and hence on $U$,

$$
(H u)(x)=(H u)\left(x^{\prime \prime}\right)=\int_{0}^{\infty} e^{i \Phi(x, y) t} a\left(x^{\prime \prime}, y, t\right) u(y) d v(y) d t+\int F\left(x^{\prime \prime}, y\right) u(y) d v(y)
$$

Hence,

$$
H(x, y) \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} a\left(x^{\prime \prime}, y, t\right) d t \text { on } U \times U
$$

We are going to prove the lemma by induction over $\ell \in \mathbb{N}$. Let $\ell=0$. Let $P$ and $Q$ be differential operators on $U$ with ord $(P)+\operatorname{ord}(Q) \leq s$. Then the symbol of $P H Q$ is of order $\leq k+\operatorname{ord}(P)+\operatorname{ord}(Q) \leq k+s \leq-2$ and hence the integral over $t$ converges. We get that the lemma holds for $\ell=0$. We assume that the lemma holds for some $\ell=\ell_{0} \in \mathbb{N}$. We are going to prove that the lemma holds for $\ell=\ell_{0}+1$. Since all the asymptotics of $a$ have vanishing order at least $2 \ell_{0}+2$ on $\operatorname{diag}(Y \times Y)$, we may assume for simplicity that $a\left(x^{\prime \prime}, y, t\right)=t^{k} b\left(x^{\prime \prime}, y\right)$ and $b$ has vanishing order at least $2 \ell_{0}+2$ on $\operatorname{diag}(Y \times Y)$. The Malgrange preparation theorem entails that there exist a neighborhood $U_{0} \subset U$ of $p$ and functions $f, g, \hat{\Phi}, \hat{b} \in \mathscr{C}^{\infty}\left(U_{0} \times U_{0}\right)$, such that $\hat{\Phi}$ and $\hat{b}$ are independent of $y_{2 n+1}$ and we have on $U_{0} \times U_{0}$,

$$
\begin{align*}
\Phi(x, y) & =f(x, y)\left(y_{2 n+1}+\hat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right) \\
b\left(x^{\prime \prime}, y\right) & =g\left(x^{\prime \prime}, y\right)\left(y_{2 n+1}+\hat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right)+\hat{b}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right) \tag{4.53}
\end{align*}
$$

We claim that $g(x, y)$ and $\hat{b}$ have vanishing order at least $2 \ell_{0}+1$ and $2 \ell_{0}+2$ on $\operatorname{diag}(Y \times Y)$, respectively. For every $j=1, \ldots, 2 \ell_{0}+1$, we have

$$
\begin{equation*}
\frac{\partial^{j} b}{\partial y_{2 n+1}^{j}}=j \frac{\partial^{j-1} g}{\partial y_{2 n+1}^{j-1}}+\frac{\partial^{j} g}{\partial y_{2 n+1}^{j}}\left(y_{2 n+1}+\hat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right) \tag{4.54}
\end{equation*}
$$

Taking $j=2 \ell_{0}+1$ in (4.54) and using that $\left.\left(y_{2 n+1}+\hat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right)\right|_{\operatorname{diag}(Y \times Y)}=0$, we conclude that $\partial_{y_{2 n+1}}^{2 \ell_{0}} g$ vanishes on $\operatorname{diag}(Y \times Y)$. Taking $j=2 \ell_{0}$ in (4.54) we conclude that $\partial_{y_{2 n+1}}^{2 \ell_{0}-1} g$ has vanishing order at least 2 on $\operatorname{diag}(Y \times Y)$. Continuing in a similar way, we can show that for every $j=1, \ldots, 2 \ell_{0}+1, \partial_{y_{2 n+1}}^{j-1} g$ has vanishing order at least $2 \ell_{0}+2-j$ on $\operatorname{diag}(Y \times Y)$. Now we have from (4.53),

$$
\begin{equation*}
\frac{\partial b}{\partial y_{2 n+1}}=g+\frac{\partial g}{\partial y_{2 n+1}}\left(y_{2 n+1}+\hat{\Phi}\left(x^{\prime \prime}, \stackrel{\circ}{y}^{\prime \prime}\right)\right) \tag{4.55}
\end{equation*}
$$

Since $\partial_{y_{2 n+1}} g$ has vanishing order at least $2 \ell_{0}$ on $\operatorname{diag}(Y \times Y)$, we deduce that $g(x, y)$ has vanishing order at least $2 \ell_{0}+1$ on diag $(Y \times Y)$. Thus, $\hat{b}=b-g\left(y_{2 n+1}+\hat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right)$ has vanishing order at least $2 \ell_{0}+2$ at diag $(Y \times Y)$. The claim follows.

By using integration by parts, we have

$$
\begin{aligned}
\int_{0}^{\infty} e^{i t \Phi(x, y)} t^{k} b\left(x^{\prime \prime}, y\right) d t & \equiv \int_{0}^{\infty} e^{i t \Phi(x, y)} t^{k}\left(g(x, y)\left(y_{2 n+1}+\hat{\Phi}\left(x^{\prime \prime}, \dot{y}^{\prime \prime}\right)\right)+\hat{b}\left(x^{\prime \prime}, \hat{y}\right)\right) d t \\
& \equiv I_{1}+I_{2}, \quad \text { where } \\
I_{1}(x, y) & \equiv \int_{0}^{\infty} e^{i t \Phi(x, y)} k i t^{k-1} \frac{g}{f}(x, y) d t \\
I_{2}(x, y) & \equiv \int_{0}^{\infty} e^{i t \Phi(x, y)} t^{k} \hat{b}\left(x^{\prime \prime}, \stackrel{y}{y}\right) d t
\end{aligned}
$$

Since $g$ has vanishing order at least $2 \ell_{0}+1$ and $k-1-\ell_{0} \leq-s-2$, we deduce by the induction assumption that $I_{1}(x, y) \in \mathscr{C}^{s}\left(U_{0} \times U_{0}\right)$. Since $\hat{b}\left(x^{\prime \prime}, y\right)$ has vanishing order at least $2 \ell_{0}+2$ on $\operatorname{diag}(Y \times Y)$, there exist $C, \hat{C}>0$ such that we have

$$
\begin{align*}
\left|\hat{b}\left(x^{\prime \prime}, \dot{y}\right)\right| & \leq C\left|\left(x^{\prime \prime}, y\right)-\left(\left(0, \underline{x}^{\prime \prime}\right),\left(0, \underline{x}^{\prime \prime}\right)\right)\right|^{2 \ell_{0}+2} \\
& \leq \hat{C}\left(\left|\hat{x}^{\prime \prime}\right|^{2}+\left|\dot{x}^{\prime \prime}-\grave{y}^{\prime \prime}\right|^{2}+\left|\hat{y}^{\prime \prime}\right|^{2}\right)^{\ell_{0}+1} \tag{4.56}
\end{align*}
$$

where $\underline{\dot{x}}^{\prime \prime}=\left(x_{2 d+1}, \ldots, x_{2 n}\right)$. From (3.93) and (4.56), there exist $c, \hat{C}>0, C_{1}, C_{2}>$ 0 such that

$$
\begin{aligned}
\left|\int_{1}^{\infty} e^{i t \Phi(x, y)} t^{k} \hat{b}\left(x^{\prime \prime}, \dot{y}\right) d t\right| \leq & \int_{1}^{\infty} e^{-t \operatorname{Im} \Phi} t^{k}\left|\hat{b}\left(x^{\prime \prime}, \dot{y}\right)\right| d t \\
\leq & \int_{1}^{\infty} e^{-t c\left(\left|\hat{x}^{\prime \prime}\right|^{2}+\left|\dot{x}^{\prime \prime}-\dot{y}^{\prime \prime}\right|^{2}+\left|\hat{y}^{\prime \prime}\right|\right)^{2}} t^{k} \\
& \times \hat{C}\left(\left|\hat{x}^{\prime \prime}\right|^{2}+\left|\dot{x}^{\prime \prime}-\dot{y}^{\prime \prime}\right|^{2}+\left|\hat{y}^{\prime \prime}\right|^{2}\right)^{\ell_{0}+1} d t \\
\leq & C_{1} \int_{1}^{\infty} t^{k-\ell_{0}-1} d t \leq C_{2} \int_{1}^{\infty} t^{-s-2} d t<\infty
\end{aligned}
$$

Let $\alpha, \beta \in \mathbb{N}^{2 n+1}$ with $|\alpha|+|\beta| \leq s$. We can repeat the procedure above with minor changes and deduce that

$$
\int_{1}^{\infty}\left|\partial_{x}^{\alpha} \partial_{y}^{\beta}\left(e^{i t \Phi(x, y)} t^{k} \hat{b}\left(x^{\prime \prime}, \stackrel{y}{y}\right)\right)\right| d t \leq C_{3} \int_{1}^{\infty} t^{k-\ell_{0}-1+|\alpha|+|\beta|} \leq C_{4} \int_{1}^{\infty} t^{-2} d t<\infty
$$

where $C_{3}, C_{4}>0$ are constants. Hence $I_{2}(x, y) \in \mathscr{C}^{s}\left(U_{0} \times U_{0}\right)$. The assertion holds thus for $\ell=\ell_{0}+1$ and the lemma follows.

In the proof of our main result, we need the following.
Theorem 4.10. Let $H \in G_{n-\frac{d}{2}, 1}(X)$. Then there exists $C>0$ such that for any $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\|H u\| \leq C\|u\| \tag{4.57}
\end{equation*}
$$

Moreover, for every $s \in \mathbb{N}^{*}$, there exist $N_{s} \in \mathbb{N}^{*}$ and $C_{s}>0$ such that for any $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\left\|H^{N_{s}} u\right\|_{s} \leq C_{s}\|u\| . \tag{4.58}
\end{equation*}
$$

Proof. From Theorem 4.2, we see that for every $N \in \mathbb{N}^{*}$, we have

$$
\begin{align*}
H^{N} & =H_{1, N}+H_{2, N} \\
H_{1, N} & \in \widehat{G}_{n-\frac{d}{2}, N}(X), \quad H_{2, N} \in G_{n-\frac{d}{2}-1,0}(X) \tag{4.59}
\end{align*}
$$

Fix $s \in \mathbb{N}^{*}$. Due to Lemma 4.9 there exists $N \gg 1$ such that $H_{1, N}(x, y) \in \mathscr{C}^{s}(X \times X)$ and for every $j \in \mathbb{N}^{*}$, there exists $C_{j}>0$ such that for every $u \in \mathscr{C}^{\infty}(X)$ we have

$$
\begin{equation*}
\left\|H_{1, N}^{j} u\right\|_{s} \leq C_{j}\|u\| \tag{4.60}
\end{equation*}
$$

Since $H_{2, N} \in G_{n-\frac{d}{2}-1,0}(X)$, Theorem 4.8 shows that there exist $K_{s} \in \mathbb{N}^{*}$ and $\widehat{C}_{s}>0$ so that

$$
\begin{equation*}
\left\|H_{2, N}^{K_{s}} u\right\|_{s} \leq \widehat{C}_{s}\|u\|, \quad \text { for } u \in \mathscr{C}^{\infty}(X) \tag{4.61}
\end{equation*}
$$

We have

$$
\begin{equation*}
H^{N K_{s}}=\left(H_{1, N}+H_{2, N}\right)^{K_{s}} \tag{4.62}
\end{equation*}
$$

Let $A:=H_{1, N}^{a_{1}} H_{2, N}^{b_{1}} \cdots H_{1, N}^{a_{p}} H_{2, N}^{b_{p}}$, where $a_{j}, b_{j} \in \mathbb{N}, j=1, \ldots, p, \sum_{j=1}^{p}\left(a_{j}+b_{j}\right)=$ $K_{s}$. We claim that

$$
\begin{equation*}
A: L^{2}(X) \rightarrow H^{s}(X) \text { is continuous. } \tag{4.63}
\end{equation*}
$$

If $a_{1}=a_{2}=\ldots=a_{p}=0$, then, $A=H_{2, N}^{K_{s}}$. From (4.61), we get (4.63). Assume that $a_{j_{0}} \neq 0$, for some $j_{0} \in\{1,2, \ldots, p\}$. Let $\Lambda_{s}$ be a classical elliptic pseudodifferential operator on $X$ of order $s$ with inverse $\Lambda_{-s}$. By the complex stationary phase formula of Melin-Sjöstrand, we have $\Lambda_{s} \circ H_{2, N} \circ \Lambda_{-s} \in G_{n-\frac{d}{2}-1,0}(X)$. From this observation and (4.46), we see that

$$
\begin{equation*}
H_{2, N}: H^{s}(X) \rightarrow H^{s}(X) \text { is continuous. } \tag{4.64}
\end{equation*}
$$

By (4.60) and (4.64) we deduce that

$$
\begin{equation*}
H_{1, N} \circ H_{2, N} \text { and } H_{2, N} \circ H_{1, N}: L^{2}(X) \rightarrow H^{s}(X) \text { are continuous. } \tag{4.65}
\end{equation*}
$$

From (4.65), we get the claim (4.63). From (4.62) and (4.63), we get (4.58) with $N_{s}=N K_{s}$. Using (4.58) we can repeat the proof of (4.46) and obtain (4.57).

Let $H \in G_{n-\frac{d}{2}, 1}(X)$. From (4.57), we can extend $I-H$ to a bounded operator in

$$
I-H: L^{2}(X) \rightarrow L^{2}(X)
$$

Theorem 4.11. Let $H \in G_{n-\frac{d}{2}, 1}(X)$ and extend $I-H$ to a bounded operator in

$$
I-H: L^{2}(X) \rightarrow L^{2}(X)
$$

by density. Then $\operatorname{Ker}(I-H)$ is a finite-dimensional subspace of $\mathscr{C}^{\infty}(X)$ and there exists $C>0$ such that

$$
\begin{equation*}
\|(I-H) u\| \geq C\|u\|, \quad \text { for any } u \in L^{2}(X), \quad u \perp \operatorname{Ker}(I-H) \tag{4.66}
\end{equation*}
$$

Proof. Fix $s \in \mathbb{N}^{*}$. Theorem 4.10 shows that we can extend $H^{N_{s}}$ to a bounded operator

$$
\begin{equation*}
H^{N_{s}}: L^{2}(X) \rightarrow H^{s}(X) \tag{4.67}
\end{equation*}
$$

by density. Now, let $u \in \operatorname{Ker}(I-H)$. Then,

$$
\begin{equation*}
\left(I-H^{N_{s}}\right) u=\left(I+H+\cdots+H^{N_{s}-1}\right)(I-H) u=0 \tag{4.68}
\end{equation*}
$$

and hence $u=H^{N_{s}} u \in H^{s}(X)$. Since $s$ is arbitrary, we deduce that $u \in \mathscr{C}^{\infty}(X)$. Moreover, from (4.68), we can apply Rellich's theorem and conclude that $\operatorname{Ker}(I-H)$ is a finite-dimensional subspace of $\mathscr{C}^{\infty}(X)$. Since the argument is standard, we omit the details.

We now prove (4.66). Assume that (4.66) is not true. For every $j \in \mathbb{N}^{*}$ we can find $u_{j} \in L^{2}(X)$ with $u_{j} \perp \operatorname{Ker}(I-H)$ and $\left\|u_{j}\right\|=1$ such that

$$
\begin{equation*}
\left\|(I-H) u_{j}\right\| \leq \frac{1}{j} \tag{4.69}
\end{equation*}
$$

Put $v_{j}:=(I-H) u_{j}$. We have for any $j \in \mathbb{N}^{*}$,

$$
\begin{equation*}
\left(I-H^{N_{s}}\right) u_{j}=\left(I+H+\ldots+H^{N_{s}-1}\right) v_{j} \tag{4.70}
\end{equation*}
$$

From (4.57) and (4.69) we see that there exists $C>0$ such that

$$
\begin{equation*}
\left\|\left(I+H+\ldots+H^{N_{s}-1}\right) v_{j}\right\| \leq \frac{C}{j}, \quad j \in \mathbb{N}^{*} \tag{4.71}
\end{equation*}
$$

By (4.67) and since $\left\|u_{j}\right\|=1$ we conclude that there exists $\widehat{C}>0$ such that for any $j \in \mathbb{N}^{*}$,

$$
\begin{equation*}
\left\|H^{N_{s}} u_{j}\right\|_{s} \leq \widehat{C} \tag{4.72}
\end{equation*}
$$

By Rellich's theorem, we can find a subsequence $H^{N_{s}} u_{j_{k}}, 1 \leq j_{1}<j_{2}<\ldots$, such that $H^{N_{s}} u_{j_{k}}$ converges to some $u$ in $L^{2}(X)$ as $k \rightarrow \infty$. From this observation, (4.70) and (4.71), we deduce that $u_{j_{k}}$ converges to $u$ in $L^{2}(X)$ with $\|u\|=1$ as $k \rightarrow \infty$. By (4.57) and (4.69), we get $u \in \operatorname{Ker}(I-H)$. Since $u_{j_{k}} \perp \operatorname{Ker}(I-H)$ for every $k$, we have $u \perp \operatorname{Ker}(I-H)$. We get a contradiction and (4.66) follows.

### 4.2. The distribution kernels of $\sigma$ and $\sigma^{*} \sigma$; proof of Theorem 1.5

We are now ready to study the distribution kernels of $\sigma$ and $\sigma \sigma^{*}$ in (4.1). We will use the same notations as before. Let $\mathscr{L}_{X_{G}, q}$ be the Levi form on $X_{G}$ at $q \in X_{G}$ induced naturally from $\mathscr{L}$. The Hermitian metric $g^{\mathbb{C} T X}$ on $\mathbb{C} T X$ restricts to a metric on $T^{1,0} X$ which in turn induces a Hermitian metric on $T^{1,0} X_{G}$. Let $\mu_{d+1}, \ldots, \mu_{n}$ be the eigenvalues of $\mathscr{L}_{X_{G}, q}$ with respect to this Hermitian metric. We set

$$
\begin{equation*}
\operatorname{det} \mathscr{L}_{X_{G}, q}=\mu_{d+1} \ldots \mu_{n} \tag{4.73}
\end{equation*}
$$

Recall that $\pi: Y=\mu^{-1}(0) \rightarrow X_{G}$ is the natural projection. Let

$$
\begin{equation*}
S_{X_{G}}: L^{2}\left(X_{G}\right) \rightarrow \operatorname{Ker} \bar{\partial}_{b, X_{G}}=H_{b}^{0}\left(X_{G}\right) \tag{4.74}
\end{equation*}
$$

be the Szegő projection on $X_{G}$ (cf. (2.14)). Since $X_{G}$ is assumed to be strictly pseudoconvex and $\bar{\partial}_{b, X_{G}}$ has closed range in $L^{2}$ on $X_{G}, S_{X_{G}}$ is smoothing away the diagonal (see [8, [27, Theorem 1.2], [30, Theorem 4.7]). Hence, for any $x, y \in Y$ with $\pi(x) \neq \pi(y)$, there are open neighborhoods $U$ of $\pi(x)$ in $X_{G}$ and $U_{1}$ of $\pi(y)$ in $X_{G}$ such that for all $\widehat{\chi} \in \mathscr{C}_{0}^{\infty}(U), \widetilde{\chi} \in \mathscr{C}_{0}^{\infty}\left(U_{1}\right)$, we have

$$
\begin{equation*}
\widehat{\chi} S_{X_{G}} \widetilde{\chi} \equiv 0 \quad \text { on } X_{G} \times X_{G} . \tag{4.75}
\end{equation*}
$$

We will use the same notations as in Sec. 4.1 Fix $p \in Y$ and let $x=$ $\left(x_{1}, \ldots, x_{2 n+1}\right)$ be the local coordinates and $\Omega_{3} \subset \mathbb{R}^{2 n+1-2 d}$ be an open set as in the discussion at the beginning of Sec. 4.1. From now on, we identify $\underline{x}^{\prime \prime}$ as local coordinates of $X_{G}$ near $q:=\pi(p) \in X_{G}$ and we identify $W:=\Omega_{3}$ with an neighborhood of $\pi(p)$ in $X_{G}$. Note that from (3.90), on $\mu^{-1}(0) / G$,

$$
\begin{align*}
\omega_{0, G}(x)= & (1+O(|x|)) d x_{2 n+1}+\sum_{j=d+1}^{n} 2 \mu_{j}\left(x_{2 j} d x_{2 j-1}-x_{2 j-1} d x_{2 j}\right) \\
& +\sum_{j=2 d+1}^{2 n} b_{j} x_{2 n+1} d x_{j}+O\left(|x|^{2}\right) \tag{4.76}
\end{align*}
$$

From (3.79), (4.76), we see that $\mu_{j}, j=d+1, \ldots, n$ in (4.76) are the same $\mu_{j}$ in (3.79) and $c_{j}=b_{2 j-1}-i b_{2 j}$. From this observation and applying [8] (more precisely Theorems 3.223 .24 to $\mu^{-1}(0) / G$, we have

$$
\begin{equation*}
S_{X_{G}}\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \equiv \int_{0}^{\infty} e^{i \varphi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) t} \beta\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}, t\right) d t \quad \text { on } W \times W \tag{4.77}
\end{equation*}
$$

where $\beta\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}, t\right) \in S_{\mathrm{cl}}^{n-d}\left(W \times W \times \mathbb{R}_{+}\right)$with

$$
\begin{equation*}
\beta_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=\frac{1}{2} \pi^{-(n-d)-1}\left|\operatorname{det} \mathscr{L}_{X_{G}, \underline{x}^{\prime \prime}}\right|, \quad \underline{x}^{\prime \prime} \in W \tag{4.78}
\end{equation*}
$$

and $\varphi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \in \mathscr{C}^{\infty}(W \times W)$ with

$$
\begin{align*}
& d_{\underline{x}^{\prime \prime}} \varphi\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=-d_{\underline{y}^{\prime \prime}} \varphi\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=-\lambda\left(\underline{x}^{\prime \prime}\right) \omega_{0, G}\left(\underline{x}^{\prime \prime}\right), \quad \lambda\left(\underline{x}^{\prime \prime}\right)>0 \\
& \operatorname{Im} \varphi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \geq c \sum_{j=2 d+1}^{2 n}\left|x_{j}-y_{j}\right|^{2}, \quad \text { for some } c>0 \\
& \bar{\partial}_{b, \underline{x}^{\prime \prime}} \varphi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \text { vanishes to infinite order at } \underline{x}^{\prime \prime}=\underline{y}^{\prime \prime} \\
& \varphi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)=-x_{2 n+1}+y_{2 n+1}+i \sum_{j=d+1}^{n}\left|\mu_{j}\right|\left|z_{j}-w_{j}\right|^{2} \\
&+\sum_{j=d+1}^{n} i \mu_{j}\left(\bar{z}_{j} w_{j}-z_{j} \bar{w}_{j}\right)  \tag{4.79}\\
&+\sum_{j=d+1}^{n} \frac{1}{2}\left(b_{2 j-1}-i b_{2 j}\right)\left(-z_{j} x_{2 n+1}+w_{j} y_{2 n+1}\right) \\
&+\sum_{j=d+1}^{n} \frac{1}{2}\left(b_{2 j-1}+i b_{2 j}\right)\left(-\bar{z}_{j} x_{2 n+1}+\bar{w}_{j} y_{2 n+1}\right) \\
&+\left(x_{2 n+1}-y_{2 n+1}\right) r\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)+O\left(\left|\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)\right|^{3}\right)
\end{align*}
$$

where $r\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \in \mathscr{C}^{\infty}(W \times W), r(0,0)=0, z_{j}=x_{2 j-1}+i x_{2 j}, j=d+1, \ldots, n$.
From Theorem 3.28, it is not difficult to see that the phase function $\Phi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)$ in Theorem 3.25 satisfies (4.79). Hence, there is a function $h \in \mathscr{C}^{\infty}(W \times W)$ with $h\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) \neq 0$ for any $\underline{x}^{\prime \prime} \in W$, such that $\varphi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)-h\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \Phi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)$ vanishes to infinite order at $\underline{x}^{\prime \prime}=\underline{y}^{\prime \prime}$ (see Theorem 3.24). We can replace the phase $\varphi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)$ by $\Phi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right)$ and we have

$$
\begin{equation*}
S_{X_{G}}\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \equiv \int_{0}^{\infty} e^{i \Phi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) t} \beta\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}, t\right) d t \quad \text { on } W \times W \tag{4.80}
\end{equation*}
$$

Theorem 4.12. If $y \notin Y$, then for any open neighborhood $D$ of $y$ with $\bar{D} \cap Y=\emptyset$, we have

$$
\begin{equation*}
\sigma \equiv 0 \quad \text { on } X_{G} \times D \tag{4.81}
\end{equation*}
$$

Let $x_{0}, y_{0} \in Y$. If $\pi\left(x_{0}\right) \neq \pi\left(y_{0}\right)$, then there are open neighborhoods $U_{G}$ of $\pi\left(x_{0}\right)$ in $X_{G}$ and $U_{1}$ of $y_{0}$ in $X$ such that

$$
\begin{equation*}
\sigma \equiv 0 \quad \text { on } U_{G} \times U_{1} \tag{4.82}
\end{equation*}
$$

Let $p \in Y$ and let $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ be the local coordinates, and $U$ an open neighborhood of $p$ as at the discussion in the beginning of Sec.4.1. Then under the
notations at (3.76), (4.77), there exists $\alpha\left(\underline{x}^{\prime \prime}, y^{\prime \prime}, t\right) \in S_{\mathrm{cl}}^{n-\frac{3}{4} d}\left(W \times U \times \mathbb{R}_{+}\right)$such that

$$
\begin{align*}
\sigma\left(\underline{x}^{\prime \prime}, y\right) \equiv & \int_{0}^{\infty} e^{i \Phi\left(\underline{x}^{\prime \prime}, y^{\prime \prime}\right) t} \alpha\left(\underline{x}^{\prime \prime}, y^{\prime \prime}, t\right) d t \quad \text { on } W \times U, \\
\alpha_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)= & 2^{-n+2 d-1} \pi^{\frac{d}{2}-n-1} \frac{1}{\sqrt{V_{\text {eff }}\left(\underline{x}^{\prime \prime}\right)}}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right|  \tag{4.83}\\
& \times\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right|^{-\frac{3}{4}}, \quad \text { for } \underline{x}^{\prime \prime} \in W .
\end{align*}
$$

Proof. By Theorem 3.20, $S_{G}$ is smoothing away $Y$, which implies 4.81) from Theorem 3.27, Let $x_{0}, y_{0} \in Y$. Assume that $\pi\left(x_{0}\right) \neq \pi\left(y_{0}\right)$. Let $V_{1}$ be a $G$-invariant neighborhood of $y_{0}$ and $x_{0} \notin V_{1}$. We have

$$
S_{G}(x, y)=\int_{G} S_{\leq \lambda_{0}}(x, g \circ y) d \mu(g)
$$

where $\lambda_{0}>0$ is as in Theorem 3.19, Since $S_{\leq \lambda_{0}}$ is smoothing away the diagonal near $Y$, for any neighborhood $U_{1}$ of $x_{0}$ in $X$ with $\bar{U}_{1} \cap V_{1}=\emptyset$, we have

$$
\begin{equation*}
S_{G} \equiv 0 \quad \text { on } U_{1} \times V_{1} \tag{4.84}
\end{equation*}
$$

Let $\widehat{V}_{G}:=\left\{\pi(y) \in X_{G}: y \in V_{1} \cap Y\right\}$. By (4.75) there is an open set $\widehat{U}_{G}$ of $\pi\left(x_{0}\right)$ in $X_{G}$ such that

$$
\begin{equation*}
S_{X_{G}} \equiv 0 \quad \text { on } \widehat{U}_{G} \times \widehat{V}_{G} \tag{4.85}
\end{equation*}
$$

The definition (4.1) of $\sigma$ and Theorem 3.27, relations (4.84) and (4.85) yield (4.82).
Fix $u=\left(u_{1}, \ldots, u_{2 n+1}\right) \in Y \cap U$. In view of (4.81) and (4.82), we only need to show that (4.83) holds near $u$. We may assume that $u=$ $\left(0, \ldots, 0, u_{2 d+1}, \ldots, u_{2 n+1}\right)=\underline{u}^{\prime \prime}$. Let $U_{2}$ be a small neighborhood of $u$. Let $\chi\left(\underline{x}^{\prime \prime}\right) \in \mathscr{C}_{0}^{\infty}\left(\Omega_{3}\right)$. By (4.5) we can extend $\chi\left(\underline{x}^{\prime \prime}\right)$ to $G \cdot \Omega_{3}$ by setting $\chi\left(g \cdot \underline{x}^{\prime \prime}\right):=\chi\left(\underline{x}^{\prime \prime}\right)$ for every $g \in G$. Assume that $\chi=1$ on some neighborhood of $\bar{U}_{2}$. Let $\chi_{1} \in \mathscr{C}_{0}^{\infty}\left(X_{G}\right)$ with $\chi_{1}=1$ on some neighborhood of $\pi\left(U_{2} \cap Y\right) \subset X_{G}$ and $\operatorname{supp} \chi_{1} \subset \pi(\{x \in$ $Y: \chi(x)=1\}$ ). We have by (4.1),

$$
\begin{align*}
\chi_{1} \sigma= & \chi_{1} S_{X_{G}} \circ E \circ \iota_{G} \circ f_{G} \circ \iota^{*} \circ S_{G} \\
= & \chi_{1} S_{X_{G}} \circ E \circ \iota_{G} \circ f_{G} \circ \iota^{*} \circ \chi S_{G} \\
& +\chi_{1} S_{X_{G}} \circ E \circ \iota_{G} \circ f_{G} \circ \iota^{*} \circ(1-\chi) S_{G} . \tag{4.86}
\end{align*}
$$

If $u \in Y$ but $u \notin\{x \in X: \chi(x)=1\}$, then $\pi(u) \notin \operatorname{supp} \chi_{1}$. From this observation, Theorem 3.27 and 4.75), we get

$$
\begin{equation*}
\chi_{1} S_{X_{G}} \circ E \circ \iota_{G} \circ f_{G} \circ \iota^{*} \circ(1-\chi) S_{G} \equiv 0 \quad \text { on } X_{G} \times X \tag{4.87}
\end{equation*}
$$

From (4.86) and (4.87), we get

$$
\begin{equation*}
\chi_{1} \sigma \equiv \chi_{1} S_{X_{G}} \circ E \circ \iota_{G} \circ f_{G} \circ \iota^{*} \circ \chi S_{G} \quad \text { on } X_{G} \times X \tag{4.88}
\end{equation*}
$$

From Theorem 3.25, (4.80) and (4.88), we can check that on $W \times U$,

$$
\begin{align*}
\chi_{1}\left(\underline{x}^{\prime \prime}\right) \sigma\left(\underline{x}^{\prime \prime}, y\right) \equiv & \int_{0}^{\infty} \int e^{i \Phi\left(\underline{x}^{\prime \prime}, \underline{v}^{\prime \prime}\right) t} \beta\left(\underline{x}^{\prime \prime}, \underline{v}^{\prime \prime}, t\right) E \circ \\
& \times\left(\int_{0}^{\infty} \chi\left(\underline{v}^{\prime \prime}\right) f_{G}\left(\underline{v}^{\prime \prime}\right) e^{i \Phi\left(\underline{v}^{\prime \prime}, y\right) s} a\left(\underline{v}^{\prime \prime}, y, s\right) d s\right) d v\left(\underline{v}^{\prime \prime}\right) d t \tag{4.89}
\end{align*}
$$

From the asymptotic formula of Melin-Sjöstrand [45, (2.28)], we have

$$
\begin{equation*}
E \circ\left(\int_{0}^{\infty} \chi\left(\underline{v}^{\prime \prime}\right) f_{G}\left(\underline{v}^{\prime \prime}\right) e^{i \Phi\left(\underline{v}^{\prime \prime}, y\right) s} a\left(\underline{v}^{\prime \prime}, y, s\right) d s\right) \equiv \int_{0}^{\infty} e^{i \Phi\left(\underline{v}^{\prime \prime}, y\right) s} b\left(\underline{v}^{\prime \prime}, y, s\right) d s \tag{4.90}
\end{equation*}
$$

where $b\left(\underline{v}^{\prime \prime}, y, s\right) \in S_{\mathrm{cl}}^{n-\frac{d}{2}-\frac{d}{4}}\left(W \times U \times \mathbb{R}_{+}\right)$with

$$
b_{0}\left(\underline{v}^{\prime \prime}, \underline{v}^{\prime \prime}\right)=\sigma_{E}^{0}\left(\underline{v}^{\prime \prime},-\omega_{0, G}\left(\underline{v}^{\prime \prime}\right)\right) \chi\left(\underline{v}^{\prime \prime}\right) f_{G}\left(\underline{v}^{\prime \prime}\right) a_{0}\left(\underline{v}^{\prime \prime}, \underline{v}^{\prime \prime}\right),
$$

where $b_{0}$ denotes the leading term of $b$ and $\sigma_{E}^{0}$ denotes the principal symbol of $E$, $a_{0}$ is as in (3.85). Since $\sigma_{E}^{0}(x, \xi)=|\xi|^{-\frac{d}{4}}$ and $\left|\omega_{0}\right|=1$, we have

$$
\begin{equation*}
b_{0}\left(\underline{v}^{\prime \prime}, \underline{v}^{\prime \prime}\right)=\chi\left(\underline{v}^{\prime \prime}\right) f_{G}\left(\underline{v}^{\prime \prime}\right) a_{0}\left(\underline{v}^{\prime \prime}, \underline{v}^{\prime \prime}\right) . \tag{4.91}
\end{equation*}
$$

From Theorem 4.3, (4.89) and (4.90), we get that there exists $\alpha\left(\underline{x}^{\prime \prime}, y^{\prime \prime}, t\right) \in$ $S_{\mathrm{cl}}^{n-\frac{3}{4} d}\left(W \times U \times \mathbb{R}_{+}\right)$such that

$$
\begin{equation*}
\sigma\left(\underline{x}^{\prime \prime}, y\right) \equiv \int_{0}^{\infty} e^{i \Phi\left(\underline{x}^{\prime \prime}, y^{\prime \prime}\right) t} \alpha\left(\underline{x}^{\prime \prime}, y^{\prime \prime}, t\right) d t \quad \text { on } W \times U \tag{4.92}
\end{equation*}
$$

Now, we compute $\alpha_{0}$ at $\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)$. From (1.19), (3.85), (4.39), (4.78) and (4.91), we have that on $\left\{\underline{x}^{\prime \prime} \in W ; \chi\left(\underline{x}^{\prime \prime}\right)=1\right\}$,

$$
\begin{align*}
\alpha_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)= & 2^{-n+1} \pi^{n-d+1}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right|^{-1}\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right| b_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) \beta_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) \\
= & 2^{-n+1} \pi^{n-d+1}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right|^{-1}\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right| f_{G}\left(\underline{x}^{\prime \prime}\right) a_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) \beta_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right) \\
= & 2^{-n+1} \pi^{n-d+1}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right|^{-1}\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right|\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right|^{-\frac{1}{4}} \sqrt{V_{\mathrm{eff}}\left(\underline{x}^{\prime \prime}\right)} \\
& \times 2^{d-1} \frac{1}{V_{\text {eff }}\left(\underline{x}^{\prime \prime}\right)} \pi^{-n-1+\frac{d}{2}}\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right|^{-\frac{1}{2}}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right| \\
& \times \frac{1}{2} \pi^{-(n-d)-1}\left|\operatorname{det} \mathscr{L}_{X_{G}, \underline{x}^{\prime \prime}}\right| \\
= & 2^{-n+2 d-1} \pi^{\frac{d}{2}-n-1} \frac{1}{\sqrt{V_{\text {eff }}\left(\underline{x}^{\prime \prime}\right)}}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right|\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right|^{-\frac{3}{4}} . \tag{4.93}
\end{align*}
$$

Here we used the fact that $\left|\operatorname{det} \mathscr{L}_{X_{G}, \underline{x}^{\prime \prime}}\right|=2^{d}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right|\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right|^{-1}$. From (4.92) and (4.93), we get (4.83).

Let $\sigma^{*}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow \mathscr{D}^{\prime}(X)$ be the formal adjoint of $\sigma$. From Theorem 4.12 we deduce that

$$
\begin{equation*}
\sigma^{*}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)^{G} \tag{4.94}
\end{equation*}
$$

Let

$$
\begin{align*}
& A_{1}:=\sigma^{*} \sigma: \mathscr{C}^{\infty}(X) \rightarrow H_{b}^{0}(X)^{G} \\
& A_{2}:=\sigma \sigma^{*}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow H_{b}^{0}\left(X_{G}\right) \tag{4.95}
\end{align*}
$$

Let $A_{1}(x, y)$ and $A_{2}(x, y)$ be the distribution kernels of $A_{1}$ and $A_{2}$, respectively. In view of Theorems 4.2, 4.3 we can repeat the proof of Theorem 4.12 with minor changes and deduce the following two theorems.

Theorem 4.13. With the notations used above, if $y \notin Y$, then for any neighborhood $D$ of $y$ with $\bar{D} \cap Y=\emptyset$, we have

$$
\begin{equation*}
A_{1} \equiv 0 \quad \text { on } X \times D \tag{4.96}
\end{equation*}
$$

Let $x, y \in Y$. If $\pi(x) \neq \pi(y)$, then there are neighborhoods $D_{1}$ of $x$ in $X$ and $D_{2}$ of $y$ in $X$ such that

$$
\begin{equation*}
A_{1} \equiv 0 \quad \text { on } D_{1} \times D_{2} \tag{4.97}
\end{equation*}
$$

Let $p \in Y$ and let $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ be the local coordinates as in the discussion in the beginning of Sec. 4.1. Then there exists an open neighborhood $U$ of $p$ and a symbol $a\left(x^{\prime \prime}, y^{\prime \prime}, t\right) \in S_{\mathrm{cl}}^{n-\frac{d}{2}}\left(U \times U \times \mathbb{R}_{+}\right)$such that the following holds under the notation (3.76),

$$
\begin{equation*}
A_{1}(x, y) \equiv \int_{0}^{\infty} e^{i \Phi\left(x^{\prime \prime}, y^{\prime \prime}\right) t} a\left(x^{\prime \prime}, y^{\prime \prime}, t\right) d t \quad \text { on } U \times U \tag{4.98}
\end{equation*}
$$

with

$$
\begin{align*}
a_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)= & 2^{-3 n+4 d-1} \pi^{-n-1} \frac{1}{V_{\mathrm{eff}}\left(\underline{x}^{\prime \prime}\right)}\left|\operatorname{det} \mathscr{L}_{\underline{x}^{\prime \prime}}\right| \\
& \times\left|\operatorname{det} \mathcal{R}_{\underline{x}^{\prime \prime}}\right|^{-\frac{1}{2}}, \quad \text { for } \underline{x}^{\prime \prime} \in U \cap Y . \tag{4.99}
\end{align*}
$$

Theorem 4.14. Let $x, y \in Y$. If $\pi(x) \neq \pi(y)$, then there are neighborhoods $D_{G}$ of $\pi(x)$ and $V_{G}$ of $\pi(y)$ in $X_{G}$ such that

$$
\begin{equation*}
A_{2} \equiv 0 \quad \text { on } D_{G} \times V_{G} \tag{4.100}
\end{equation*}
$$

Let $p \in Y$ and let $x=\left(x_{1}, \ldots, x_{2 n+1}\right)$ be the local coordinates as in the discussion at the beginning of Sec. 4.1. Then there exists $\widehat{a}\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}, t\right) \in S_{\mathrm{cl}}^{n-d}\left(W \times W \times \mathbb{R}_{+}\right)$ such that by using the notations (3.76), (4.77), we have

$$
\begin{equation*}
A_{2}\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \equiv \int_{0}^{\infty} e^{i \Phi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) t} \widehat{a}\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}, t\right) d t \quad \text { on } W \times W \tag{4.101}
\end{equation*}
$$

with

$$
\begin{equation*}
\widehat{a}_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=2^{-3 n+\frac{5}{2} d-1} \pi^{-n+\frac{d}{2}-1}\left|\operatorname{det} \mathscr{L}_{X_{G}, \underline{x}^{\prime \prime}}\right|, \quad \text { for } \underline{x}^{\prime \prime} \in W . \tag{4.102}
\end{equation*}
$$

Set

$$
\begin{align*}
Q & :=-C_{0} \sigma^{*} \sigma+S_{G}:=-C_{0} A_{1}+S_{G} \\
\mathscr{C}^{\infty}(X) & \rightarrow H_{b}^{0}(X)^{G}, \quad \text { with } C_{0}=2^{3(n-d)} \pi^{d / 2} . \tag{4.103}
\end{align*}
$$

Since $A_{1}=A_{1} S_{G}=S_{G} A_{1}$, it is clear that

$$
\begin{equation*}
C_{0} A_{1}=S_{G}-Q=S_{G}-Q S_{G}=(I-Q) S_{G}=S_{G}(I-Q) \tag{4.104}
\end{equation*}
$$

and

$$
\begin{equation*}
Q^{*}=Q \tag{4.105}
\end{equation*}
$$

where $Q^{*}$ is the formal adjoint of $Q$. From Theorems 3.25, 4.11 and 4.13, we get:
Theorem 4.15. The operator $Q$ belongs to the class $G_{n-\frac{d}{2}, 1}(X)$ and hence $I-Q$ extends by density to a bounded self-adjoint operator $I-Q: L^{2}(X) \rightarrow L^{2}(X)$.

By Theorem4.15 there exists $C>0$ such that for every $u \in H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)$, we have

$$
\begin{equation*}
(\sigma u, \sigma u)_{X_{G}}=\left(\sigma^{*} \sigma u, u\right)=\frac{1}{C_{0}}((I-Q) u, u) \leq C\|u\|^{2} \tag{4.106}
\end{equation*}
$$

From (4.106), we deduce:
Corollary 4.16. There exists $C>0$ such that

$$
\begin{equation*}
(\sigma u, \sigma u)_{X_{G}} \leq C\|u\|^{2}, \quad \text { for any } u \in H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X) \tag{4.107}
\end{equation*}
$$

Hence we can extend $\sigma$ by density to a bounded operator

$$
\sigma: H_{b}^{0}(X)^{G} \rightarrow H_{b}^{0}\left(X_{G}\right)
$$

From now on, we consider $\sigma$ as a bounded operator $\sigma: H_{b}^{0}(X)^{G} \rightarrow H_{b}^{0}\left(X_{G}\right)$.
Theorem 4.17. Ker $\sigma$ is a finite-dimensional subspace of $H_{b}^{0}(X) \cap \mathscr{C}^{\infty}(X)$.
Proof. From Theorem4.11 we see that $\operatorname{Ker}(I-Q)$ is a finite-dimensional subspace of the space $\mathscr{C}^{\infty}(X)$. Note that $\operatorname{Ker} \sigma \subset H_{b}^{0}(X)^{G} \cap \operatorname{Ker}(I-Q)$, so the theorem follows.

From Theorem 4.14 we see that $\sigma \sigma^{*}$ is a complex Fourier integral operator with the same type as $S_{X_{G}}$. It is known that $\sigma \sigma^{*}$ is a pseudodifferential operator of order zero type $\left(\frac{1}{2}, \frac{1}{2}\right)$ (see [27, Proposition 5.18]). Set $C_{1}=\pi^{\frac{d}{2}} 2^{3 n-\frac{5}{2} d}$. Then
the leading of the symbol of $-C_{1} \sigma \sigma^{*}+S_{X_{G}}$ vanishes on the diagonal $x=y$, and it is known that $-C_{1} \sigma \sigma^{*}+S_{X_{G}}$ is a pseudodifferential operator of order $-\frac{1}{2}$ and type $\left(\frac{1}{2}, \frac{1}{2}\right)$ (see [27, Proposition 5.18]). By the classical Calderon-Vaillancourt theorem [26, Chap. XVIII], we have for every $s \in \mathbb{R}$,

$$
\begin{align*}
\sigma \sigma^{*}: H^{s}\left(X_{G}\right) & \rightarrow H^{s}\left(X_{G}\right) \quad \text { is continuous, } \\
-C_{1} \sigma \sigma^{*}+S_{X_{G}}: H^{s}\left(X_{G}\right) & \rightarrow H^{s+\frac{1}{2}}\left(X_{G}\right) \quad \text { is continuous. } \tag{4.108}
\end{align*}
$$

From (4.108), we have a result similar to Theorem 4.10 Hence, we can apply the proof of Corollary 4.16 with minor changes and deduce that there exists $\widehat{C}>0$ such that

$$
\begin{equation*}
\left(\sigma^{*} v, \sigma^{*} v\right) \leq \widehat{C}\|v\|_{X_{G}}^{2}, \quad \text { for any } v \in H_{b}^{0}\left(X_{G}\right) \cap \mathscr{C}^{\infty}\left(X_{G}\right) \tag{4.109}
\end{equation*}
$$

Therefore we can extend $\sigma^{*}$ by density to a bounded operator

$$
\sigma^{*}: H_{b}^{0}\left(X_{G}\right) \rightarrow H_{b}^{0}(X)^{G}
$$

We repeat the proof of Theorem 4.17 with minor changes and deduce:
Theorem 4.18. $\operatorname{Ker} \sigma^{*}$ is a finite-dimensional subspace of $H_{b}^{0}\left(X_{G}\right) \cap \mathscr{C}^{\infty}\left(X_{G}\right)$.
Finally, we obtain:
Theorem 4.19. $\operatorname{Ker} \sigma$ and $(\operatorname{Im} \sigma)^{\perp}$ are finite-dimensional subspaces of $H_{b}^{0}(X)^{G} \cap$ $\mathscr{C}^{\infty}(X)$ and $H_{b}^{0}\left(X_{G}\right) \cap \mathscr{C}^{\infty}\left(X_{G}\right)$, respectively.

Proof. We only need to prove that $(\operatorname{Im} \sigma)^{\perp}$ is a finite-dimensional subspace of $\mathscr{C}^{\infty}\left(X_{G}\right)$. Note that $(\operatorname{Im} \sigma)^{\perp} \subset \operatorname{Ker} \sigma^{*}$. From this observation and Theorem 4.18, the theorem follows.

Theorem 4.19 implies Theorem 1.5 .

## 5. Proof of Theorem 1.2

The main goal of this section is to prove Theorem 1.2, Recall that the Riemannian metrics $g^{T X}$ on $X$ and $g^{T X_{G}}$ on $X_{G}$ are given by Convention 2.8, Let $\Delta^{X}$ and $\Delta^{X_{G}}$ be the (positive) Laplacians on ( $X, g^{T X}$ ) and ( $X_{G}, g^{T X_{G}}$ ), respectively. For $s \in \mathbb{R}$ we consider the classical pseudodifferential operators of order $s$ on $X$ and $X_{G}$, respectively,

$$
\begin{equation*}
\Lambda_{s}=\left(1+\Delta^{X}\right)^{s / 2}, \quad \widehat{\Lambda}_{s}=\left(1+\Delta^{X_{G}}\right)^{s / 2} \tag{5.1}
\end{equation*}
$$

They are self-adjoint and positive with respect to the inner products $(\cdot, \cdot)_{X}$ and $(\cdot, \cdot)_{X_{G}}$, respectively. In particular, the maps

$$
\begin{equation*}
\Lambda_{s}: H^{\ell}(X) \rightarrow H^{\ell-s}(X), \quad \widehat{\Lambda}_{s}: H^{\ell}\left(X_{G}\right) \rightarrow H^{\ell-s}\left(X_{G}\right) \tag{5.2}
\end{equation*}
$$

are injective for any $\ell \in \mathbb{R}$. For $u, v \in H^{s}(X), u^{\prime}, v^{\prime} \in H^{s}\left(X_{G}\right)$, we define the inner products

$$
\begin{equation*}
(u, v)_{s}:=\left(\Lambda_{s} u, \Lambda_{s} v\right)_{X}, \quad\left(u^{\prime}, v^{\prime}\right)_{s}:=\left(\widehat{\Lambda}_{s} u, \widehat{\Lambda}_{s} v\right)_{X_{G}} \tag{5.3}
\end{equation*}
$$

and let $\|\cdot\|_{s},\|\cdot\|_{X_{G}, s}$ be the corresponding norms.
Recall that the space $G_{k, \ell}(X)$ is given by Definition 4.5 and $S_{G} \in G_{n-\frac{d}{2}, 0}(X)$.
Theorem 5.1. Let $A \in G_{n-\frac{d}{2}+\gamma, 0}(X), \gamma \in \mathbb{R}$. Then $A$ can be extended continuously to $A: H^{s}(X) \rightarrow H^{s-\gamma}(X)$ for every $s \in \mathbb{R}$.

Proof. For every $s \in \mathbb{R}$, put $B:=\Lambda_{s-\gamma} A \Lambda_{-s}$. As $\Lambda_{s}$ is $G$-invariant, we see from complex stationary phase formula of Melin-Sjöstrand that $B \in G_{n-\frac{d}{2}, 0}(X)$. Near $Y$, write $B(x, y) \equiv \int_{0}^{\infty} e^{i \Phi(x, y) t} b(x, y, t) d t$ as in 4.42) and let $b_{0}(x, y)$ be the leading term of $b(x, y, t)$. Let $P$ be a classical pseudodifferential operator of order 0 on $X$ with

$$
\begin{equation*}
\sigma_{P}\left(x, d_{x} \Phi(x, x)\right) a_{0}(x, x)=b_{0}(x, x) \tag{5.4}
\end{equation*}
$$

for every $x \in Y$, where $\sigma_{P}$ denotes the principal symbol of $P$ and $a_{0}$ is the leading term of the expansion of $S_{G}$ (see (3.85)). The complex stationary phase formula yields

$$
\begin{equation*}
B=P S_{G}+R, \quad R \in G_{n-\frac{d}{2}, 1}(X) \tag{5.5}
\end{equation*}
$$

From Theorem 4.10, (5.5) and the fact that $S_{G}$ is $L^{2}$ bounded, we deduce that $B=\Lambda_{s-\gamma} A \Lambda_{-s}$ is $L^{2}$ bounded. This implies that there exists $C_{1}>0$ such that for every $u \in \mathscr{C}^{\infty}(X)$,

$$
\begin{equation*}
\|A u\|_{s-\ell}=\left\|\Lambda_{s-\gamma} A \Lambda_{-s} \Lambda_{s} u\right\| \leq C_{1}\left\|\Lambda_{s} u\right\|=C_{1}\|u\|_{s} . \tag{5.6}
\end{equation*}
$$

The theorem follows.

From Theorems 3.27, 5.1 and note that $S_{G} \in G_{n-\frac{d}{2}, 0}(X)$, we deduce the following regularity property of the $G$-invariant Szegő projector.

Corollary 5.2. For every $s \in \mathbb{R}, S_{G}: H^{s}(X) \rightarrow H^{s}(X)$ is continuous, and in particular, $H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)$ is dense in $H_{b}^{0}(X)_{s}^{G}$ in $H^{s}(X)$.

Due to Theorem 3.27 the map $\sigma_{G}$ given by (1.5) has a well-defined extension

$$
\sigma_{G}: \mathscr{C}^{\infty}(X) \rightarrow H_{b}^{0}\left(X_{G}\right) \cap \mathscr{C}^{\infty}\left(X_{G}\right), \quad u \mapsto \iota_{G} \circ \iota^{*} \circ S_{G} u
$$

Theorem 5.3. For every $s \in \mathbb{Z}$ there exists $C_{s}>0$ such that

$$
\begin{equation*}
\left\|\sigma_{G} u\right\|_{X_{G}, s-\frac{d}{4}}^{2} \leq C_{s}\|u\|_{s}^{2}, \quad \text { for every } u \in \mathscr{C}^{\infty}(X) \tag{5.7}
\end{equation*}
$$

Moreover, the map (1.5) can be extended continuously by density to a bounded operator

$$
\begin{equation*}
\sigma_{G, s}:=\sigma_{G}: H_{b}^{0}(X)_{s}^{G} \rightarrow H_{b}^{0}\left(X_{G}\right)_{s-\frac{d}{4}} \tag{5.8}
\end{equation*}
$$

for every $s \in \mathbb{R}$.
Proof. Fix $s \in \mathbb{R}$. Let $\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow \mathscr{D}^{\prime}(X)$ be the formal adjoint of $\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}: \mathscr{C}^{\infty}(X) \rightarrow \mathscr{C}^{\infty}\left(X_{G}\right)$. For $u \in \mathscr{C}^{\infty}(X)$, we have

$$
\begin{align*}
\left\|\sigma_{G} u\right\|_{X_{G}, s-\frac{d}{4}}^{2} & =\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G} u, \quad \hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G} u\right)_{X_{G}} \\
& =\left(\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*}\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right) u, u\right) \tag{5.9}
\end{align*}
$$

We repeat the proof of Theorem 4.13 and conclude that

$$
\begin{equation*}
\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*}\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right) \in G_{n-\frac{d}{2}+2 s, 0}(X) \tag{5.10}
\end{equation*}
$$

From Theorem 5.1 and (5.10), we deduce that there exist $C, C_{1}>0$ such that

$$
\begin{align*}
\left|\left(\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*}\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right) u, u\right)\right| & \leq C\left\|\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*}\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right) u\right\|_{-s}\|u\|_{s} \\
& \leq C_{1}\|u\|_{s}^{2} \tag{5.11}
\end{align*}
$$

for every $u \in \mathscr{C}^{\infty}(X)$. From (5.9) and (5.11) we get (5.7).
By Theorem 3.27 and Corollary 5.2 we see that $H_{b}^{0}(X)^{G} \cap \mathscr{C}^{\infty}(X)$ is dense in $H_{b}^{0}(X)_{s}^{G}$, for every $s \in \mathbb{R}$. From this observation and (5.7), we deduce that for every $s \in \mathbb{R}$ the map (1.5) can be extended continuously by density to a bounded operator $\sigma_{G, s}$ as in (5.8).

Let $f_{G} \in \mathscr{C}^{\infty}(Y)^{G}$ be as in (1.19). We identify $f_{G}$ with a smooth function on $X_{G}$.

Theorem 5.4. For every $s \in \mathbb{R}$, $\operatorname{Ker} \sigma_{G, s}$ is a finite-dimensional subspace of $\mathscr{C}^{\infty}(X)$. Moreover, $\operatorname{Ker} \sigma_{G, s}$ is independent of $s$.

Proof. Let $E$ be a classical pseudodifferential operator on $X_{G}$ with principal symbol $\sigma_{E}(x, \xi)=|\xi|^{-\frac{d}{4}}$. Let for every $s \in \mathbb{R}$,

$$
\begin{equation*}
\hat{\sigma}_{G}:=S_{X_{G}} \circ E \circ f_{G} \circ \sigma_{G}: H_{b}^{0}(X)_{s}^{G} \rightarrow H_{b}^{0}\left(X_{G}\right)_{s} \tag{5.12}
\end{equation*}
$$

We repeat the proof of Theorem 4.17 and deduce that $\operatorname{Ker} \hat{\sigma}_{G}$ is a finite-dimensional subspace of $\mathscr{C}^{\infty}(X)$. Since $\operatorname{Ker} \sigma_{G} \subset \operatorname{Ker} \hat{\sigma}_{G}$, the theorem follows.

Theorem 5.5. With $\sigma_{G, s}$ as in (5.8), $\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$ in (1.7) is a finite-dimensional subspace of $\mathscr{C}^{\infty}\left(X_{G}\right)$ for every $s \in \mathbb{R}$.

Proof. Fix $s \in \mathbb{R}$. By Corollary 5.2 and Theorem 5.3 we can extend $\sigma_{G, s}$ in (5.8) to $H^{s}(X)$ by

$$
\begin{equation*}
\sigma_{G}: H^{s}(X) \rightarrow H^{s-\frac{d}{4}}\left(X_{G}\right), \quad u \mapsto \sigma_{G} S_{G} u . \tag{5.13}
\end{equation*}
$$

We have $\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}: \mathscr{C}^{\infty}(X) \rightarrow \mathscr{C}^{\infty}\left(X_{G}\right)$. Let $\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*}: \mathscr{D}^{\prime}\left(X_{G}\right) \rightarrow \mathscr{D}^{\prime}(X)$ be the formal adjoint of $\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}$. We repeat the proof of Theorem4.12 with minor changes and deduce that $\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*}: \mathscr{C}^{\infty}\left(X_{G}\right) \rightarrow \mathscr{C}^{\infty}(X)$. Let

$$
\begin{equation*}
F_{s}:=S_{X_{G}} f_{G}^{2} \hat{\Lambda}_{-2 s} \sigma_{G}\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*} \hat{\Lambda}_{s-\frac{d}{4}} S_{X_{G}}: \mathscr{D}^{\prime}\left(X_{G}\right) \rightarrow \mathscr{D}^{\prime}\left(X_{G}\right) \tag{5.14}
\end{equation*}
$$

For any $u \in\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$ and $v \in \mathscr{C}^{\infty}(X)$, by (1.7), we have

$$
\begin{align*}
\left(v, S_{G}\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G}\right)^{*} \hat{\Lambda}_{s-\frac{d}{4}} S_{X_{G}} u\right) & =\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G} S_{G} v, \hat{\Lambda}_{s-\frac{d}{4}} S_{X_{G}} u\right)_{X_{G}} \\
& =\left(\sigma_{G} v, u\right)_{X_{G}, s-\frac{d}{4}}=0 . \tag{5.15}
\end{align*}
$$

In view of (5.15), we see that

$$
\begin{equation*}
\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp} \subset \operatorname{Ker} F_{s} \cap H_{b}^{0}\left(X_{G}\right)_{s} \tag{5.16}
\end{equation*}
$$

We repeat the proof of Theorem 4.14 and conclude as in (4.104) that $F_{s}=C(I-$ $R) S_{X_{G}}$, where $C \neq 0$ is a constant and $R$ is a complex Fourier integral operator with the same phase and order as $S_{X_{G}}$ and vanishing leading term on the diagonal. More precisely, in the local coordinates $\underline{x}^{\prime \prime}$ of $X_{G}$ defined in an open set $W$ of $X_{G}$, we have

$$
\begin{equation*}
R\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) \equiv \int_{0}^{\infty} e^{i \Phi\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}\right) t} r\left(\underline{x}^{\prime \prime}, \underline{y}^{\prime \prime}, t\right) d t \tag{5.17}
\end{equation*}
$$

where $\Phi$ is as in (4.101), $r \in S_{\mathrm{cl}}^{n-d}\left(W \times W \times \mathbb{R}_{+}\right), r_{0}\left(\underline{x}^{\prime \prime}, \underline{x}^{\prime \prime}\right)=0$ for every $\underline{x}^{\prime \prime} \in W$, where $r_{0}$ is the leading term of $r$. We repeat now the proof of Theorem 4.11 with minor changes and deduce that $\operatorname{Ker}(I-R)$ is a finite-dimensional subspace of $\mathscr{C}^{\infty}\left(X_{G}\right)$. From (5.16) we deduce that $\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp} \subset \operatorname{Ker}(I-R)$. The theorem follows.

Theorem 5.6. $\operatorname{dim}\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$ is independent of $s$.
Proof. Fix $s \in \mathbb{R}$ and let $u \in\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$. By Theorem 5.5we have $u \in \mathscr{C}^{\infty}\left(X_{G}\right) \cap$ $H_{b}^{0}\left(X_{G}\right)$. We have an orthogonal decomposition

$$
\begin{equation*}
u=\sigma_{G, \frac{d}{4}} v+w, \quad v \in H_{b}^{0}(X)_{\frac{d}{4}}^{G}, w \in\left(\operatorname{Im} \sigma_{G, \frac{d}{4}}\right)^{\perp} \tag{5.18}
\end{equation*}
$$

We define a linear map

$$
\begin{equation*}
\gamma_{s}:\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp} \rightarrow\left(\operatorname{Im} \sigma_{G, \frac{d}{4}}\right)^{\perp}, \quad \gamma_{s} u:=w \in\left(\operatorname{Im} \sigma_{G, \frac{d}{4}}\right)^{\perp} \tag{5.19}
\end{equation*}
$$

We claim that $\gamma_{s}$ is injective. Assume that $\gamma_{s} u=0$. Then there exists $v \in H_{b}^{0}(X)_{\frac{d}{4}}^{G}$ such that $u=\sigma_{G, \frac{d}{4}} v$. Let $\left(v_{j}\right)$ be a sequence in $\mathscr{C}^{\infty}(X)$ with $v_{j} \rightarrow v$ in $H^{\frac{d}{4}}(X)$ as $j \rightarrow \infty$. By Corollary 5.2 we have $S_{G} v_{j} \rightarrow S_{G} v=v$ in $H^{\frac{d}{4}}(X)$ as $j \rightarrow \infty$.

By Theorem 5.3 we have $\sigma_{G, \frac{d}{4}} S_{G} v_{j} \rightarrow \sigma_{G, \frac{d}{4}} v=u$ in $L^{2}\left(X_{G}\right)$ as $j \rightarrow \infty$. Since $u \in\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$, we have as $j \rightarrow \infty$
$0=\left(\hat{\Lambda}_{s-\frac{d}{4}} \sigma_{G} S_{G} v_{j}, \quad \hat{\Lambda}_{s-\frac{d}{4}} u\right)_{X_{G}}=\left(\sigma_{G} S_{G} v_{j},\left(\hat{\Lambda}_{s-\frac{d}{4}}\right)^{2} u\right)_{X_{G}} \rightarrow\left(u,\left(\hat{\Lambda}_{s-\frac{d}{4}}\right)^{2} u\right)_{X_{G}}$.
Hence $\left(u,\left(\hat{\Lambda}_{s-\frac{d}{4}}\right)^{2} u\right)_{X_{G}}=0$ and thus $\hat{\Lambda}_{s-\frac{d}{4}} u=0$. Since $\hat{\Lambda}_{s-\frac{d}{4}}$ is injective, we get $u=0$ and hence $\gamma_{s}$ is injective. Since $\gamma_{s}$ is injective, $\operatorname{dim}\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp} \leq$ $\operatorname{dim}\left(\operatorname{Im} \sigma_{G, \frac{d}{4}}\right)^{\perp}$.

Similarly, we can repeat the procedure above and conclude that dim (Im $\left.\sigma_{G, \frac{d}{4}}\right)^{\perp} \leq \operatorname{dim}\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}$. Thus, $\operatorname{dim}\left(\operatorname{Im} \sigma_{G, s}\right)^{\perp}=\operatorname{dim}\left(\operatorname{Im} \sigma_{G, \frac{d}{4}}\right)^{\perp}$. The theorem follows.

Proof of Theorem 1.2, Theorems 5.4, 5.5 and 5.6 yield Theorem 1.2 for the case when $\operatorname{dim} X_{G} \geq 3$.

Assume now that $\operatorname{dim} X_{G}=1$. Then $X_{G}$ is a union of circles. For simplicity, suppose that $X_{G}=S^{1}$. The circle $X_{G}$ admits a natural $S^{1}$ action $e^{i \theta}: S^{1} \times X_{G} \rightarrow$ $X_{G},\left(e^{i \theta}, z\right) \rightarrow e^{i \theta} z$. For every $m \in \mathbb{Z}$, put

$$
L_{m}^{2}\left(X_{G}\right):=\left\{u \in L^{2}\left(X_{G}\right):\left(e^{i \theta}\right)^{*} u=e^{i m \theta} u, \text { for every } e^{i \theta} \in S^{1}\right\}
$$

It is clear that $L^{2}\left(X_{G}\right)=\oplus_{m \in \mathbb{Z}} L_{m}^{2}\left(X_{G}\right)$. By definition $H_{b}^{0}\left(X_{G}\right)=\oplus_{m \in \mathbb{N}} L_{m}^{2}\left(X_{G}\right)$. The Szegő projection on $X_{G}$ is the orthogonal projection: $S_{X_{G}}$ : $L^{2}\left(X_{G}\right) \rightarrow H_{b}^{0}\left(X_{G}\right)$. The $S^{1}$ action induces a smooth vector field $\frac{\partial}{\partial \theta}$ on $X_{G}$. Fix a point $p \in X_{G}$. Let $x$ be local coordinate of $X_{G}$ such that $x(p)=0$ and $\frac{\partial}{\partial x}=\frac{\partial}{\partial \theta}$. Then

$$
\begin{equation*}
S_{X_{G}}(x, y) \equiv \frac{1}{2 \pi} \int_{0}^{\infty} e^{i t(x-y)} d t \tag{5.20}
\end{equation*}
$$

In particular, $S_{X_{G}}(x, y)$ is a Fourier integral operator with complex phase. Therefore, the above proof of Theorem 1.2 in the case $\operatorname{dim} X_{G} \geq 3$ works also when $\operatorname{dim} X_{G}=1$.

Theorem 5.7. Let $X$ be a three-dimensional compact orientable $C R$ manifold and let $G$ be a compact Lie group acting on $X$ such that the $G$-action preserves $J$ and $\omega_{0}$. We assume that $X$ is pseudoconvex of finite type and that $\bar{\partial}_{b}$ has closed range in $L^{2}$ on $X$. Then the conclusions of Theorem 1.2 hold.

Proof. It was shown in [12, Proposition 4.1] that there exists a bounded linear operator $\mathcal{G}: \operatorname{Im} \bar{\partial}_{b} \rightarrow L^{2}(X)$ such that $S=I-\mathcal{G} \bar{\partial}_{b}$. Furthermore $\mathcal{G}$ is smoothing away the diagonal and maps smooth functions to smooth functions, where $S$ denotes the Szegő projector. Hence, the Szegő projector is smoothing outside the diagonal and preserves the space of smooth functions. Moreover, $G$ is a circle and $\operatorname{dim} X_{G}=1$ in this case, thus the arguments above apply again.

Example 5.8. If $X$ be a compact pseudoconvex three-dimensional CR manifold of finite type admitting a transversal CR circle action, then $\bar{\partial}_{b}$ has closed range in $L^{2}$,
see [31, $\S 5.2$. Let $M$ be a compact Riemann surface and ( $L, h^{L}$ ) be semi-positive line bundle over $M$ whose curvature $R^{L}$ vanishes to finite order at any point. Then the Grauert tube $X$ of $L^{*}(2.19)$ is a compact pseudoconvex three-dimensional CR manifold of finite type (cf. [31, [41, Proposition 11]) admitting a transversal CR circle action. If $G$ is a compact Lie group acting holomorphically on $M$ and whose action lifts to $\left(L, h^{L}\right)$, Theorem 5.7 applies to $X$.

## 6. An Almost Complex Version of Theorem 1.4

In this section, we will prove a version of Theorem 1.4 in the case of an almost complex manifold. We will mostly follow [52] and adopt the notations therein.

Let $\left(L, h^{L}\right)$ be a Hermitian line bundle with Hermitian connection $\nabla^{L}$ and associated curvature $R^{L}$ on a compact almost complex manifold $(M, J)$.

Let $G$ be a compact Lie group with Lie algebra $\mathfrak{g}$ acting (on left) on $M$, whose action lifts on $L$ such that $h^{L}, \nabla^{L}$ are $G$-equivariant. Then the moment map $\mu: M \rightarrow \mathfrak{g}^{*}$ is defined by the Kostant formula

$$
\begin{equation*}
2 \pi \sqrt{-1}\langle\mu, \xi\rangle:=\nabla_{\xi_{M}}^{L}-L_{\xi}, \quad \text { for } \xi \in \mathfrak{g} . \tag{6.1}
\end{equation*}
$$

For any $\xi \in \mathfrak{g}$, we have

$$
\begin{equation*}
d\langle\mu, \xi\rangle=i_{\xi_{M}} \omega, \quad \text { with } \omega=\frac{\sqrt{-1}}{2 \pi} R^{L} . \tag{6.2}
\end{equation*}
$$

We assume that the almost complex structure $J$ on $M$ is $G$-invariant and $R^{L}$ is $J$-invariant, $G$ acts freely on $\mu^{-1}(0)$ and $\omega(\cdot, J \cdot)$ defines a metric on $\left.T M\right|_{\mu^{-1}(0)}$.

By choosing any $G$ - and $J$-invariant Riemannian metric $g^{T M}$ on $T M$, we can define an associated Dirac operator $D^{L}$ on $\Lambda\left(T^{*(0,1)} M\right) \otimes L$ and $D_{ \pm}^{L}$ its restriction on $\Omega^{0, \text { even/odd }}(M, L)$ (cf. [52, Definition 1.1]). Its index as a finite-dimensional virtual representation of $G$,

$$
\begin{equation*}
\operatorname{Ind}\left(D_{+}^{L}\right)=\operatorname{Ker}\left(D_{+}^{L}\right)-\operatorname{Ker}\left(D_{-}^{L}\right) \in R(G) \tag{6.3}
\end{equation*}
$$

does not depend on the choice of $g^{T M}$.
Moreover $\left(L, h^{L}, \nabla^{L}\right), J, \omega$ on $M$ induce canonically $\left(L_{G}, h^{L_{G}}, \nabla^{L_{G}}\right), J_{G}, \omega_{G}$ on $M_{G}=\mu^{-1}(0) / G$. In particular, $\left(M_{G}, \omega_{G}\right)$ is a compact symplectic manifold with compatible almost complex structure $J_{G}$. Thus $\operatorname{Ind}\left(D_{+}^{L_{G}}\right)$ is well defined as a virtual vector space.

Theorem 6.1. Let $(M, J)$ be a compact almost complex manifold and $\left(L, h, \nabla^{L}\right)$ be a Hermitian line bundle with connection on $M$. Let $G$ be a compact Lie group acting (on left) on $M$, whose action lifts on $L$ such that $J, h^{L}$ and $\nabla^{L}$ are $G$ equivariant. We assume that $G$ acts freely on $\mu^{-1}(0)$ and $\omega(\cdot, J \cdot)$ defines a metric on $\left.T M\right|_{\mu^{-1}(0)}$. Then there exists $m_{0} \in \mathbb{N}$ such that for any $m \geq m_{0}$ we have

$$
\begin{equation*}
\operatorname{Ind}\left(D_{+}^{L^{m}}\right)^{G}=\operatorname{Ind}\left(D_{+}^{L_{G}^{m}}\right) . \tag{6.4}
\end{equation*}
$$

Note that by the main result of 44, [52, Theorem 0.1 ], if $\omega(\cdot, J \cdot)>0$ on the whole $M$, then Theorem 6.1 holds for $m_{0}=1$.

Proof. We adapt directly the notation and argument from 52. We fix a $G$ - and $J$-invariant metric $g^{T M}$ on $T M$ such that near $\mu^{-1}(0), g^{T M}$ is given by $\omega(\cdot, J \cdot)$. We fix an $\operatorname{Ad}_{G}$-invariant scalar product on $\mathfrak{g}$ and identify $\mathfrak{g}$ and $\mathfrak{g}^{*}$ via this product.

Let $h_{1}, \ldots, h_{d}$ be an orthonormal basis of $\mathfrak{g}$. Set
$X^{\mathcal{H}}(x)=2(\mu(x))_{M}(x)=2 \sum_{i} \mu_{i}(x) V_{i}(x)$ with $\mu=\sum_{i} \mu_{i}(x) h_{i}$ and $V_{i}=h_{i, M}$.

Following [52, Definition 1.2] we set for $T \in \mathbb{R}$,

$$
\begin{equation*}
D_{T}^{L^{m}}=D^{L^{m}}+\frac{\sqrt{-1}}{2} T c\left(X^{\mathcal{H}}\right): \Omega^{0, *}(M, L) \rightarrow \Omega^{0, *}(M, L) \tag{6.6}
\end{equation*}
$$

where $c(\cdot)$ is the Clifford action. Then we have (cf. [52, (1.26)])

$$
\begin{equation*}
\left(D_{T}^{L^{m}}\right)^{2}=\left(D^{L^{m}}\right)^{2}+\frac{\sqrt{-1}}{2} T \sum_{j} c\left(e_{j}\right) c\left(\nabla_{e_{j}} X^{\mathcal{H}}\right)-\sqrt{-1} T \nabla_{X \mathcal{H}}+\frac{T^{2}}{4}\left|X^{\mathcal{H}}\right|^{2} \tag{6.7}
\end{equation*}
$$

where $\left\{e_{j}\right\}_{j}$ is an orthonormal frame of $\left(T M, g^{T M}\right)$. Now as in [52, (1.27)],

$$
\begin{equation*}
\nabla_{X \mathcal{H}}=2 \sum_{i} \mu_{i} L_{V_{i}}+4 m \sqrt{-1} \mathcal{H}+A \tag{6.8}
\end{equation*}
$$

where $A$ is an endomorphism of $\Lambda\left(T^{*(0,1)} M\right)$ and does not depend on $m$.
We fix a sufficiently small $G$-invariant open neighborhoods $U^{\prime} \Subset U$ of $\mu^{-1}(0)$ such that $G$ acts freely on $U$ and $\omega(\cdot, J \cdot)>0$ on $U$. Note that on $\Omega^{0, *}(M, L)^{G}$, $L_{V_{i}}=0$. From (6.7) and (6.8), there exists $m_{0}>0$ such that for any $m \geq m_{0}$, the assertion of [52, Theorem 2.1] holds for $U^{\prime}$ : there exist $C>0, b>0$ such that for any $T \geq 1$ and any $s \in \Omega^{0, *}\left(M, L^{m}\right)$ with $\operatorname{supp} s \subset M \backslash U^{\prime}$, we have

$$
\begin{equation*}
\left\|D_{T}^{L^{m}} s\right\|_{0}^{2} \geq C\left(\|s\|_{1}^{2}+(T-b)\|s\|_{0}^{2}\right) . \tag{6.9}
\end{equation*}
$$

Thus we are on $U$ exactly in same situation as considered in [52, §3(b)-3(e)]. Hence Theorem 6.1 follows as in [52, (3.36), (3.37)].

It is an interesting question to show that in the holomorphic situation, under the assumption of this section, we have

$$
\begin{equation*}
H^{j}\left(M, L^{m}\right)^{G}=0 \quad \text { for any } j>0, m \gg 1 \tag{6.10}
\end{equation*}
$$

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